# Dissociative dynamics of $\mathrm{O}_{2}$ on $\mathrm{Ag}(110)$ 

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## Electronic Supplementary Information

## 1. THREE DIMENSIONAL ATOMIC POTENTIAL ENERGY SURFACE

The interaction of the O atom with the frozen $\mathrm{Ag}(110)$ surface is described with a three-dimensional (3D) adiabatic potential energy surface (PES) that depends on the position $\mathbf{R}$ of the O atom over the surface. This 3D PES is simply used as an auxiliary utility to reduce the corrugation of the $6 \mathrm{D} \mathrm{O}_{2} / \mathrm{Ag}(110) \mathrm{PES}$. In this respect, it is important to keep the $\operatorname{Ag}(110)$ frozen, as done in the 6D PES.

The 3D PES $V^{3 D}(\mathbf{R})$ is constructed by interpolating with the 3 D corrugation reducing procedure (CRP) [1] the spin-polarized DFT energy grid described in Sec. II of the manuscript. In the 3 D CRP, the function $I^{3 D}$ to be interpolated is calculated as

$$
\begin{equation*}
I^{3 D}(\mathbf{R})=V^{3 D}(\mathbf{R})-\sum_{i=1}^{n} V^{1 D}\left(\left|\mathbf{R}-\mathbf{R}_{i}\right|\right) . \tag{1}
\end{equation*}
$$

where $V^{1 D}$ is a one-dimensional potential energy describing the interaction between the O atom and the $i$ th -Ag atom that is located at $\mathbf{R}_{\mathbf{i}}$. In our case, we use the potential energy of the O atom on the top site of $\operatorname{Ag}(110)$ for $V^{1 D}$. The summation in Eq. (1) runs over first and second layer atoms that give a non zero contribution to $V^{1 D}$. Compared to $V^{3 D}$, the resulting $I^{3 D}$ is a less corrugated function that can be easily interpolated over $X, Y$ and $Z$ through a third order spline interpolation.

The dependence of the $\mathrm{O} / \operatorname{Ag}(110)$ PES on the distance $Z$ between the atom and the surface is shown in Fig. 1 for the top site, which is used to evaluate $V^{1 D}$, and for other high-symmetry sites. In all cases the interaction of O with $\mathrm{Ag}(110)$ along the $Z$ direction is characterized by a well followed by a repulsive wall. The systematic inspection of the $(X, Y)$-cuts of the 3D PES at fixed distances $Z$ allows us to determine which sites are truly minima in 3D. From this analysis we confirm that there is a minimum at the long-bridge site with the O atom $0.56 \AA$ above the surface and a well depth of 3.09 eV . However, the absolute minimum corresponds to the atom located at $(X=1.41 \AA, Y=1.47$ $\AA, Z=0.78 \AA$ ) with an adsorption energy of 3.13 eV . It is worthy to remark that surface relaxation may alter these results for O adsorption. Finally, we note that in common with the $\mathrm{O}_{2} / \mathrm{Ag}(110) \mathrm{PES}$, none of the calculated sites show energy barriers in the entrance channel.

## 2. INTERPOLATION OF THE 6D $\mathrm{O}_{2} / \mathrm{Ag}(110)$ PES

In this part of the Supplementary Information we describe in more detail the interpolation steps followed to obtain the value of the 6D PES at any point of the configurational space, i.e., $V^{6 D}(X, Y, Z, r, \theta, \varphi)$.

Once the ( $r_{i}, Z_{i}$ )-cuts of the 34 configurations ( $X_{0}, Y_{0}, \theta_{0}, \varphi_{0}$ ) that conform the DFT energy grid described in Sec. II of the manuscript are calculated, we start obtaining the corresponding ( $r_{i}, Z_{i}$ )-cuts of the interpolation function $I^{6 D}$ for each of these configurations. This consists in subtracting the 3 D potential $V^{3 D}$ as indicated in Eq. (1) of the manuscript. Next, a conventional 2D cubic spline algorithm (as in Ref. [1]) is used to interpolate the $\left(r_{i}, Z_{i}\right)$-cuts of each $I^{6 D}\left(X_{0}, Y_{0}, \theta_{0}, \varphi_{0}\right)$ for the requested values $(r, Z)$. The latter are then interpolated over $\varphi$ and $\theta$ with Fourier expansions that take into account the symmetry of every site $\left(X_{0}, Y_{0}\right)$. In particular, the Fourier functions used are as follows:


