

Supporting Information

(1) The bulk structures of marcasite- and pyrite-type NiS₂

The marcasite-type bulk NiS₂ [Fig. S1(a)] has an orthorhombic lattice in $Pn\bar{m}$ symmetry with two Ni atoms and four S atoms, occupying the $2c$ (0.0, 0.5, 0.0) and $4g$ (0.1832, 0.1108, 0.0) Wyckoff positions, respectively. Each Ni atom is six-fold coordinated with six S atoms, while each S atom is four-fold coordinated with three Ni atoms and one S atom. The calculated lattice parameters $a = 4.60$ Å, $b = 5.57$ Å, $c = 3.55$ Å, with the bond lengths of $d_{\text{Ni-S}} = 2.38$ Å and $d_{\text{S-S}} = 2.09$ Å. We can obtain the P-NiS₂ monolayer by cleaving along the (100) plane.

The pyrite-type bulk NiS₂ [Fig. S1(b)] has a cubic lattice in $Pa\bar{3}$ symmetry with four Ni atoms and eight S atoms, occupying the $4b$ (0.5, 0.5, 0.5) and $8c$ (0.1069, 0.1069, 0.1069) Wyckoff positions, respectively. Each Ni atom is six-fold coordinated with six S atoms, while each S atom is four-fold coordinated with three Ni atoms and one S atom. The calculated lattice parameters $a = b = c = 5.62$ Å, with the bond lengths of $d_{\text{Ni-S}} = 2.37$ Å and $d_{\text{S-S}} = 2.08$ Å. The bulk NiS₂ is generally considered to be pyrite-type structure.^[1-5] We can obtain the O-NiS₂ monolayer by cleaving along the (001) plane.

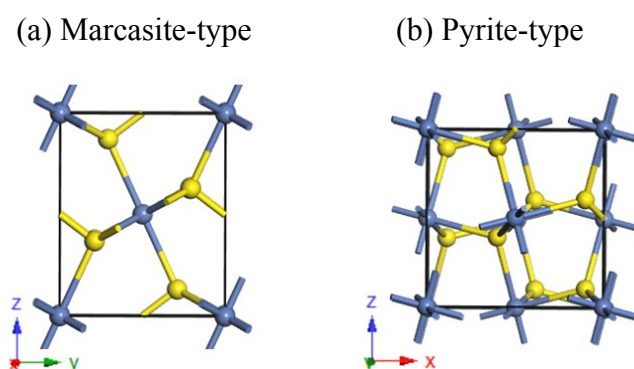


Fig. S1: The bulk structures of (a) marcasite-type NiS₂ and (b) pyrite-type NiS₂. The unit cell is marked by black lines. The blue and yellow balls represent the Ni and S atoms, respectively.

(2) The monolayer structures of O-, T- and H-NiS₂ ^[6-10]

The O-NiS₂ monolayer [Fig. S2(a)] has a monoclinic lattice in $P2_1/c$ symmetry with two Ni atoms and four S atoms, occupying the $2b$ (0.0, 0.0, 0.5) and $4e$ (0.1166,

0.3762, 0.4713) Wyckoff positions, respectively. Each Ni atom is four-fold coordinated with four S atoms, while each S atom is three-fold coordinated with two Ni atoms and one S atom, forming an intriguing pentagonal ring network known as the Cairo pentagonal tiling. The calculated lattice parameters $a = 5.22 \text{ \AA}$, $b = 5.33 \text{ \AA}$, the buckling height $h = 0.57 \text{ \AA}$, with the bond lengths of $d_{\text{Ni-S}} = 2.17, 2.18 \text{ \AA}$ and $d_{\text{S-S}} = 2.13 \text{ \AA}$.

The T-NiS₂ monolayer [Fig. S2(b)] has a trigonal lattice in $P\bar{3}m1$ symmetry with one Ni atom and two S atoms, occupying the $1b$ (0.0, 0.0, 0.5) and $2d$ (0.3333, 0.6667, 0.5583) Wyckoff positions, respectively. Each Ni atom is six-fold coordinated with six S atoms, while each S atom is three-fold coordinated with three Ni atoms. The calculated lattice parameters $a = b = 3.35 \text{ \AA}$, the buckling height $h = 1.17 \text{ \AA}$, with the bond lengths of $d_{\text{Ni-S}} = 2.26 \text{ \AA}$ and $d_{\text{S-S}} = 3.03 \text{ \AA}$.

The H-NiS₂ monolayer [Fig. S2(c)] has a hexagonal lattice in $P\bar{6}m2$ symmetry with one Ni atom and two S atoms, occupying the $1d$ (0.3333, 0.6667, 0.5000) and $2g$ (0.0000, 0.0000, 0.4474) Wyckoff positions, respectively. Each Ni atom is six-fold coordinated with six S atoms, while each S atom is three-fold coordinated with three Ni atoms. The calculated lattice parameters $a = b = 3.54 \text{ \AA}$, the buckling height $h = 1.05 \text{ \AA}$, with the bond lengths of $d_{\text{Ni-S}} = 2.30 \text{ \AA}$ and $d_{\text{S-S}} = 2.10 \text{ \AA}$.

