

## Henry constant and Isothermic Heat at Zero-Loading for Gas Adsorption in Carbon Nanotubes

by

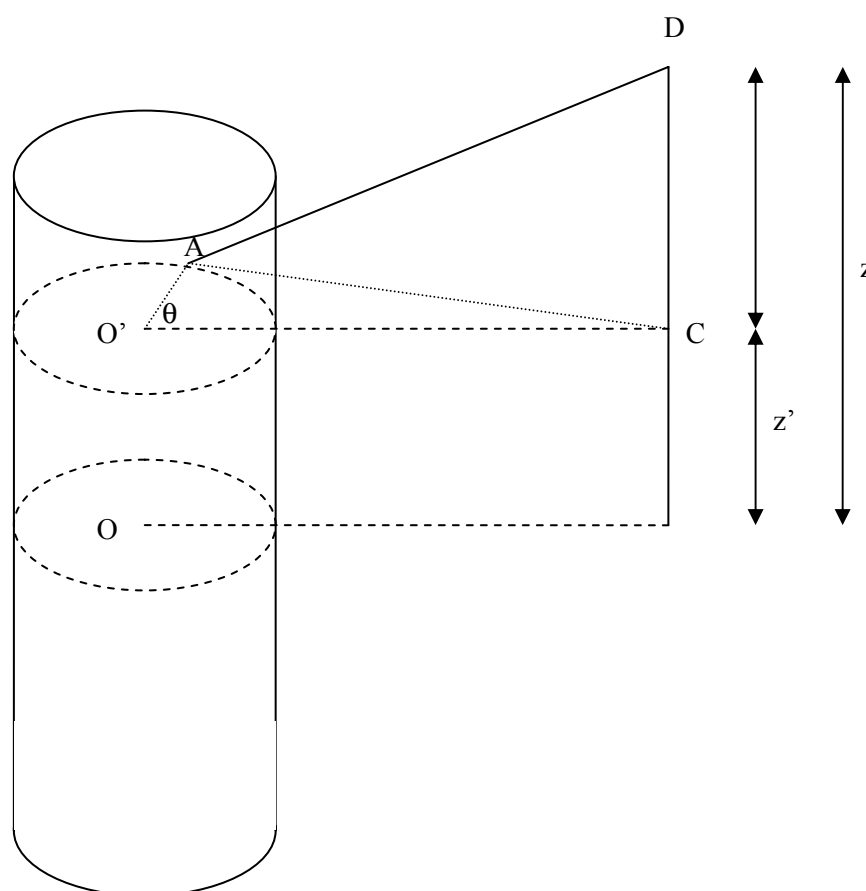
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### Supplementary Material

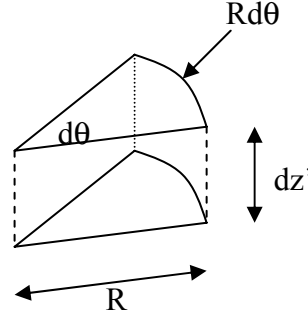
#### Solid-Fluid potential for a finite cylindrical pore



We choose the origin at the centre of the finite cylinder whose radius and length are R and L, respectively. Let D be the position of a point in space and A be an arbitrary point on the cylinder. The plane orthogonal to the cylinder defines the various points and an angle as shown in Figure A1.1. The coordinates of the point D are r and z. The distance AC is calculated from  $d_{AC}^2 = r^2 + R^2 - 2rR \cos \theta$ , and therefore the distance between the adsorbate and an arbitrary point D on the carbon nanotube is

$$d_{AD}^2 = r^2 + R^2 - 2rR \cos \theta + (z - z')^2 \quad (\text{A1.1})$$

Consider a differential area around the point A as shown with a length  $dz'$  and an angle  $d\theta$ . The area is  $dA = R d\theta dz'$ . Let  $\rho_s$  be the surface carbon atom density, the number of carbon atoms in this differential area is  $dN_C = \rho_s R d\theta dz'$ .



Therefore the potential energy of an adsorbate with all carbon atoms in the differential area  $dA$  is (assuming pairwise additivity):

$$d\varphi = \rho_s R d\theta dz' \left\{ 4\epsilon^{(\alpha,\beta)} \left[ \left( \frac{\sigma^{(\alpha,\beta)}}{d_{AD}} \right)^{12} - \left( \frac{\sigma^{(\alpha,\beta)}}{d_{AD}} \right)^6 \right] \right\} \quad (\text{A1.2})$$

for  $-L/2 \leq z' \leq L/2$  and  $0 \leq \theta \leq 2\pi$ . Substituting the distance  $d_{AD}$  of eq. (A1.1) into eq. (A1.2) and then integrating with respect to  $z'$  and  $\theta$ , we obtain the potential energy between a single particle and the finite carbon nanotube.

