

Electronic supporting information (ESI) for the manuscript:

**Self-assembly, metal binding ability, and magnetic properties of
dinickel(II) and dicobalt(II) triple mesocates**

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Electronic spectra

The electronic absorption spectra of **4** and **5** in water are compared to that of the proligand Et₂H₂L in THF in Fig. S1. Complexes **4** and **5** exhibit a very strong band in the UV region located at $\lambda_{\text{max}} = 250$ ($\epsilon = 88500 \text{ M}^{-1} \text{ cm}^{-1}$) and 255 nm ($\epsilon = 92000 \text{ M}^{-1} \text{ cm}^{-1}$), respectively [Fig. S1(a) and (b)]. This intense UV band is unambiguously assigned to an intraligand (IL) $\pi\text{-}\pi^*$ transition within the aromatic benzene ring, which occurs at $\lambda_{\text{max}} = 270$ nm ($\epsilon = 29500 \text{ M}^{-1} \text{ cm}^{-1}$) for Et₂H₂L with approximately one third-fold intensity relative to that in **4** and **5** as expected from the ligand/metal molar ratio of 3:2 [Fig. S1(c)]. In addition, **4** and **5** show several characteristic d–d bands in the visible and NIR regions.

Complex **4** shows a weak NIR band located at $\lambda_{\text{max}} = 945$ nm ($\epsilon = 35 \text{ M}^{-1} \text{ cm}^{-1}$) and a visible band at $\lambda_{\text{max}} = 635$ nm ($\epsilon = 60 \text{ M}^{-1} \text{ cm}^{-1}$), which is responsible for its green color [Fig. S1(a)]. These two d–d bands are assigned to the ${}^3\text{T}_{2\text{g}}(\text{F}) \leftarrow {}^3\text{A}_{2\text{g}}(\text{F})$ (ν_1) and ${}^3\text{T}_{1\text{g}}(\text{F}) \leftarrow {}^3\text{A}_{2\text{g}}(\text{F})$ (ν_2) transitions respectively, of the octahedral high-spin d⁸ Ni^{II} ion (*O_h* point group).¹ The third spin-allowed d–d band corresponding to the ${}^3\text{T}_{1\text{g}}(\text{P}) \leftarrow {}^3\text{A}_{2\text{g}}(\text{F})$ (ν_3) transition appears as a poorly resolved shoulder at λ_{max} ca. 375 nm ($\epsilon = 280 \text{ M}^{-1} \text{ cm}^{-1}$) in the low-energy tail of the intense UV band. Instead, the extremely weak absorption feature at $\lambda_{\text{max}} = 788$ nm ($\epsilon = 25 \text{ M}^{-1} \text{ cm}^{-1}$) is often ascribed to the spin-forbidden d–d band corresponding to the ${}^1\text{E}_{\text{g}}(\text{F}) \leftarrow {}^3\text{A}_{2\text{g}}(\text{F})$ transition of octahedral high-spin d⁸ Ni^{II} ion (*O_h* point group).¹

Complex **5** shows a weak NIR band at $\lambda_{\text{max}} = 1045$ nm ($\epsilon = 15 \text{ M}^{-1} \text{ cm}^{-1}$) and a visible band at $\lambda_{\text{max}} = 510$ nm ($\epsilon = 100 \text{ M}^{-1} \text{ cm}^{-1}$), which is responsible for its pink color [Fig. S1(b)]. These two d–d bands are commonly assigned to the ${}^4\text{T}_{2\text{g}}(\text{F}) \leftarrow {}^4\text{T}_{1\text{g}}(\text{F})$ (ν_1) and ${}^4\text{T}_{1\text{g}}(\text{P}) \leftarrow {}^4\text{T}_{1\text{g}}(\text{F})$ (ν_2) transitions respectively, of the octahedral high-spin d⁷ Co^{II} ion (*O_h* point group).¹ The third spin-allowed d–d band which corresponds to the ${}^4\text{A}_{2\text{g}}(\text{F}) \leftarrow {}^4\text{T}_{1\text{g}}(\text{F})$ (ν_3) transition, being essentially a two-electron transition from $(t_{2\text{g}})^5(e_{\text{g}})^2$ to $(t_{2\text{g}})^3(e_{\text{g}})^4$ configurations, is expected to be weak and probably not observed.¹ In fact, the fine structure of the main visible band at $\lambda_{\text{max}} = 510$ nm, with two distinct shoulders at $\lambda_{\text{max}} = 473$ and 543 nm ($\epsilon = 90 \text{ M}^{-1} \text{ cm}^{-1}$), may likely arise from the term splitting of the ${}^4\text{T}_{1\text{g}}(\text{P}) \leftarrow {}^4\text{T}_{1\text{g}}(\text{F})$ (ν_2)