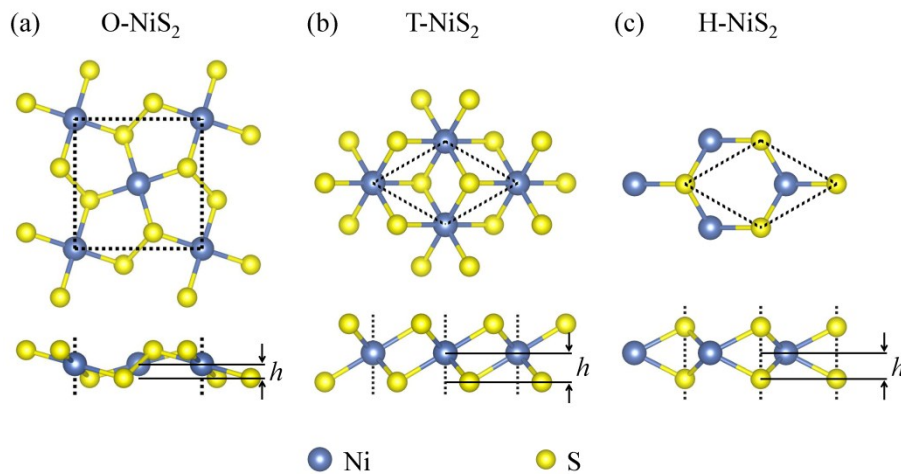


Fig. S2: Top and side views of (a) O-NiS₂ monolayer in $P21/c$ symmetry, (b) T-NiS₂ monolayer in $P\bar{3}m1$ symmetry, and (c) H-NiS₂ monolayer in $P\bar{6}m2$ symmetry. The unit cell is marked by black dashed lines. h is the buckling height. The blue and yellow

balls represent the Ni and S atoms, respectively.

