

Supporting Information for

A Unique Substituted Co(II)–Formate Coordination Framework Exhibits Weak Ferromagnetic Single-Chain-Magnet Like Behavior

Jiong-Peng Zhao,^{a,b} Qian Yang,^a Zhong-Yi Liu,^c Ran Zhao,^a Bo-Wen Hu,^a Miao Du,^c Ze Chang,^a

and Xian-He Bu^{*a}

* Corresponding author. Fax: (+86) 22-23502458. E-mail: buxh@nankai.edu.cn

Table S1. Selected bond lengths [$\text{\AA}$] and angles [$^\circ$] for **1**.

Co1—O6	1.968 (4)	Co2—O2	2.053 (5)
Co1—O1	2.094 (5)	Co2—O7 ⁱⁱ	2.086 (5)
Co1—O3	2.094 (5)	Co2—O7 ⁱⁱⁱ	2.094 (5)
Co1—O4	2.162 (5)	Co2—O7	2.126 (5)
Co1—N1 ⁱ	2.178 (6)	Co2—O4	2.153 (5)
Co1—O5	2.271 (5)	Co2—O5	2.154 (5)
Co1 ^{iv} —O6—Co1	118.4 (3)	Co2—O7—Co2 ⁱⁱⁱ	90.56 (19)
Co2—O4—Co1	101.9 (2)	Co2 ⁱⁱ —O7—Co2 ⁱⁱⁱ	102.9 (2)
Co2—O5—Co1	98.4 (2)	Co2 ⁱⁱ —O7—Co2	96.1 (2)

Symmetry Code: ⁱ $y-1, -x+1, -z+3/2$; ⁱⁱ $-y+3/2, -x+3/2, -z+3/2$; ⁱⁱⁱ $y-1/2, x+1/2, -z+3/2$; ^{iv} $x, y, -z+2$

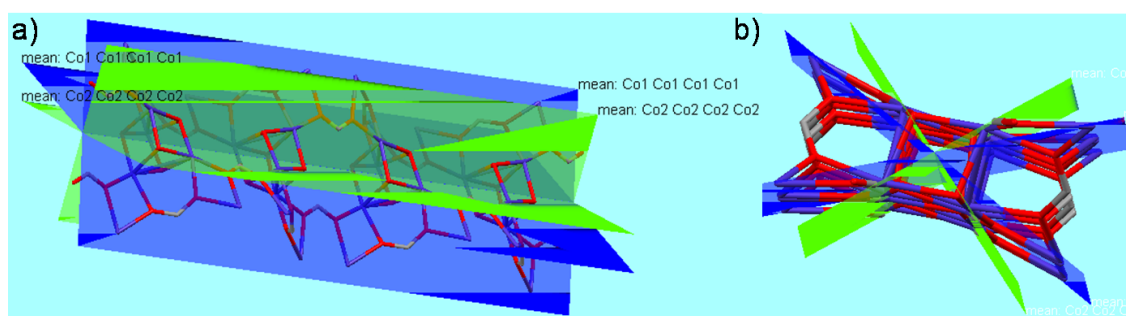


Fig. S1 a) Side and b) perspective views of four different planes defined by Co1 and Co2 ions in the 1D chain. The dihedral angle between Co1-plane and Co2-plane in the neighbor rings is 67.51° , and the dihedral angle between the same-atom planes in the neighbor rings is 49.43° for Co1 and 85.60° for Co2.