$$\varphi = 4\epsilon^{(\alpha,\beta)} \rho_s R \int_0^{2\pi} \int_{-L/2}^{L/2} \sum_{n=3,6} a_n \frac{[\sigma^{(\alpha,\beta)}]^n}{[(z-z')^2 + R^2 + r^2 - 2rR \cos \theta]^n} d\theta dz' \quad (\text{A1.3})$$

where  $a_3 = -1$  and  $a_6 = 1$ . Interchanging the integration and summation operation, we get:

$$\varphi = 4\epsilon^{(\alpha,\beta)} \rho_s R \sum_{n=3,6} a_n [\sigma^{(\alpha,\beta)}]^{2n} \int_0^{2\pi} I(\theta; r, z; n, R, L) d\theta \quad (\text{A1.4})$$

Here the integrand  $I(\theta; r, z; n, R, L)$  is a function of  $\theta$ , and the above integral can be integrated numerically. This integrand  $I$  is:

$$I = \int_{-L/2}^{L/2} \frac{1}{[(z-z')^2 + R^2 + r^2 - 2rR \cos \theta]^n} dz' \quad (\text{A1.5})$$

Depending on the value of  $r^2 + R^2 - 2rR \cos \theta$ , the integrand  $I$  can be evaluated analytically as below. For  $r^2 + R^2 - 2rR \cos \theta = 0$ , the function  $I$  takes the form:

$$I = (1 - 2n)^{-1} \left[ \left( \frac{L}{2} - z \right)^{1-2n} + \left( \frac{L}{2} + z \right)^{1-2n} \right] \quad (\text{A1.6a})$$

When  $r^2 + R^2 - 2rR \cos \theta \neq 0$ , the function I is:

$$I = \frac{H(x_2; n) - H(x_1; n)}{(R^2 + r^2 - 2rR \cos \theta)^{n-(1/2)}} \quad (\text{A1.6b})$$

where  $x_1$  and  $x_2$  are given by:

$$x_1 = \frac{-\left(\frac{L}{2} + z\right)}{\sqrt{R^2 + r^2 - 2rR \cos \theta}}; \quad x_2 = \frac{\left(\frac{L}{2} - z\right)}{\sqrt{R^2 + r^2 - 2rR \cos \theta}} \quad (\text{A1.6c})$$

The function  $H(x; n)$  is given by:

$$\begin{aligned} H(x; n) &= \int (1 + x^2)^{-n} dx \\ &= \frac{x}{(2n-1)} \sum_{j=1}^{n-1} \frac{(2n-1)(2n-3)\cdots(2n-2j+1)}{2^j (n-1)(n-2)\cdots(n-j)(1+x^2)^{n-j}} + \frac{(2n-3)!!}{2^{n-1}(n-1)!} \arctan(x) \end{aligned} \quad (\text{A1.6d})$$

When the pore becomes infinite, the potential can be expressed in terms of the hypergeometric function [1]. It has the form:

$$\varphi = 4\pi^2 \rho_s [\sigma^{(\alpha, \beta)}]^2 \varepsilon^{(\alpha, \beta)} (I_6 - I_3) \quad (\text{A1.7a})$$

where  $I_6$  and  $I_3$  are given by

$$I_6 = \frac{63}{128} \left( \frac{\sigma^{(\alpha, \beta)}}{R} \right)^{10} \left[ 1 - \left( \frac{r}{R} \right)^2 \right]^{-10} \times F \left[ -\frac{9}{2}; -\frac{9}{2}; 1; \left( \frac{r}{R} \right)^2 \right] \quad (\text{A1.7b})$$

$$I_3 = \frac{3}{4} \left( \frac{\sigma^{(\alpha, \beta)}}{R} \right)^4 \left[ 1 - \left( \frac{r}{R} \right)^2 \right]^{-4} \times F \left[ -\frac{3}{2}; -\frac{3}{2}; 1; \left( \frac{r}{R} \right)^2 \right] \quad (\text{A1.7c})$$

Here  $R$  is the pore radius and  $r$  is the distance from the center to the adsorbate center.

The potential energy equation of (A1.7) has been used by a number of workers [2-6] in their analysis of solids having cylindrical pores, such as carbon nanotube and MCM-41.

1. G. J. Tjatjopoulos, D. L. Feke, and J. A. Mann, Jr., *J. Phys. Chem.*, 1998, **92**, 4006-7.
2. Ravikovitch, P.I.; Haller, G.; Neimark, A. *Adv. Coll. Int. Sci.* 76 (1998) 203.
3. Neimark, A.; Ravikovitch, P.I.; Grun, M.; Schuth, F.; Unger, K. J. *Coll. Int. Sci.* 207 (1998) 159.
4. Gordon, P.A.; Saeger, R.B. *Ind. Eng. Chem. Res.* 38 (1999) 4647.
5. Neimark, A.; Vishnyakov, A. *Phys. Rev. E* 62 (2000) 4611.
6. Ohba, T.; Kaneko, K. *J. Phys. Chem.* 106 (2002) 7171.