

An azido-bridged disc-like heptanuclear cobalt(II) cluster: towards single-molecule magnet

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Supporting Materials

§ [Magnetic Analysis]

The exchange Hamiltonian for a S_6 symmetric heptanuclear model is

$$H = -2J_{\text{intra}} \left(\sum_{i=2}^7 S_1 \cdot S_i + \sum_{j=2}^6 S_j \cdot S_{j+1} + S_2 \cdot S_7 \right) + g_0 \mu_B H_0 \sum_{i=1}^7 S_i^z \quad (\text{S1})$$

where J_{intra} is the exchange integral between adjacent Co atoms. Then, based on the Kambe's method,

$$H = -2J_{\text{intra}} \left\{ \frac{1}{2} \sum_{i=2}^7 (S_{1i}^2 - S_1^2 - S_i^2) + \frac{1}{2} \sum_{j=2}^6 (S'^2 - S_j^2 - S_{j+1}^2) + \frac{1}{2} (S_{27}^2 - S_2^2 - S_7^2) \right\} + g_0 \mu_B H_0 \sum_{i=1}^7 S_i^z \quad (\text{S2})$$

where each S_i is $1/2$, $S_{1i} = S_1 + S_i$, $S' = S_j + S_{j+1}$, $S_{27} = S_2 + S_7$. Thus, the energy expression is easy to get:

$$E = -J_{\text{intra}} \left\{ \sum_{i=2}^7 S_{1i}(S_{1i} + 1) + S_{23}(S_{23} + 1) + L + S_{72}(S_{72} + 1) \right\} \quad (\text{S3})$$

where $S_{ij} = S_i + S_j$, $S_i + S_{j-1}$, ..., $|S_i - S_j|$, here $S_{ij} = 1$ or 0 . By solving the exchange matrix, the molecular magnetic susceptibility expression in the van Vleck equation was shown in eq S4:

$$\chi_m = \left(\frac{Ng^2\beta^2}{3kT} \right) \left(\frac{A}{B} \right) \quad (\text{S4})$$

Here N is Avogadro's number, g the Landé g factor, β the electron Bohr magneton, k the Boltzmann constant, T the temperature, and,

$$\begin{aligned} A &= 126 \exp(24y) + 315 \exp(18y) + 90 \exp(14y) + 187.5 \exp(12y) + 12 \exp(10y) + 108 \exp(8y) \\ &\quad + 12 \exp(6y) + 13.5 \exp(4y); \\ B &= 8 \exp(24y) + 36 \exp(18y) + 24 \exp(14y) + 42 \exp(12y) + 12 \exp(10y) + 48 \exp(8y) \\ &\quad + 16 \exp(6y) + 18 \exp(4y) \end{aligned}$$

with $y = \frac{J_{\text{intra}}}{kT}$;

If intermolecular interaction (J_{inter}) are considered by mean field approximation, then:

$$\chi_m T = \frac{\chi_m T}{1 - 2J_{\text{inter}}\chi_m / (N\beta^2 g^2)} \quad (\text{S5})$$

The experimental data (2–30 K) were fitted to eq 4, and the best fitting of the data gave:

$$J_1 = 10.4 \text{ cm}^{-1}, J_2 = -0.05 \text{ cm}^{-1}, g = 5.9, \text{ with } R = 1.5 \times 10^{-5}; \{R = \sum[(\chi_M T)_{\text{obs}} - (\chi_M T)_{\text{calc}}]^2 / \sum(\chi_M T)_{\text{obs}}^2\}.$$

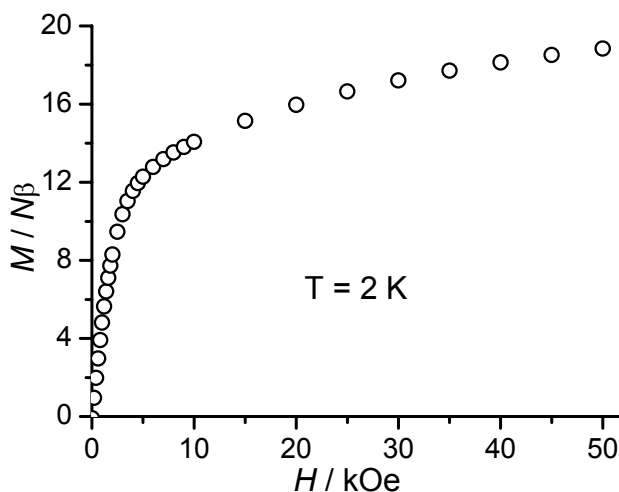


Figure S1 *dc* field-dependence of magnetization of **1** at 2 K..

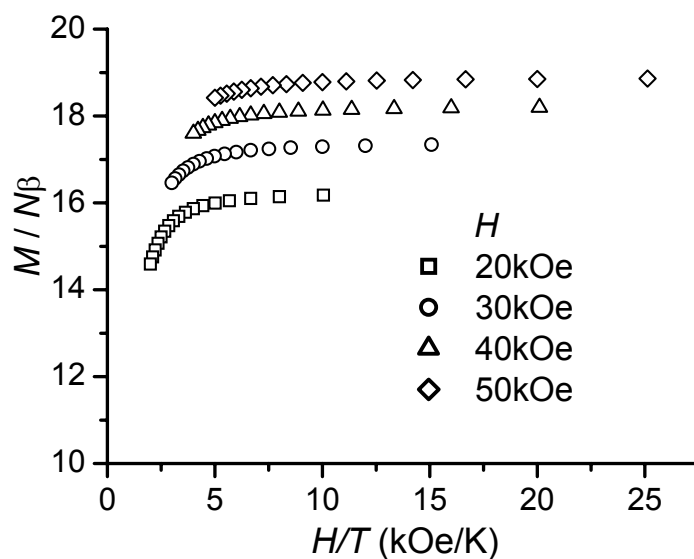


Figure S2 Magnetization data in applied fields from 20 to 50 kOe at temperatures between 2 and 10K.

