

Electronic Supplementary Information for

## Size dependence in two-dimensional lateral heterostructures of transition metal dichalcogenides

Hao Jin,<sup>a</sup> Vincent Michaud-Rioux,<sup>b</sup> Zhi-Rui Gong,<sup>a</sup> Langhui Wan,<sup>a</sup> Yadong Wei,<sup>\*a</sup> and  
Hong Guo,<sup>a,b</sup>

<sup>a</sup> Shenzhen Key Laboratory of Advanced Thin Films and Applications, College of  
Physics and Optoelectronic Engineering, Shenzhen University, Shenzhen 518060,  
China; E-mail: ywei@szu.edu.cn

<sup>b</sup> Centre for the Physics of Materials and Department of Physics, McGill University,  
Montréal H3A 2T8, Canada.

For armchair (zigzag) boundary, it is one-unit length along  $x$  ( $y$ ) direction (see Fig. 1a  
of the manuscript), i.e. 5.51 Å (3.18 Å). While the length along  $y$  ( $x$ ) direction varies.  
To better illustrate the size of the supercell employed in this work, we list the length  
and the number of atoms in Table S1 and S2.

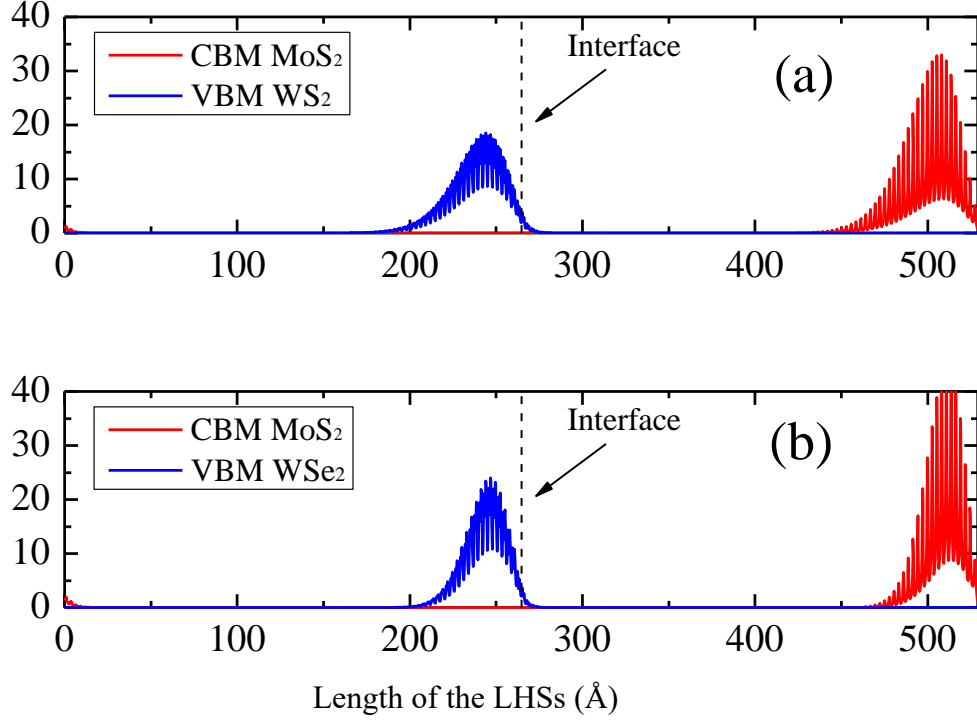
**Table S1.** The length (as labelled in Fig. 1 of the manuscript) and the number of atoms  
employed in the supercell along armchair direction.

Length (Å)	Number of atoms
19.1	36
38.2	72
57.3	108
76.4	144
114.6	216
152.8	288
191.0	360
229.2	432
305.5	576

<b>343.7</b>	<b>648</b>
<b>458.3</b>	<b>864</b>

**Table S2.** The length (as labelled in Fig. 1 of the manuscript) and the number of atoms employed in the supercell along zigzag direction.

<b>Length (Å)</b>	<b>Number of atoms</b>
<b>44.1</b>	<b>48</b>
<b>88.2</b>	<b>96</b>
<b>132.3</b>	<b>144</b>
<b>176.4</b>	<b>192</b>
<b>264.6</b>	<b>288</b>
<b>352.8</b>	<b>384</b>
<b>441.0</b>	<b>480</b>
<b>529.2</b>	<b>576</b>
<b>705.6</b>	<b>768</b>
<b>882.0</b>	<b>960</b>
<b>1058.4</b>	<b>1152</b>
<b>1411.2</b>	<b>1536</b>
<b>1764.0</b>	<b>1920</b>
<b>2116.8</b>	<b>2304</b>
<b>2756.3</b>	<b>3000</b>
<b>3528.0</b>	<b>3840</b>
<b>4233.6</b>	<b>4608</b>



**Fig. S1.** Plane integrated modular square wave functions of the VBM ( $|\psi_v(r)|^2$ ) and CBM ( $|\psi_c(r)|^2$ ) for (a) MoS<sub>2</sub>/WS<sub>2</sub> and (b) MoS<sub>2</sub>/WSe<sub>2</sub> LHSs.

As shown in Fig. R1, the plane integrated modular square wave functions of the VBM ( $|\psi_v(r)|^2$ ) and CBM ( $|\psi_c(r)|^2$ ) of the MoS<sub>2</sub>/WSe<sub>2</sub> LHS are similar to those of MoS<sub>2</sub>/WS<sub>2</sub> LHS. As have been discussed in the manuscript, those wave functions determine the band structures of the LHSs. As such, we believe that MoS<sub>2</sub>/WSe<sub>2</sub> LHS can show similar trend as compared with MoS<sub>2</sub>/WS<sub>2</sub> LHS. Consequently, our conclusions are not limited to MoS<sub>2</sub>/WS<sub>2</sub> LHS, but can be applied to other TMDs LHSs.