

## **Supporting information**

# **Unravelling the Conformations of di-(perylene bisimide acrylate) by combining Time-Resolved Fluorescence-Anisotropy experiments and Molecular Modelling**

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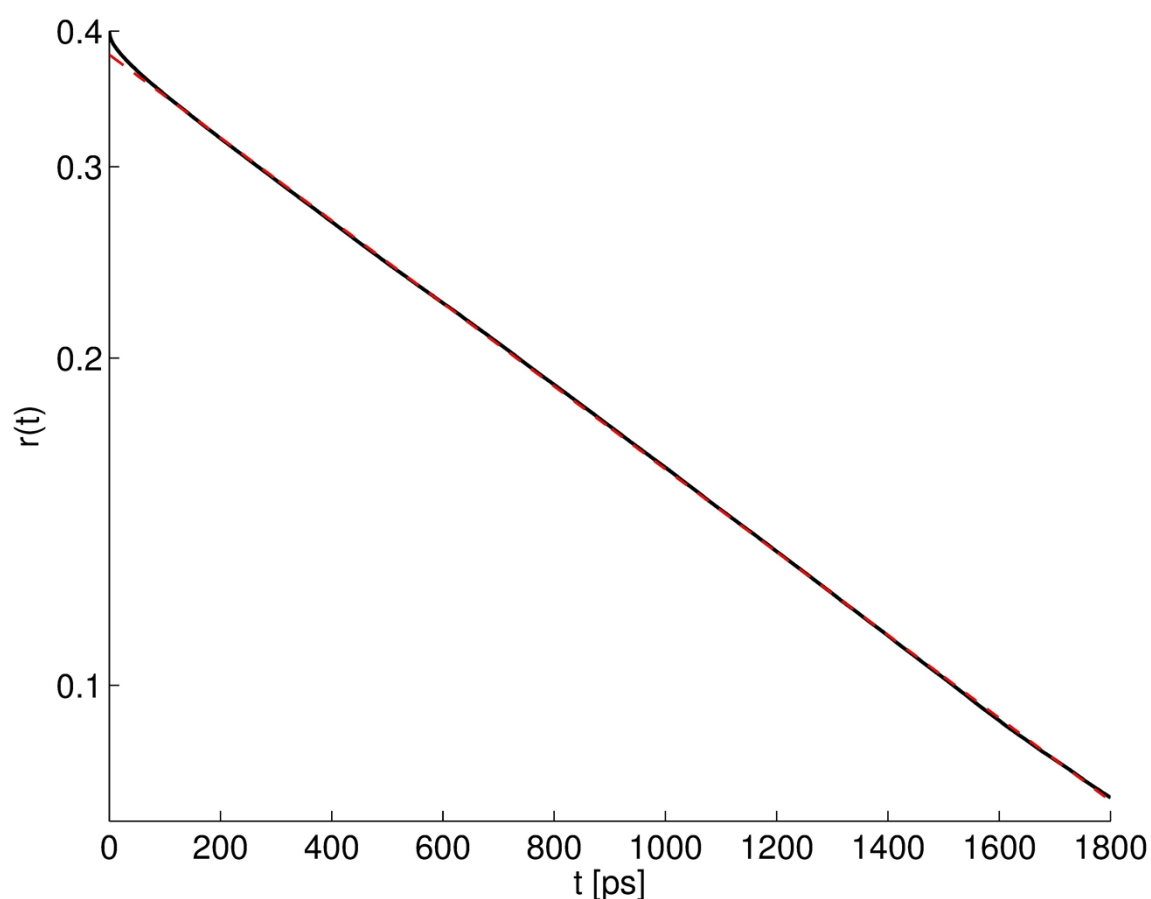
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## I. Time rescaling in MD simulations

Since the focus of our work is comparing the photophysical properties of monomeric PerAcr with those of dimeric  $(\text{PerAcr})_2$  we choose the long-time anisotropy decay of the former as a reference time scale. From simulations on a single monomer we obtain the anisotropy decay depicted in figure SI\_1. It features a clear mono-exponential tail with a time constant  $\theta_{\text{sim}}=1140\text{ps}$ . Comparing with the experimental value  $\theta_{\text{exp}}=460\text{ps}$  yields a scaling factor of 2.5. All times obtained from MD simulations were divided by this constant factor without any further adjustments.