(1) (a) E. Koenig, *Struct. Bond.*, 1972, **9**, 175; (b) Y. S. Dou, *J. Chem. Ed.*, 1990, **67**, 134.

transition due to spin-orbit coupling and/or trigonal distortion of the octahedral geometry.

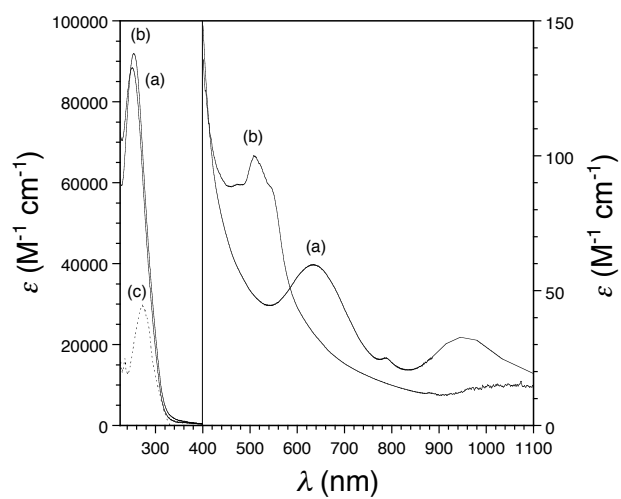


Fig. S1 UV-visible-NIR spectra of **4** (a) and **5** (b) in water and of $\text{Et}_2\text{H}_2\text{L}$ in THF (c).

Magnetic properties

The $\chi_{\text{M}}T$ vs. T plots (χ_{M} being the molar magnetic susceptibility per dinuclear unit and T the temperature) of **4** and **5** are compared in Fig. S2.

The $\chi_{\text{M}}T$ vs. T plot for **4** is characteristic of moderately weak ferromagnetically coupled Ni^{II} units with a nonnegligible zero-field splitting (ZFS) of the six-coordinate, octahedral high-spin d^8 Ni^{II} ions ($^3\text{A}_{2\text{g}}$) (Fig. S2). At room temperature, the $\chi_{\text{M}}T$ value of $2.32 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ is close to that expected for two magnetically non-interacting high-spin Ni^{II} ions [$\chi_{\text{M}}T = 2 \times (N\beta^2 g_{\text{Ni}}^2/3k)S_{\text{Ni}}(S_{\text{Ni}} + 1) = 2.31 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ with $S_{\text{Ni}} = 1$ and $g_{\text{Ni}} = 2.15$]. Upon cooling, $\chi_{\text{M}}T$ continuously increases to reach a maximum of $3.21 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ at 3.8 K, and then it slightly decreases down to $3.04 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ at 2.0 K. The increase of $\chi_{\text{M}}T$ in the high-temperature region indicates a weak ferromagnetic intradimer interaction. Yet the maximum value of $\chi_{\text{M}}T$ is slightly below than that expected for a ground quintet spin state for the Ni^{II}_2 unit [$\chi_{\text{M}}T = (N\beta^2 g^2/3k)S(S + 1) = 2N\beta^2 g^2/k = 3.47 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ with $S = 2$ and $g = g_{\text{Ni}} = 2.15$]. The slight decrease of $\chi_{\text{M}}T$ in the low-temperature region is most likely due to ZFS effects. In fact, the antiferromagnetic interdimer interactions through the diamagnetic Na^{I} ions are certainly negligible given the large intermolecular metal-metal separation of *ca.* 8 Å.