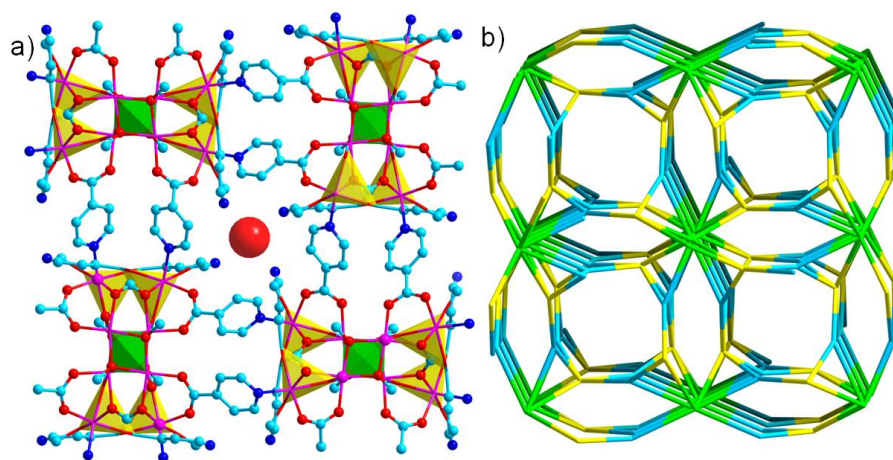


Fig. S2 (a) The 3D framework of **1** (C in cyan, O in red, N in blue, Co in pink). (b) The (3,6,8)-connected topological network of **1**. The cyan, yellow, and green spheres represent the 3-, 6-, and 8-connected nodes, respectively.

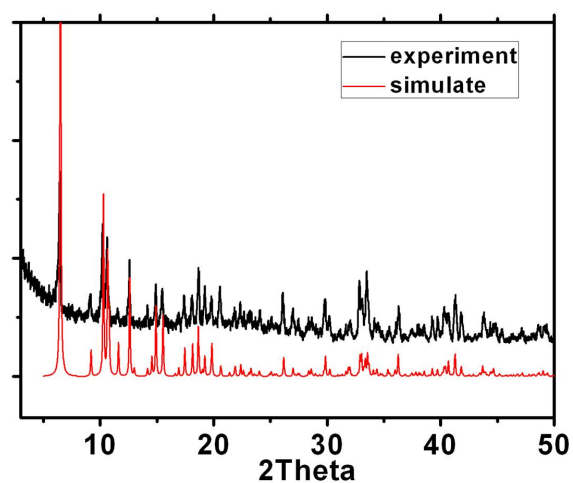


Fig. S3 X-ray powder diffraction (XRPD) patterns of **1**.

Magnetic measurements

The magnetic measurements were performed by using an MPMS XL-5 SQUID magnetometer with polycrystals of **1**. With the help of a ^3He system (at Tianjin Normal University), measurements can be performed down to 0.5 K. Diamagnetic corrections were estimated by using Pascal constants and background corrections by experimental measurement on sample holders.

Co1 and Co2 ions are all in a slightly compressed octahedral geometry with O1, O3 and O2, O7 in the apical positions. And the axis of the neighbouring Co^{II} octahedrons tilt to each other in the ring and parallel Co^{II} octahedrons could be found in in interval rings (Fig 2a). The bridges could be classified as four different type, μ -O atom, *syn,syn*, *syn.anti* and *anti,anti* carboxylate (including formate) (Fig S4 and Table S2). The interactions conducted by the μ -O atom dependent on the Co-O-Co angle. Large Co-O-Co angles lead to antiferromagnetic interactions while small ones

induce ferromagnetic coupling with critical angle about 90° . The *syn,syn* carboxylate conduct stronger antiferromagnetic interactions between metal ions, while others conduct weak ferro- or antiferromagnetic interactions. Thus, the coupling J_1 , J_2 , J_3 , J_4 and J_5 are antiferromagnetic,^{1,2,3,4,5} while the interactions of J_6 , J_6 , J_8 , and J_9 should be weak ferro- or antiferromagnetic.⁶ Taking spin-canting into account two factors need to be considered: the antisymmetric interaction (Dzyaloshinsky-Moriya interaction, DMI) and the magnetic anisotropy.⁶ The ions in the Co1 dimer and the Co2 tetrahedron are all asymmetric and the neighbouring Co^{II} octahedrons tilt to each other. Furthermore the Co^{II} ions show significant anisotropy, thus spin canting may occur in the asymmetric arranged Co^{II} chain. Also, the spin competition resulted from antiferromagnetic interactions between the Co2 tetrahedrons may contribute to the spin canting. The interactions conducted by the *syn,syn* carboxylate and μ -O atom are dominate, and the mainly canted antiferromagnetic interactions in the chains could be predigested in Fig 2b (blue line).