FIG. 1. Dependence of 3D atomic PES on distance $Z$ of the oxygen atom from the $\mathrm{Ag}(110)$ for most symmetric sites on the surface and for the site where the absolute minima of the 3D PES is located (see text).
a) The configurations calculated on top site $(X=0, Y=0)$ and hollow site $(X=a / 2, Y=a /(2 \sqrt{2}))$ (both sites have the same symmetry) are: (1) $\theta=0^{\circ}$; (2) $\theta=45^{\circ}$ and $\varphi=0^{\circ}$; (3) $\theta=45^{\circ}$ and $\varphi=35.26^{\circ}$; (4) $\theta=45^{\circ}$ and $\varphi=90^{\circ}$; (5) $\theta=90^{\circ}$ and $\varphi=0^{\circ}$; (6) $\theta=90^{\circ}$ and $\varphi=35.26^{\circ}$; (7) $\theta=90^{\circ}$ and $\varphi=90^{\circ}$. In the following the notation $I(j)$ refers to the value of the interpolation function $I^{6 D}$ at the orientation conditions of item (j) above. Then, the Fourier functions $I_{\theta}(\varphi)$ for $\varphi$ interpolation on top and hollow sites and $\theta=0^{\circ}, 45^{\circ}, 90^{\circ}$ are:

$$
\begin{equation*}
I_{0}(\varphi)=I(1) \tag{2}
\end{equation*}
$$

$$
\begin{align*}
& I_{45}(\varphi)=\frac{1}{16}(2 I(2)+9 I(3)+5 I(4)+8(I(2)-I(4)) \cos 2 \varphi+(6 I(2)-9 I(3)+3 I(4)) \cos 4 \varphi)  \tag{3}\\
& I_{90}(\varphi)=\frac{1}{16}(2 I(5)+9 I(6)+5 I(7)+8(I(5)-I(7)) \cos 2 \varphi+(6 I(5)-9 I(6)+3 I(7)) \cos 4 \varphi) \tag{4}
\end{align*}
$$

The Fourier function for $\theta$ interpolation on top (T) and hollow (H) sites is:

$$
\begin{equation*}
I_{T, H}(\theta, \varphi)=\frac{1}{4}\left(\left(I_{0}(\varphi)+2 I_{45}(\varphi)+I_{90}(\varphi)+2\left(I_{0}(\varphi)-I_{90}(\varphi)\right) \cos 2 \theta+\left(I_{0}(\varphi)-2 I_{45}(\varphi)+I_{90}(\varphi)\right) \cos 4 \theta\right)\right. \tag{5}
\end{equation*}
$$

b) The configurations calculated on long-bridge site ( $X=a / 2, Y=0$ ) and short-bridge site $(X=0, Y=a /(2 \sqrt{2})$ ) (both sites have the same symmetry) are: (1) $\theta=0^{\circ}$; (2) $\theta=45^{\circ}$ and $\varphi=0^{\circ}$; (3) $\theta=45^{\circ}$ and $\varphi=90^{\circ}$; (4) $\theta=90^{\circ}$ and $\varphi=0^{\circ}$; (5) $\theta=90^{\circ}$ and $\varphi=90^{\circ}$. The Fourier functions of $\varphi$ interpolation on long-bridge and short-bridge sites are:

$$
\begin{gather*}
I_{0}(\varphi)=I(1)  \tag{6}\\
I_{45}(\varphi)=\frac{1}{2}(I(2)+I(3))+\frac{1}{2}(I(2)-I(3)) \cos 2 \varphi \tag{7}
\end{gather*}
$$

$$
\begin{equation*}
I_{90}(\varphi)=\frac{1}{2}(I(4)+I(5))+\frac{1}{2}(I(4)-I(5)) \cos 2 \varphi \tag{8}
\end{equation*}
$$

The Fourier function for $\theta$ interpolation on long-bridge (LB) and short-bridge ( SB ) sites is:

$$
\begin{equation*}
I_{L B, S B}(\theta, \varphi)=\frac{1}{4}\left(\left(I_{0}(\varphi)+2 I_{45}(\varphi)+I_{90}(\varphi)+2\left(I_{0}(\varphi)-I_{90}(\varphi)\right) \cos 2 \theta+\left(I_{0}(\varphi)-2 I_{45}(\varphi)+I_{90}(\varphi)\right) \cos 4 \theta\right)\right. \tag{9}
\end{equation*}
$$