The magnetic susceptibility data of **4** were analyzed through a spin Hamiltonian for a dimer model which takes into account the axial ZFS of the $^3\text{A}_{2\text{g}}$ ground state of the octahedral high-spin Ni^{II} ions [eqn (S1) with $S_1 = S_2 = S_{\text{Ni}} = 1$], where J is the magnetic coupling parameter, D is the single-ion axial magnetic anisotropy parameter, and g is the Landé factor.² A good fit was obtained through the appropriate analytical expression^{2b} with $J = +3.6 \text{ cm}^{-1}$, $D = -3.5 \text{ cm}^{-1}$, and $g = 2.14$ (solid line in Fig. S2). The calculated values of J and D for **4** agree with those found for **2** ($J = +3.2 \text{ cm}^{-1}$ and $D = -3.4 \text{ cm}^{-1}$).³

(2) (a) A. P. Ginsberg, R. L. Martin, R. W. Brookes and R. C. Sherwood, *Inorg. Chem.*, 1972, **11**, 2884; (b) G. De Munno, M. Julve and F. Lloret, *J. Chem. Soc., Dalton Trans.*, 1993, 1179.

(3) D. Cangussu, E. Pardo, M.-C. Dul, R. Lescouëzec, P. Herson, Y. Journaux, E. F. Pedroso, C. L. M. Pereira, H. O. Stumpf, M. C. Muñoz, R. Ruiz-García, J. Cano, M. Julve and F. Lloret, *Inorg. Chim. Acta*, 2008, **361**, 3394.

$$\mathbf{H} = -J \mathbf{S}_1 \cdot \mathbf{S}_2 + D \sum_{i=1,2} \mathbf{S}_{zi}^2 + g\beta H \sum_{i=1,2} \mathbf{S}_i \quad (\text{S1})$$

The $\chi_M T$ vs. T plot for **5** agrees with a weak ferromagnetically coupled Co^{II} units with an important spin-orbit coupling (SOC) of the octahedral high-spin d^7 Co^{II} ions ($^4\text{T}_{1g}$) (Fig. S2). At room temperature, the $\chi_M T$ value of $5.98 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ is close to that expected for two magnetically non-interacting high-spin Co^{II} ions with an important orbital contribution, as compared with the spin-only value [$\chi_M T = 2 \times (N\beta^2 g_{\text{Co}}^2 / 3k) S_{\text{Co}}(S_{\text{Co}} + 1) = 3.75 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ with $S_{\text{Co}} = 3/2$ and $g_{\text{Co}} = 2.0$]. Upon cooling, $\chi_M T$ continuously decreases to reach a minimum of $3.90 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ at 12.0 K, and then it slightly increases up to $4.43 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ at 2.0 K. The increase of $\chi_M T$ in the low-temperature region is indicative of a weak ferromagnetic intradimer interaction, whereas the decrease of $\chi_M T$ in the high-temperature region is due to the SOC of the orbitally degenerate, high-spin Co^{II} ion in a distorted octahedral coordination geometry, as found in related mononuclear high-spin octahedral cobalt(II) complexes.⁴

The magnetic susceptibility data of **5** were analyzed through a spin Hamiltonian for a dimer model which includes the SOC of the $^4\text{T}_{1g}$ ground state of the octahedral high-spin Co^{II} ions, together with the splitting of the T_{1g} orbital term into a singlet (A_1) and a doublet (E) terms due to the axial distortion [eqn (S2) with $S_1 = S_2 = S_{\text{Co}} = 3/2$ and $L_1 = L_2 = L_{\text{Co}} = 1$], where J is the magnetic coupling parameter, λ is the spin-orbit coupling parameter, Δ is the axial orbital splitting, and α is an orbital reduction factor defined as $\alpha = A\kappa$.³ The κ parameter considers the reduction of the orbital momentum caused by covalency effects ($0 < \kappa \leq 1$), while the A parameter takes into account the admixture between the excited $^4\text{T}_{1g}(\text{P})$ and the ground $^4\text{T}_{1g}(\text{F})$ levels when using the T_1 and P term isomorphism (A varies between 1.0 and 1.5 for the strong and weak crystal field limits, respectively).⁴