References:

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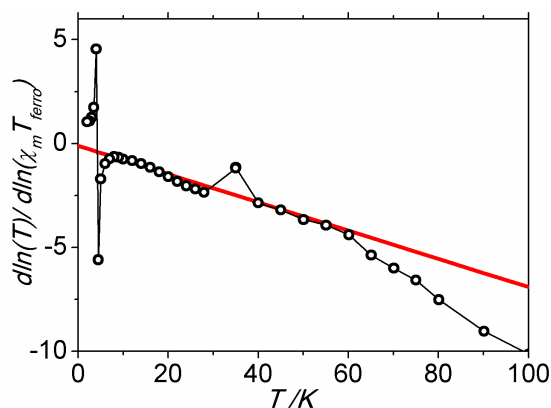


Fig. S6 Plot of the magnetic data as $d\ln(T)/d\ln(\chi_m T_{ferro})$ vs T for **1**. The intercepts of the red linear fitting lines with the horizontal and vertical zero axes give T_c and $1/\gamma$ respectively in which $\chi_m T_{ferro} = \chi_m T_{ferro} - C_2 \exp(\beta/T)$, C_2 and β attributed to the antiferromagnetic contribution of the fitting of the date by $\chi_m T = C_1 \exp(\alpha/T) + C_2 \exp(\beta/T)$. A direct fit of $d\ln(T)/d\ln(\chi_m T_{ferro})$ vs T plots the temperatures from 7 – 50 K yields $T_c = 0$ K and $g = \infty$, which is consistent with the theoretical 1D-Ising value. (ref: J. Souletie, P. Rabu and M. Drillon, *Scaling Theory Applied to Low Dimensional Magnetic System in Magnetism: Molecules to Materials*, vol. V ed. J. S. Miller and M. Drillon, Wiley-VCH Weinheim, Germany, 2005)

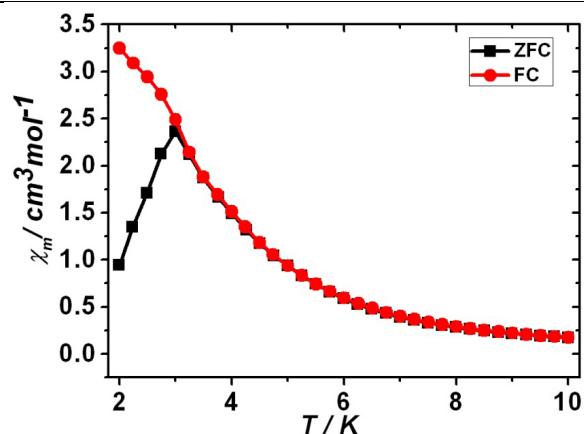


Fig. S7 Plot of ZFC and FC of **1** under 20 Oe external field.