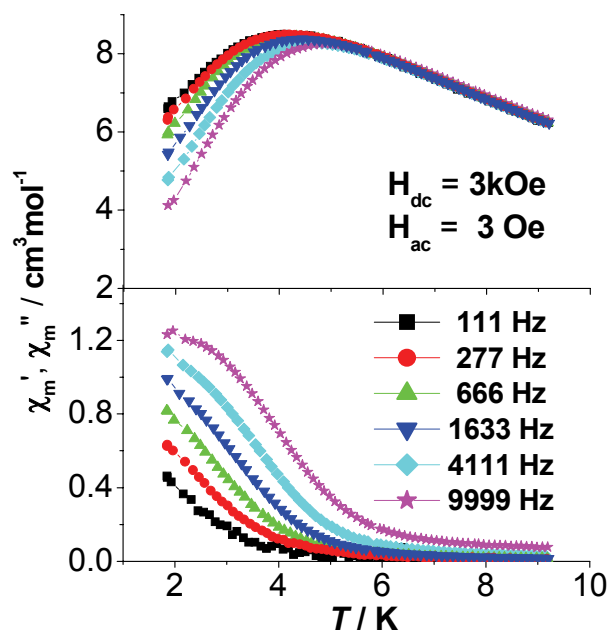


Figure S3 Temperature dependence of the real (top) and imaginary (bottom) components of the *ac* susceptibility for **1** with an oscillating field 3 Oe in frequency of 111-9999 Hz in an applied static field of 3 kOe.

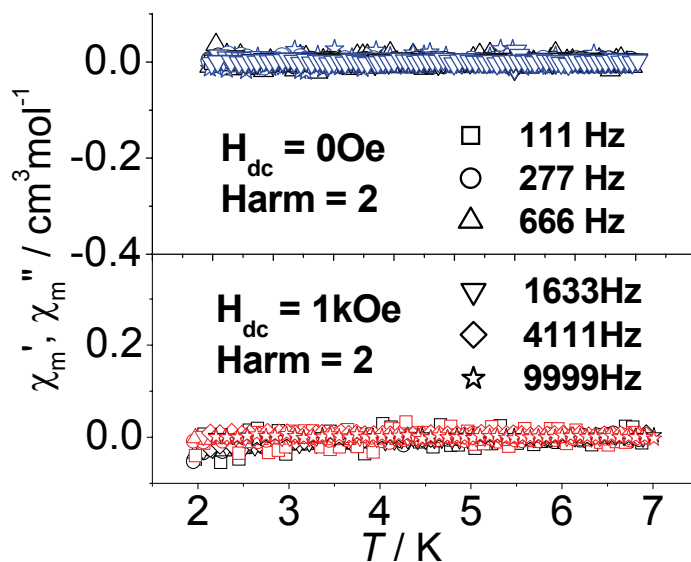
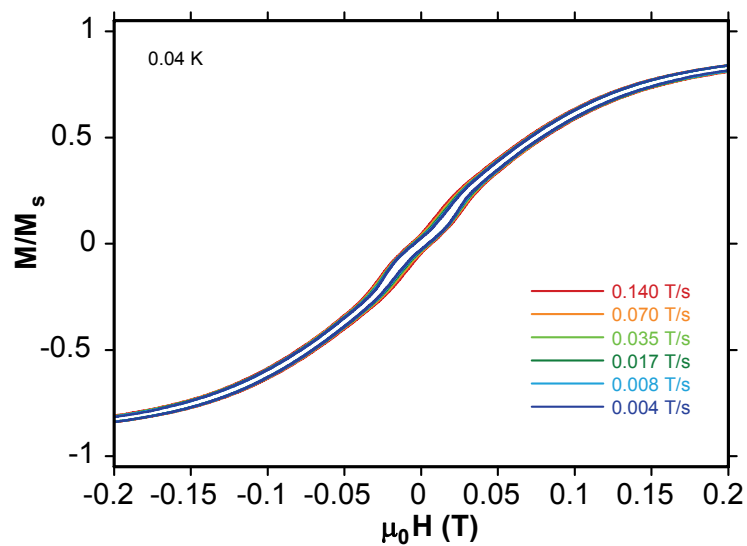


Figure S4 Temperature dependence of the real and imaginary components of the nonlinear second harmonic *ac* susceptibility for **1** in zero applied static field (up) and 1 kOe (down) with an oscillating field 5 Oe in frequency of 111-9999 Hz.



(b)

Figure S5 Single crystal magnetization (M) versus field (H) hysteresis loops for **1** at the indicated field sweep rates and temperatures. The small features at low fields (0.03 T) are probably due to tiny antiferromagnetic interactions between molecules.