$$\mathbf{H} = -J \mathbf{S}_1 \cdot \mathbf{S}_2 + \alpha\lambda \sum_{i=1,2} \mathbf{L}_i \cdot \mathbf{S}_i + \Delta \sum_{i=1,2} \mathbf{L}_{zi}^2 + \beta H \sum_{i=1,2} (\alpha \mathbf{L}_i + g_e \mathbf{S}_i) \quad (\text{S2})$$

A good fit was obtained by full-matrix diagonalization techniques⁵ with $J = +1.1 \text{ cm}^{-1}$, $\lambda = -132 \text{ cm}^{-1}$, $\Delta = 109 \text{ cm}^{-1}$, and $\alpha = A\kappa = 1.10$ (solid line in Fig. S2). The value of A can be calculated from the

(4) F. Lloret, M. Julve, J. Cano, R. Ruiz-García and E. Pardo, *Inorg. Chim. Acta*, 2008, **361**, 3432.

(5) J. Cano, *VPMAG package*, University of Valencia, Valencia, Spain, 2003.

crystal field parameters obtained from the analysis of the electronic spectrum of **5** through eqn (S3) and (S4) ($Dq = 10725 \text{ cm}^{-1}$ and $B = 745 \text{ cm}^{-1}$).¹ The calculated A value is 1.37 and then $\kappa = 0.79$.

$$A = (1.5 - c^2)/(1 + c^2) \quad (\text{S3})$$

$$c = 0.75 + 1.875(B/Dq) - [1 + 1.8(B/Dq) + 2.25(B/Dq)^2]^{1/2} \quad (\text{S4})$$

The calculated values of J , λ , Δ , and κ for **5** agree with those found for **3** ($J = +1.0 \text{ cm}^{-1}$, $\lambda = -116.3 \text{ cm}^{-1}$, $\Delta = 108 \text{ cm}^{-1}$, and $\kappa = 0.79$ with $A = 1.4$).⁴ The absolute values of the spin-orbit coupling parameter for these dicobalt(II) triple mesocates are indeed lower than that of the free ion ($\lambda = -180 \text{ cm}^{-1}$) because of the metal-ligand covalency. Likewise, the strongly reduced values of the orbital reduction parameter reveal a large covalency of the M–N and M–O bonds (to be compared with $\kappa = 1$ for the free ion), which is ultimately responsible for the strong delocalization of the spin density of the metal onto the phenylenediamidate bridges.

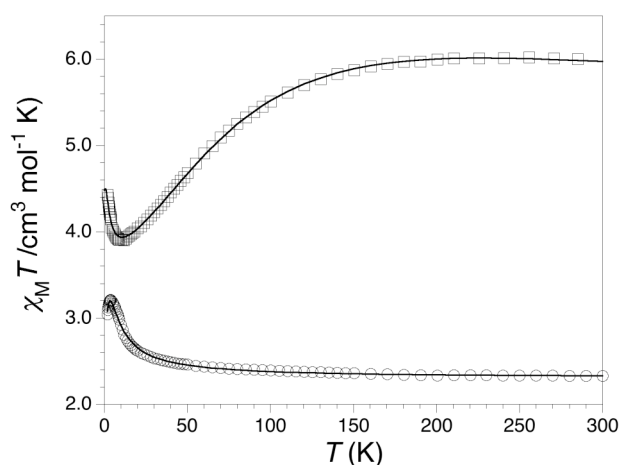
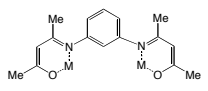
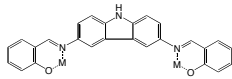
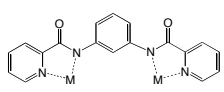
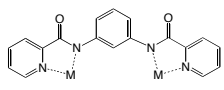
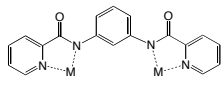
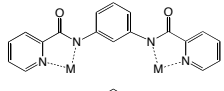
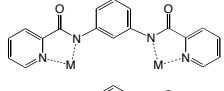
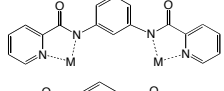
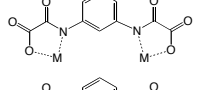
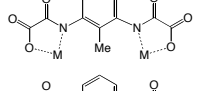
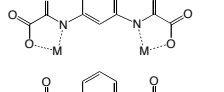
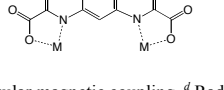


