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## Supporting Information

Title: All-organic Room Temperature Thermally Switchable Dielectric System

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Layer	${\cal E}'_i$	$\sigma_i / S cm^{-1}$	d <sub>i</sub> / mm
I	2.5	3.8X10 <sup>-14</sup>	0.056
II	4.5	8.5X10 <sup>-13</sup>	0
III	4.8	3.1X10 <sup>-9</sup>	4.838

**Table S1.** The initial parameters of each layer in the system.



**Figure S1**. The (a) real and (b) imaginary dielectric spectrum of the system at different stages. The dark line, red line and blue line represent the liquid state, phase transition and solid state, respectively. (c) and (d) are the magnified plot of the relaxation  $\gamma$  in Figure S1 (a) and Figure S1 (b).

Here, the thickness of the insulating film  $(d_1)$  is  $10^{-7}$  mm.

The strong relaxation  $\alpha$  comes from the interface between the insulating film and the liquid mixture (Figure S1(a)). During the phase transition, the relaxation  $\alpha$  is replaced by the relaxation  $\beta$  that stems from the interface between the solid mixture and the liquid mixture (Figure S1(b)). Whereas the relaxation  $\gamma$  caused by the interface between the insulating film and the solid mixture is too small that it is overlapped by the relaxation  $\beta$ . At the solid mixture thickness ( $d_2$ ) of 0.001 mm, the strength of the relaxation  $\beta$  boosts the dielectric constant of the system, which is responsible to the peak trend in Figure 1(b). As the solid mixture becomes thicker, the strength of the relaxation  $\beta$  declines. After complete solidification, the relaxation  $\beta$  vanishes and the tiny strength of the relaxation  $\gamma$  is shown in Figure S1(c) and (d).



**Figure S2.** Influence of the parameters of (a)-(b) the insulating layer and (c)-(e) the liquid layer on the dielectric constant of the system ( $\varepsilon'_s$ ) and the frequency of the relaxation  $\alpha$  ( $f_\alpha$ ). Here, the simulation focuses on the first stage of the system.



**Figure S3.** The influence of the conductivity of the insulating film ( $\sigma_1$ ) on (a) the dielectric constant of the system ( $\varepsilon'_s$ ) and (b) the frequency of the relaxation  $\alpha$  ( $f_\alpha$ ) at the first stage.



**Figure S4.** (a) Dependence of the dielectric constant of the system ( $\varepsilon'_s$ ), the frequency of the relaxation  $\gamma$  ( $f_{\gamma}$ ) and on the thickness of layer I ( $d_1$ ) after the whole mixture becomes solid. (b) Dependence of the switching ratio on the thickness of layer I ( $d_1$ )

In the first stage, the "three-layer" model degenerates into the "two-layer" model<sup>1</sup> composed by layer I and layer III. The dielectric constant ( $\varepsilon'_s$ ) and relaxation frequency (f) of the system is determined by the following equations,<sup>2</sup>

$$\varepsilon_{s}^{'} = d \left( \frac{d_{1}\varepsilon_{1}}{\sigma_{1}^{2}} + \frac{d_{3}\varepsilon_{3}}{\sigma_{3}^{2}} \right) / \left( \frac{d_{1}}{\sigma_{1}} + \frac{d_{3}}{\sigma_{3}} \right)^{2}$$
(S1)  
$$f = \frac{1}{2\pi\varepsilon_{0}} \left( \frac{\sigma_{1}}{d_{1}} + \frac{\sigma_{3}}{d_{3}} \right) / \left( \frac{\varepsilon_{1}^{'}}{d_{1}} + \frac{\varepsilon_{3}^{'}}{d_{3}} \right)$$
(S2)

where  $\varepsilon_0$  is the dielectric constant of vacuum. According to Equation S1 and S2, the dielectric