*Fig.SI\_1: Anisotropy decay of a single PerAcr monomer (solid line) with a fit to the mono-exponential tail (red dashed line).*

## II. Förster-distance for PBI-PBI homo-ET in (PerAcr)<sub>2</sub>

Förster resonance energy transfer (FRET) refers to the transfer of electronic excitation energy between a donor (D) and an acceptor (A) molecule, that are only weakly coupled. This allows to take advantage of time-dependent perturbation theory and to derive under several further approximations (restriction of the electrostatic interaction to the leading dipole-dipole term) an expression from Fermi's golden rule for the energy transfer rate from D to A. This yields that the rate depends on the D-A distance as  $R^{-6}$ , the reciprocal lifetime of the excited state of the unperturbed donor, and further scaling factors that can be lumped together in a constant termed the Förster radius  $R_0$ , which is given as:<sup>1</sup>

$$R_0^6 = \frac{9(\ln 10)\kappa^2 Q_D}{128\pi^5 N n^4} \int_0^\infty F_D(\lambda) \epsilon_A(\lambda) \lambda^4 d\lambda \quad (1)$$

Here  $N$  denotes Avogadro's constant,  $n$  refers to the refractive index of the solvent;  $F_D$  is the (normalised) fluorescence spectrum of the donor,  $Q_D$  is the fluorescence quantum yield of the unperturbed donor,  $\epsilon_A$  denotes the extinction coefficient of the acceptor, and  $\kappa^2$  is the orientational factor of the two interacting transition-dipole moments. Under the assumption of dynamic random averaging of the donor/acceptor orientation this yields  $2/3$ . The Förster radius corresponds to that distance between the donor and the acceptor where the probability is 50% for energy transfer from D to A during the lifetime of the donor excited state. Using the fluorescence and absorption spectra reported for PerAcr earlier,<sup>2</sup> we find a value of  $R_0 = 5$  nm for the Förster radius between the two PBIs in (PerAcr)<sub>2</sub>.

### III. Derivation of Eq. 9 of the main text

Let  $p_a(t')$  denote the probability that the excitation is located on monomer  $a$  at time  $t'$ . With the reaction rate  $k$ , its rate of change is given by

$$\begin{aligned}\dot{p}_a(t') &= -p_a(t')k + (1-p_a(t'))k \\ &= -2p_a(t')k + k\end{aligned}\tag{2}$$

which can be solved by the ansatz

$$p_a(t') = \frac{1}{2} + C_0 e^{-2kt'}\tag{3}$$

where  $C_0$  is a constant. Assuming that  $p_a(t'=0) = 1$  fixes  $C_0 = \frac{1}{2}$ .

Thus, the probability that the excitation is located on monomer  $a$  at time  $\Delta t$ , given that it was on  $a$  at  $t'=0$ , is

$$p_a(\Delta t) = \frac{1}{2} + \frac{1}{2} e^{-2k\Delta t}.\tag{4}$$

This is the resting probability used in the main text.

#### IV. Calculation of the hopping constant $k_0$ and the orientation factor $\kappa$ in Eq. 10 of the main text

The reaction rate  $k(t)$  is given by <sup>3</sup>

$$k(t) = \frac{3}{2} \frac{\kappa(t)^2}{\tau} \left( \frac{R_0}{R} \right)^6. \quad (5)$$

The fluorescence life time  $\tau$  for the ISO state has been determined to be 3.7ns <sup>2</sup> and the Förster radius  $R_0=5$ nm which leads to  $k_0=6.33$  nm<sup>6</sup>/ps.

The orientation factor  $\kappa(t)$  is calculated from <sup>3</sup>

$$\kappa(t) = \left( \cos \alpha(t) - 3 \cos \theta_D(t) \cos \theta_A(t) \right)^2 \quad (6)$$

where  $\alpha$  is the angle between the two dipole vectors as in Fig. 3 of the main text and  $\theta_D$  ( $\theta_A$ ) is the angle between the acceptor (donor) dipole with the vector connecting the centres-of-mass of the two monomers.

## References

- 1 S. E. Braslavsky et al., *Photochem. Photobiol. Sci.* 7, 1444-1448 (2008).
- 2 F. Spreitler, M. Sommer, M. Thelakkat, J. Köhler, *Phys. Chem. Chem. Phys.* 14, 7971-7980 (2012).
- 3 J. R. Lakowicz, *Principles of Fluorescence Spectroscopy* (Springer, ed. 3, 2006).