Fig. S2 Temperature dependence of $\chi_M T$ for **4** (○) and **5** (□). The solid lines are the best fit curves.

Table S1 Selected magnetostructural data for double- and triple-stranded, dinuclear helicates and mesocates of late and middle 3d metal ions with coordinating group-substituted aromatic diamines, diimines, and diamides as bridging ligands.

Entry	Complex	Type	Topology	M	Geometry	L	$d^M / \text{\AA}$	J^M / cm^{-1}	Ref.
1	$\text{Na}_4[\text{Cu}_2\text{L}_2] \cdot 18\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	O_h/D_{4h}		6.93	< 1.0	16a
2	$\text{Na}_4[\text{Cu}_2\text{L}_2] \cdot 20\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	O_h/D_{4h}		n.a.	< 1.0	16b
3	$\text{Na}_4[\text{Ni}_2\text{L}_2] \cdot 15\text{H}_2\text{O}$	mesocate	double	$\text{Ni}^{\text{II}} (d^8, \text{hs})$	O_h		6.95	-1.4	16b
4	$\text{Na}_4[\text{Co}_2\text{L}_2] \cdot 12\text{H}_2\text{O}$	mesocate	double	$\text{Co}^{\text{II}} (d^7, \text{hs})$	O_h		n.a.	-0.9	16b
5	$\text{Na}_4[\text{Mn}_2\text{L}_2] \cdot 10\text{H}_2\text{O}$	mesocate	double	$\text{Mn}^{\text{II}} (d^5, \text{hs})$	O_h		n.a.	-0.4	16b
6	$[\text{Cu}_2\text{L}_2] \cdot 2\text{CHCl}_3$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	T_d/D_{4h}		7.44	-1.0	16c
7	$[\text{Cu}_2\text{L}_2]$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	T_d/D_{4h}		n.a.	-0.9	16c
8	$[\text{Cu}_2\text{L}_2]$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	T_d/D_{4h}		n.a.	-1.1	16c
9	$[\text{Ni}_2\text{L}_2] \cdot 3\text{H}_2\text{O}$	mesocate	double	$\text{Ni}^{\text{II}} (d^8, \text{hs})$	T_d		n.a.	-2.8	16c
10	$[\text{Ni}_2\text{L}_2] \cdot 3\text{H}_2\text{O}$	mesocate	double	$\text{Ni}^{\text{II}} (d^8, \text{hs})$	T_d		n.a.	-2.7	16c
11	$[\text{Ni}_2\text{L}_2]$	mesocate	double	$\text{Ni}^{\text{II}} (d^8, \text{hs})$	T_d		n.a.	-2.9	16c
12	$[\text{Co}_2\text{L}_2] \cdot \text{CHCl}_3$	mesocate	double	$\text{Co}^{\text{II}} (d^7, \text{hs})$	T_d		7.26	-1.3	16c
13	$[\text{Co}_2\text{L}_2] \cdot 4\text{H}_2\text{O}$	mesocate	double	$\text{Co}^{\text{II}} (d^7, \text{hs})$	T_d		n.a.	-1.2	16c
14	$[\text{Co}_2\text{L}_2]$	mesocate	double	$\text{Co}^{\text{II}} (d^7, \text{hs})$	T_d		n.a.	-1.4	16c
15	$[\text{Fe}_2\text{L}_2(\text{NO}_3)_2]$	mesocate	double	$\text{Fe}^{\text{III}} (d^5, \text{hs})$	O_h		7.90	+0.4	16d
16	$[\text{Cu}_2\text{L}_2] \cdot 2\text{MeCN}$	helicate	double	$\text{Cu}^{\text{II}} (d^9)$	T_d		12.02	-0.025 ^c	17
17	$\text{Li}_4[\text{Cu}_2\text{L}_2] \cdot 10\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	D_{4h}		n.a.	-95	18a
18	$\text{Na}_4[\text{Cu}_2\text{L}_2] \cdot 11\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	D_{4h}		7.91	-81	18a
19	$(\text{Ph}_4\text{P})_4[\text{Cu}_2\text{L}_2] \cdot 8\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	D_{4h}		n.a.	-94	18a
20	$\text{Li}_4[\text{Cu}_2\text{L}_2] \cdot 12\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (d^9)$	D_{4h}		n.a.	-11.5	18a

