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## Photophobic and phototropic movement of a self-oscillating

gel

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## **Supporting Information**

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# 1. Experimental details for synthesizing the BZ gel

1.1 Materials.

N-Isopropyl acrylamide (NIPAAm, Tokyo Chemical Industry Co., Inc.), N,N'-methylenebisacrylamide (MBA, Sigma-Aldrich), 2-acrylamido-2methylpropanesulfonic acid (AMPS, Sigma-Aldrich), 2,2'-azobis (isobutyronitrile) (AIBN, Sigma-Aldrich), ruthenium (4-vinyl-4'-methyl-2,2`-bipylridine) bis(2,2`bipyridine) bis (hexafluorophosphate) (Ru(bpy)<sub>3</sub>, were synthesized according to established protocols<sup>1-3</sup>). Other reagents were purchased from Sinopharm Chemical Co. Ltd.

#### **1.2 Fabrication of BZ gel in capillary.**

The BZ gel was synthesized in a capillary of inner diameter of  $1.0 \pm 0.1$  mm according to previous work <sup>4</sup>. NIPAAm (0.305 g), Ru(bpy)<sub>3</sub> (37.8 mg), AIBN ( 4.6 mg) and MBA (2.2 mg) were dissolved in 1.0 mL of methanol, and AMPS (10.2 mg) was dissolved in 1.0 mL of distilled water (both solvents were previously purged with N<sub>2</sub>). The solutions were mixed and injected into the capillary. The open end of the capillary was sealed with PVC and then polymerized at 60.0 °C for 24 h. After gelation, the capillary was cut into small sections (5.0-6.0 mm in length). The resulting BZ gel was soaked in pure methanol for a week to remove unreacted monomers and then hydrated by immersing in a graded series of methanol/water mixtures for 1 day each in 100%, 75%, 50%, 25% and 0% methanol in water. Finally, we obtained the pure one dimensional poly(NIPAAm-co-Ru(bpy)<sub>3</sub>-co-AMPS) gel.



## 2. Experimental OPR curves

**Fig. S1** Dependence of oscillatory frequency of bulk system on light intensity. a: I = 1.65  $\mu$ W cm<sup>-2</sup>, frequency = 0.01524 s<sup>-1</sup>; b: I = 52  $\mu$ W cm<sup>-2</sup>, frequency = 0.01813 s<sup>-1</sup>; c: I = 280  $\mu$ W cm<sup>-2</sup>, frequency = 0.01951 s<sup>-1</sup>; d: I = 2810  $\mu$ W cm<sup>-2</sup>, frequency = 0.01413 s<sup>-1</sup>; e: I = 4410  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.01413 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.01413 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.01413 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, frequency = 0.00984 s<sup>-1</sup>; f: I = 5200  $\mu$ W cm<sup>-2</sup>, f: I = 5200

## **3. Model and Simulation**

Amemiya *et al.*<sup>5</sup> introduced a three-variable photosensitive Oregonator model to describe both photoinduction and photoinhibition in the ruthenium-catalysed BZ reaction. The equations were nondimensionalized by using the Tyson scaling<sup>6</sup>.

$$\begin{cases} \varepsilon \frac{\mathrm{dx}}{\mathrm{d\tau}} = x(1-x) + y(q-x) - \varepsilon k_f x + \phi P_2 \\ \varepsilon \frac{\mathrm{dy}}{\mathrm{d\tau}} = -y(q+x) + fz + \varepsilon k_f(y_0 - y) + \phi P_1 \\ \frac{\mathrm{dz}}{\mathrm{d\tau}} = x - z - k_f z + \phi (\frac{P_1}{2} + P_2) \end{cases}$$
(I)

In the equations,  $y_0$  is the nondimensionalized concentration of Br<sup>-</sup> in the feed flow solution;  $k_f$  is the flow rate,  $\phi$  is the light flux,  $P_1$  and  $P_2$  are the factors for photoinhibition and photoinduction, respectively. By applicating a steady state approximation for y, the above equations can be reduced to a two-variable Oregonator model (II). Because our experiments were carried out in a batch system, we set  $k_f = 0$ . The resulting equations have the following form:

$$\begin{cases} \varepsilon \frac{dx}{d\tau} = x(1-x) + (fz + \phi P_1) (q - x) / (q + x) + \phi P_2 \\ \frac{dz}{d\tau} = x - z + \phi (\frac{P_1}{2} + P_2) \end{cases}$$
(II)

In our numerical simulations, we set f = 1.02,  $\varepsilon = 0.03314$  and  $q = 1.0 \times 10^{-4}$ . Values for  $P_1$  and  $P_2$  were taken from the literature<sup>5</sup>, i.e.,  $P_1 = 0.0124$ ,  $P_2 = 0.77$ . The ODEs were numerically integrated with an explicit fourth-order Runge-Kutta method with time step  $1.0 \times 10^{-5}$ . Adjusting the parameter  $\phi$  results in changes in the oscillation frequency similar to those seen in the experiments with changing illumination intensity.

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