constant of the system ( $\varepsilon'_s$ ) and the frequency of the relaxation  $\alpha$  ( $f_\alpha$ ) are profoundly affected by the dielectric constant ( $\varepsilon'_1$ ) and conductivity ( $\sigma_1$ ) of the insulating film (Figure S2(a) and (b)). Specifically, as  $\varepsilon_l$  increases from 1 to 10,  $\varepsilon'_s$  increases from 44.9 to 449.3 in an approximately linear manner, and  $f_{\alpha}$  decreases concavely from 60.8 to 6.4 as shown in Figure S2 (a). Figure S2 (b) shows the dependence of  $\varepsilon'_s$  and f on  $\sigma_l$ . With the increasing of  $\sigma_l$ ,  $\varepsilon'_s$  begins at a plateau value of 218.5 ( $\sigma_l$  is 1.0X10<sup>-15</sup> S cm<sup>-1</sup>) and quickly goes down to its low plateau of 4.9 ( $\sigma_l$  is 1.0X10<sup>-7</sup> S cm<sup>-1</sup> <sup>1</sup>). During the descending, it is interesting to find that  $\varepsilon'_s$  attains a minimum of 4.75 when  $\sigma_1$  equals to 1.6X10<sup>-9</sup> S cm<sup>-1</sup>, as indicated in Figure S3(a)). Whereas  $f_{\alpha}$  changes with the increasing of  $\sigma_1$  in a different way. Initially, f is as low as 25.1 Hz ( $\sigma_1$  is 1.s0X10<sup>-15</sup> S cm<sup>-1</sup>), and then it holds a straight increasing when  $\sigma_1$  is above 1.0X10<sup>-11</sup> S cm<sup>-1</sup>, as Figure S2(b) and Figure S3(b) show. The liquid mixture (layer III) also plays an important role in determining the dielectric property of the system. As shown in Figure S2(c),  $\varepsilon'_s$  and  $f_\alpha$  is positively correlated with the thickness of the liquid mixture  $(d_3)$ . Specifically, as  $d_3$  increases,  $\varepsilon'_s$  maintains a low plateau of 2.5 at first, and then keep a continuous growth after  $d_3$  is larger than 1X10<sup>-2</sup> mm). While  $f_{\alpha}$  goes up from the low plateau of 0.027 Hz ( $d_3$  is 1X10<sup>-7</sup> mm) and reaches the plateau value of 1125.6 Hz at the thickness of 100 mm. Contrarily,  $\varepsilon'_s$  shows strong immunity and  $f_{\alpha}$  suffers a slight falloff to the variation of the dielectric constant of the liquid mixture ( $\varepsilon'_3$ ) (Figure S2(d)). Besides, by changing the conductivity of the liquid mixture ( $\sigma_3$ ), there is a dramatic rise in both  $\varepsilon'_s$  and  $f_\alpha$  (Figure S2(e)). Concretely, as  $\sigma_3$ increases,  $\varepsilon'_s$  is enhanced from 4.9 ( $\sigma_3$  is 1.0X10<sup>-16</sup> S cm<sup>-1</sup>) to 218.5 ( $\sigma_3$  is 1.0X10<sup>-7</sup> S cm<sup>-1</sup>), and  $f_{\alpha}$ keeps rising up when  $\sigma_3$  is beyond 1.0X10<sup>-12</sup> S cm<sup>-1</sup>.

In the third stage where the mixture totally turns into solid, the dielectric constant of the system goes into the low state as shown in Figure 1(b). The system is composed by the insulating film (layer I) and the solid mixture (layer II). The dielectric constant of the system ( $\varepsilon'_s$ ) and the frequency of the relaxation  $\gamma$  ( $f_{\gamma}$ ) can be simulated by the following equations:

$$\varepsilon_{s} = d \left( \frac{d_{1}\varepsilon_{1}}{\sigma_{1}^{2}} + \frac{d_{2}\varepsilon_{2}}{\sigma_{2}^{2}} \right) / \left( \frac{d_{1}}{\sigma_{1}} + \frac{d_{2}}{\sigma_{2}} \right)^{2}$$

$$f = \frac{1}{2\pi\varepsilon_{0}} \left( \frac{\sigma_{1}}{d_{1}} + \frac{\sigma_{2}}{d_{2}} \right) / \left( \frac{\varepsilon_{1}}{d_{1}} + \frac{\varepsilon_{2}}{d_{2}} \right)$$
(S3)
(S4)

According to Equation S3 and S4, by altering the thickness of the insulating film ( $d_1$ ),  $\varepsilon'_s$  increases from 4.5 ( $d_1$  is 1X10<sup>-7</sup> mm) and peaks at 15.8 ( $d_1$  is 0.2 mm), then it decreases rapidly to

2.6 when  $d_1$  equals to 100 mm as shown in Figure S4. While  $f_{\gamma}$  monotonically increases from 0.027 Hz ( $d_1$  is 1X10<sup>-7</sup> mm) to 0.332Hz ( $d_1$  is 100 mm).



**Figure S5.** Dependence of (a) the imaginary part of the dielectric constant and (b) the dielectric loss of the system between 50 °C and -30 °C in the isochronous measurement (1Hz, 2 K min<sup>-1</sup>).

References:

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