21	$\text{Na}_4[\text{Cu}_2\text{L}_2] \cdot 12\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		12.19	-9.5	18a
22	$(\text{Bu}_4\text{N})_4[\text{Cu}_2\text{L}_2] \cdot 10\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		n.a.	-8.7	18a
23	$\text{Li}_4[\text{Cu}_2\text{L}_2] \cdot 12\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		n.a.	-20.5	18b
24	$(\text{Ph}_4\text{P})_4[\text{Cu}_2\text{L}_2] \cdot 8\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		8.33	-20.7	18b
25	$\text{Li}_4[\text{Cu}_2\text{L}_2] \cdot 12\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		n.a.	-21.2	18b
26	$(\text{Ph}_4\text{P})_4[\text{Cu}_2\text{L}_2] \cdot 4\text{H}_2\text{O}$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		n.a.	-23.0	18b
27	$(\text{Bu}_4\text{N})_4[\text{Cu}_2\text{L}_2]$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		12.48	-23.9	18c
28	$[\text{Co}_2\text{L}_3]$	mesocate	triple	$\text{Co}^{\text{III}} (\text{d}^6, \text{ls})$	O_h		6.73	-21.1 ^d	19a
29	$[\text{Fe}_2\text{L}_3]$	mesocate	triple	$\text{Fe}^{\text{III}} (\text{d}^5, \text{hs})$	O_h		6.90	-11.0 ^d	19a
30	$[\text{Mn}_2\text{L}_3]$	mesocate	triple	$\text{Mn}^{\text{IV}} (\text{d}^3)$	O_h		6.75	n.a. ^d	19b
31	$[\text{Cu}_2\text{L}_2]$	mesocate	double	$\text{Cu}^{\text{II}} (\text{d}^9)$	D_{4h}		6.70	n.a. ^e	20a
32	$[\text{Ni}_2\text{L}_2]$	mesocate	double	$\text{Ni}^{\text{II}} (\text{d}^8, \text{ls})$	D_{4h}		n.a.		20a
33	$[\text{Co}_2\text{L}_3]$	mesocate	triple	$\text{Co}^{\text{III}} (\text{d}^6, \text{ls})$	O_h		6.72	+26 ^d	20a
34	$[\text{Fe}_2\text{L}_3]$	mesocate	triple	$\text{Fe}^{\text{III}} (\text{d}^5, \text{hs})$	O_h		6.92	n.a. ^d	20a
35	$[\text{Cr}_2\text{L}_3]$	mesocate	triple	$\text{Cr}^{\text{III}} (\text{d}^3)$	O_h		n.a.	n.a. ^e	20a
36	$[\text{Mn}_2\text{L}_3]$	mesocate	triple	$\text{Mn}^{\text{IV}} (\text{d}^3)$	O_h		6.76	-6.5 ^d	20b
37	$[\text{Fe}_2\text{L}_3](\text{PF}_6)_4$	helicate	triple	$\text{Fe}^{\text{II}} (\text{d}^6, \text{hs})$	O_h		11.40	n.a. ^f	21
38	$[\text{Fe}_2\text{L}_3](\text{BF}_4)_4$	helicate	triple	$\text{Fe}^{\text{II}} (\text{d}^6, \text{hs})$	O_h		11.56	n.a. ^f	21
39	$[\text{Fe}_2\text{L}_3](\text{ClO}_4)_4$	helicate	triple	$\text{Fe}^{\text{II}} (\text{d}^6, \text{hs})$	O_h		11.58	n.a. ^f	21
40	$[\text{Ni}_2\text{L}_3](\text{PF}_6)_4$	helicate	triple	$\text{Ni}^{\text{II}} (\text{d}^8, \text{hs})$	O_h		11.54	<1.0	21
41	$[\text{Co}_2\text{L}_3](\text{PF}_6)_4$	helicate	triple	$\text{Co}^{\text{II}} (\text{d}^7, \text{hs})$	O_h		11.54	<1.0	21
42	$[\text{Mn}_2\text{L}_3](\text{PF}_6)_4$	helicate	triple	$\text{Mn}^{\text{II}} (\text{d}^5, \text{hs})$	O_h		11.66	<1.0	21