c) The configurations calculated on the non-symmetrical top-hollow site $(X=a / 4, Y=a /(4 \sqrt{2}))$ are: (1) $\theta=0^{\circ}$; (2) $\theta=45^{\circ}$ and $\varphi=35.26$; (3) $\theta=45^{\circ}$ and $\varphi=90^{\circ}$; (4) $\theta=45^{\circ}$ and $\varphi=144.74^{\circ}$; (5) $\theta=45^{\circ}$ and $\varphi=215.26^{\circ}$; (6) $\theta=45^{\circ}$ and $\varphi=270^{\circ}$; (7) $\theta=45^{\circ}$ and $\varphi=324.74^{\circ}$; (8) $\theta=90^{\circ}$ and $\varphi=35.26^{\circ}$; (9) $\theta=90^{\circ}$ and $\varphi=90^{\circ}$; (10) $\theta=90^{\circ}$ and $\varphi=144.74^{\circ}$. The Fourier functions of $\varphi$ interpolation on the top-hollow site are:

$$
\begin{align*}
& I_{0}(\varphi)=I(1)  \tag{10}\\
& a_{0}=\frac{1}{16}(3(I(2)+I(4)+I(5)+I(7))+2(I(3)+I(6))) \\
& a_{1}=\frac{\sqrt{6}}{8}(I(2)-I(4)-I(5)+I(7)) \\
& a_{2}=\frac{3}{16}(I(2)+I(4)+I(5)+I(7)-2(I(3)+I(6))) \\
& b_{1}=\frac{1}{32 \sqrt{3}}(9(I(2)+I(4)-I(5)-I(7))+10 \sqrt{3}(I(3)-I(6))) \\
& b_{2}=\frac{3 \sqrt{2}}{16}(I(2)-I(4)+I(5)-I(7)) \\
& b_{3}=\frac{3}{32}(\sqrt{3}(I(2)+I(4)-I(5)-I(7))+2(I(6)-I(3))) \\
& I_{45}(\varphi)=a_{0}+a_{1} \cos \varphi+a_{2} \cos 2 \varphi+b_{1} \sin \varphi+b_{2} \sin 2 \varphi+b_{3} \sin 3 \varphi  \tag{11}\\
& I_{135}(\varphi)=a_{0}-a_{1} \cos \varphi+a_{2} \cos 2 \varphi-b_{1} \sin \varphi+b_{2} \sin 2 \varphi-b_{3} \sin 3 \varphi  \tag{12}\\
& I_{90}(\varphi)=\frac{1}{8}(3 I(8)+3 I(10)+2 I(9)+3 \sqrt{2}(I(8)-I(10)) \sin 2 \varphi+3(I(8)+I(10)-2 I(9)) \cos 2 \varphi) \tag{13}
\end{align*}
$$

The Fourier function for $\theta$ interpolation on the top-hollow (T-H) site is:

$$
\begin{align*}
I_{T-H}(\theta, \varphi)= & \frac{1}{4}\left(\left(I_{0}(\varphi)+I_{45}(\varphi)+I_{90}(\varphi)+I_{135}(\varphi)\right)+\left(I_{0}(\varphi)-I_{45}(\varphi)+I_{90}(\varphi)-I_{135}(\varphi)\right) \cos 4 \theta\right) \\
& +\frac{1}{2}\left(\left(I_{0}(\varphi)-I_{90}(\varphi)\right) \cos 2 \theta+\left(I_{45}(\varphi)-I_{135}(\varphi)\right) \sin 2 \theta\right) \tag{14}
\end{align*}
$$

Once the $(r, Z, \theta, \varphi)$ interpolation is performed, one is left with one $I^{6 D}$ value for each of the five sites that conform the DFT energy grid: top $I_{T}$, hollow $I_{H}$, long-bridge $I_{L B}$, short-bridge $I_{S B}$ and top-hollow $I_{T-H}$. Then, we use the 2 D scattered cubic spline interpolation from Ref. [2] to get values in between each of top, hollow, long bridge and short bridge sites giving us a rectangular grid of nine values. On this grid we use a 2 D periodic cubic spline as in Ref. 3 to get the final value $I^{6 D}(X, Y, Z, r, \varphi, \theta)$ from which we finally get the intended $V^{6 D}(X, Y, Z, r, \varphi, \theta)$ by adding the corresponding $V^{3 D}$ values.

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    [2] H. Akima, ACM Trans. Math. Softw. 22, 362 (1996).
    [3] M. Alducin, R. Díez Muiño, H. F. Busnengo, and A. Salin, J. Chem. Phys. 125, 144705 (2006).

