## **Supplementary Information**

## Lithium-Ion Diffusion Mechanisms in the Battery Anode Material Li<sub>1+x</sub>V<sub>1-x</sub>O<sub>2</sub>

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Table S1. Interatomic Buckingham potentials used in this study.

Interaction	A (eV)	ρ (Å)	C (eVÅ <sup>6</sup> )	γ (e)	k (eVÅ <sup>-2</sup> )
Li <sup>+</sup> 0 <sup>2-</sup>	632.1018	0.2906	0.0	1.00	99999
$V^{4+} \dots O^{2-}$	1380.82	0.30170	0.0	3.04	196.30
$V^{3+} \dots O^{2-}$	1380.82	0.30170	0.0	2.04	196.30
O <sup>2-</sup> O <sup>2-</sup>	22764.30	0.1490	20.0	-2.96	30.0

The interatomic potentials are based on the Born model for polar solids where the interactions between ions are represented in terms of a long-range Coulombic term plus an analytical function representing short- range repulsive and van der Waals interactions. For this study, the short-range interactions were modelled using the Buckingham potential:

$$V_{ij}(r) = A_{ij} \exp\left(-\frac{r}{\rho_{ij}}\right) - \frac{C_{ij}}{r^6}$$

where r is the interatomic distance and A,  $\rho$  and C are empirically derived parameters. Charged defects will polarise nearby ions in the lattice and therefore, to calculate the defect energies accurately we included the electronic polarisability in the model, which is incorporated via the shell model.

Property	Structural composition		
Elastic Constant	LiVO <sub>2</sub>	Li <sub>1.07</sub> V <sub>0.93</sub> O <sub>2</sub>	
$(10^{11} \text{ dyne cm}^{-2})$			
C 11	47.44	47.91	
<i>c</i> <sub>12</sub>	17.79	19.14	
C 13	6.63	6.82	
C 33	28.75	30.34	
C 44	6.26	6.57	
C 66	14.28	14.39	
Dielectric			
Constant			
$< \varepsilon_{\text{static}} >$	5.82	6.09	
<pre> chigh freq &gt;</pre>	4.39	4.45	

Table S2. Calculated elastic and dielectric properties of LiVO2 and  $Li_{1.07}V_{0.93}O_2$ .

## **Defect Reactions:**

Li Frenkel defect:

$$Li_{Li}^{x} \rightarrow Vac_{Li}^{'} + Li_{i}^{\bullet}$$

Antisite defect:

$$Li_{Li}^{x} + V_{V}^{x} \rightarrow Li_{V}^{"} + V_{Li}^{\bullet \bullet}$$

Using these defect reaction equations, the energies of the isolated point defects (vacancy, interstitial, substitutional) were combined to derive the total Frenkel and anti-site defect formation energies.

The point defects were modelled using the Mott-Littleton approach, in which a defect is introduced into the energy minimised lattice, and the surrounding ions partitioned into two regions. An inner sphere of ions immediately surrounding the point defect (region 1) is then relaxed explicitly whilst the crystal bulk (region 2) is treated by computationally less expensive quasi-continuum methods.