43	[Cu ₂ L ₂]	mesocate	double	Cu ^{II} (d ⁹)	D _{4h} /T _d		7.27	+14.6	22a
44	[Cu ₂ L ₂] · 2CHCl ₃ · 2H ₂ O	mesocate	double	Cu ^{II} (d ⁹)	D _{4h} /T _d		10.78	-2.2	22b
45	[Cu ₂ L ₂] · 2H ₂ O	mesocate	double	Cu ^{II} (d ⁹)	D _{4h} /T _d		7.56	+21.1	23a
46	[Ni ₂ (HL) ₃]PF ₆ · 2H ₂ O	mesocate	triple	Ni ^{II} (d ⁸ , hs)	O _h		6.93	+3.6	23a
47	[Co(H ₂ O) ₆][Ni ₂ L ₃] · THF · 10H ₂ O	mesocate	triple	Ni ^{II} (d ⁸ , hs)	O _h		6.92	+3.6	23a
48	[Ag ₂ (H ₂ O)][Ni ₂ L ₃] · 11H ₂ O	mesocate	triple	Ni ^{II} (d ⁸ , hs)	O _h		6.96	+2.9	23a
49	[Ni ₂ (HL) ₃]ClO ₄ · 15H ₂ O	mesocate	triple	Ni ^{II} (d ⁸ , hs)	O _h		6.93	+3.1	23b
50	[Co ₂ L ₃] · 19H ₂ O	mesocate	triple	Co ^{III} (d ⁶ , ls)	O _h		6.85		23b
51	Na ₄ [Cu ₂ L ₂] · 10H ₂ O	mesocate	double	Cu ^{II} (d ⁹)	D _{4h}		6.82	+16.0	25
52	(Bu ₄ N) ₄ [Cu ₂ L ₂] · 4H ₂ O	mesocate	double	Cu ^{II} (d ⁹)	D _{4h} /T _d		7.20	+16.4	13i
53	Li ₄ [Ni ₂ L ₃] · 28H ₂ O	mesocate	triple	Ni ^{II} (d ⁸ , hs)	O _h		6.86	+3.2	26b
54	Li ₄ [Co ₂ L ₃] · 37H ₂ O	mesocate	triple	Co ^{II} (d ⁷ , hs)	O _h		6.85	+1.0	26a,b

^a Intermetallic distance. ^b Magnetic coupling parameter ($H = -JS_1 \cdot S_2$ with $S_1 = S_2 = S_M$). ^c Intermolecular magnetic coupling. ^d Radical-radical coupling.

^e Metal-radical coupling. ^f Spin crossover magnetic behavior.