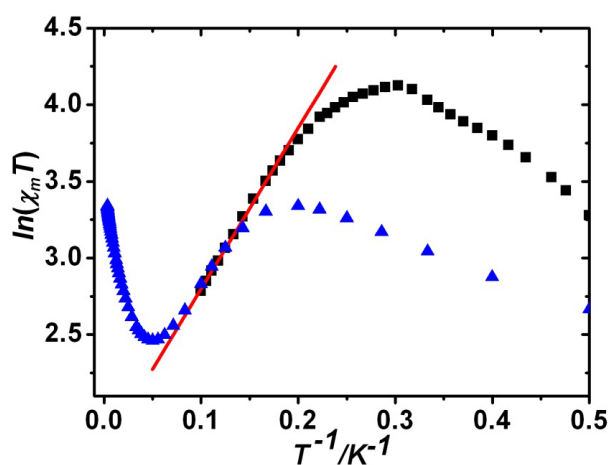


Fig. S8 Logarithm of $\chi_m T$ vs $1/T$ plots (assuming eight Co^{II} ions per formula unit) of **1**. The dc susceptibility ($\blacktriangle$) was collected at 0.2 T and the ac susceptibility ($\blacksquare$) was obtained at a frequency of 100 Hz under oscillating field of 3 Oe. The red line is the result of fitting by $\chi_m T = C_{\text{eff}} \exp(\Delta/k_B T)$ between 15 K and 5 K with the energy gap of $\Delta/k_B = 9.60$ K and $C_{\text{eff}} = 6.84 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$.

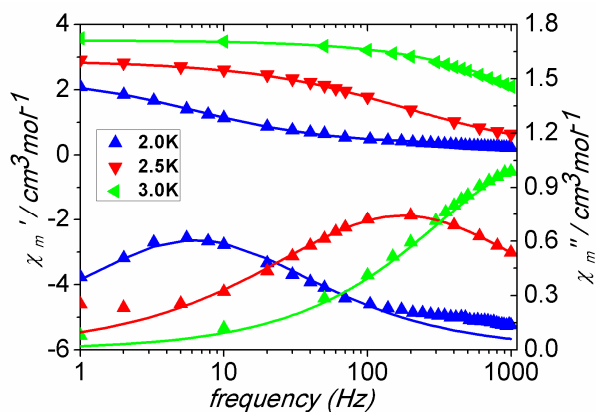


Fig. S9 Frequency dependence of ac susceptibilities at 2.0 K, 2.5 K and 3 K in zero applied dc field for **1**. The red solid line represents the least-squares fitting of the data.

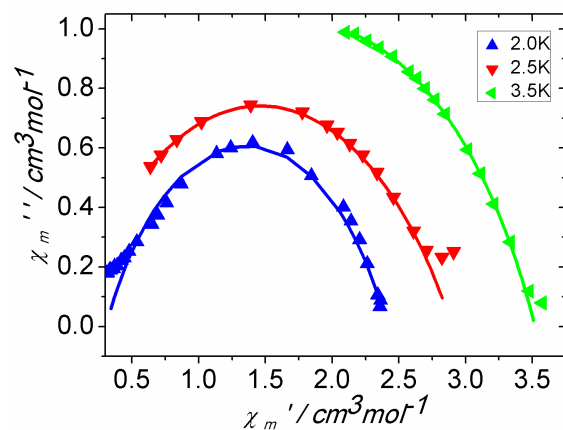


Fig. S10 Cole-Cole plot of χ_M'' vs χ_M' at 2.0 K, 2.5 K and 3 K in zero applied dc field for **1**. The red solid line represents the least-squares fitting of the data to a distribution of single relaxation processes.

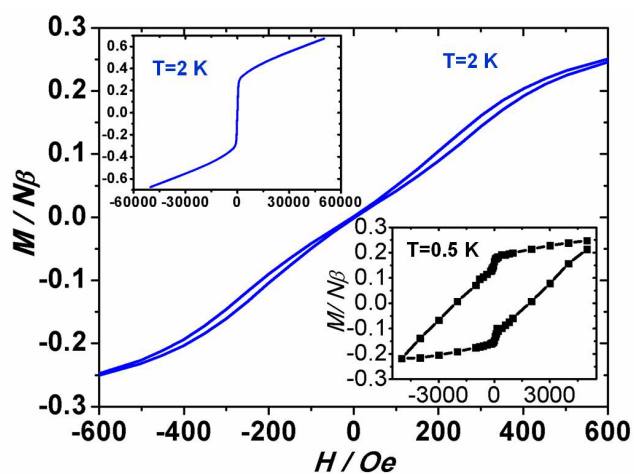


Fig. S11 Plots of hysteresis loop of **1** at 2 K and 0.5 K.