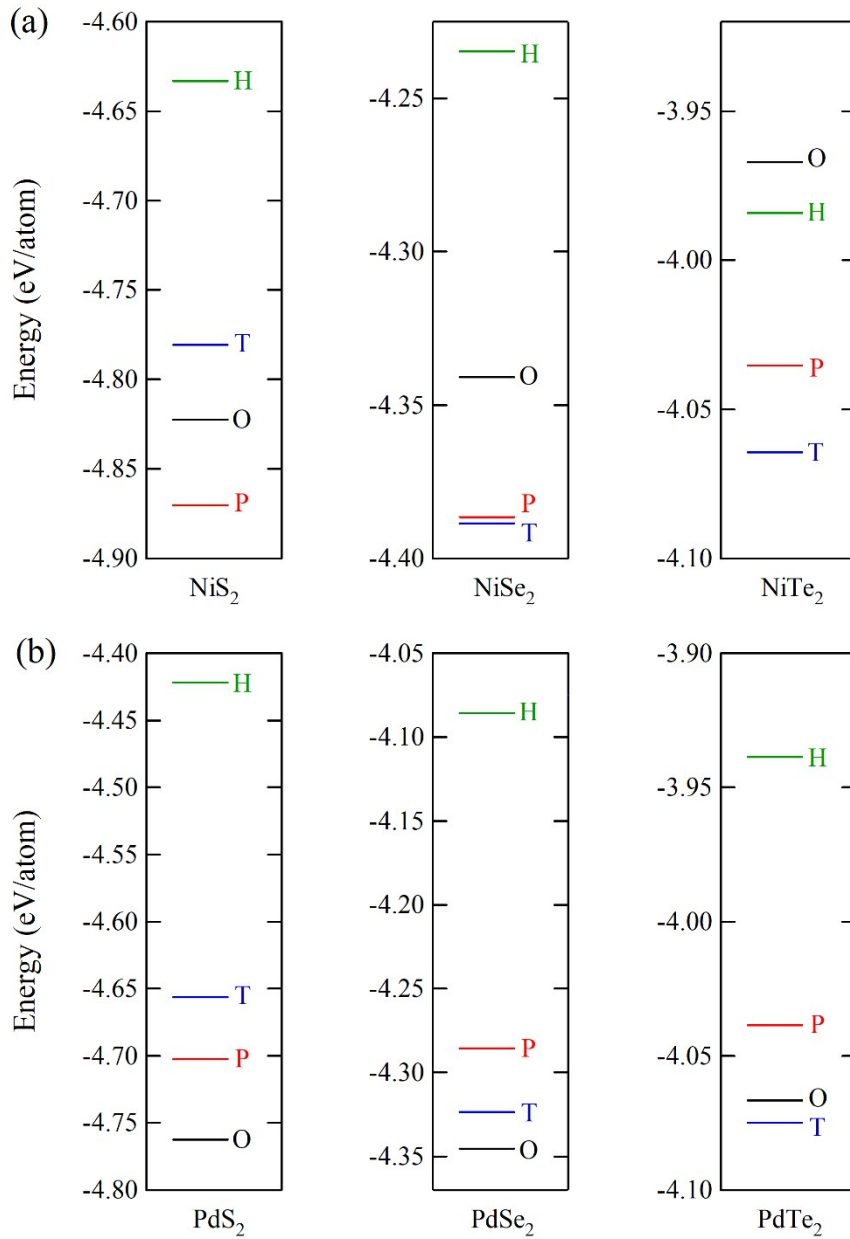


Fig. S3: Energy per atom for (a) NiX₂ and (b) PdX₂ (X = S, Se, Te) monolayers in P-, O-, T- and H-structures.

(3) Total energy per atom for NiX₂ and PdX₂ (X = S, Se, Te) monolayers

Fig. S3(a) presents the total energy per atom for P-NiX₂ (X = S, Se, Te) pentagonal network structures (red) in comparison with those for O-, T- and H-NiX₂ monolayers. For NiS₂, the energetic stability sequence is estimated to be: H-NiS₂ < T-NiS₂ < O-NiS₂ < P-NiS₂; for NiSe₂, the energetic stability sequence is estimated to be: H-NiSe₂

$< \text{O-NiSe}_2 < \text{P-NiSe}_2 < \text{T-NiSe}_2$; and for NiTe_2 , the energetic stability sequence is estimated to be: $\text{O-NiTe}_2 < \text{H-NiTe}_2 < \text{P-NiTe}_2 < \text{T-NiTe}_2$.

Fig. S3(b) presents the total energy per atom for P-PdX_2 ($X = \text{S, Se, Te}$) pentagonal network structures (red) in comparison with those for O-, T- and H- PdX_2 monolayers. For PdS_2 , the energetic stability sequence is estimated to be: $\text{H-PdS}_2 < \text{T-PdS}_2 < \text{P-PdS}_2 < \text{O-PdS}_2$; for PdSe_2 , the energetic stability sequence is estimated to be: $\text{H-PdSe}_2 < \text{P-PdSe}_2 < \text{T-PdSe}_2 < \text{O-PdSe}_2$; and for PdTe_2 , the energetic stability sequence is estimated to be: $\text{H-PdTe}_2 < \text{P-PdTe}_2 < \text{O-PdTe}_2 < \text{T-PdTe}_2$.

It is found that, for pentagonal network structure, the P-NiX_2 is always more favourable in energy than the O-NiX_2 , while the O-PdX_2 is always more favourable in energy than the P-PdX_2 is. If we consider all four possible structures, the T-structure becomes more stable in NiSe_2 , NiTe_2 , and PdTe_2 monolayers. Experimentally, multilayer O-PdSe_2 ,^[11] T-NiSe_2 ,^[12] T-NiTe_2 ,^[13] and T-PdTe_2 ^[14] have been successfully synthesized, which are in good agreement with our calculated results shown in Fig. S3.

References

- [1] J. Yin, Y. Li, F. Lv, M. Lu, K. Sun, W. Wang, L. Wang, F. Cheng, Y. Li, P. Xi and S. Guo, *Adv. Mater.*, 2017, **29**, 1704681.
- [2] N. Jiang, Q. Tang, M. Sheng, B. You, D.-E. Jiang and Y. Sun, *Catal. Sci. Technol.*, 2016, **6**, 1077-1084.
- [3] T. Wang, X. Guo, J. Zhang, W. Xiao, P. Xi, S. Peng and D. Gao, *J. Mater. Chem. A*, 2019, **7**, 4971-4976.
- [4] R. L. Kautz, M. S. Dresselhaus, D. Adler and A. Linz, *Phys. Rev. B*, 1972, **6**, 2078-2082.
- [5] C. Schuster, M. Gatti and A. Rubio, *Eur. Phys. J. B*, 2012, **85**, 325.
- [6] W. Xiong, K. Huang and S. Yuan, *J. Mater. Chem. C*, 2019, **7**, 13518-13525.
- [7] H. Yang, Y. Li, Z. Yang, X. Shi, Z. Lin, R. Guo, L. Xu, H. Qu and S. Zhang, *Vacuum*, 2020, **174**, 109176.
- [8] C. Ataca, H. Sahin and S. Ciraci, *J. Phys. Chem. C*, 2012, **116**, 8983-8999.
- [9] P. Miró, M. Ghorbani-Asl and T. Heine, *Angew. Chem. Int. Ed.*, 2014, **53**, 3015-3018.

- [10] J. A. Reyes-Retana and F. Cervantes-Sodi, *Sci. Rep.*, 2016, **6**, 24093.
- [11] A. D. Oyedele, S. Yang, L. Liang, A. A. Puretzky, K. Wang, J. Zhang, P. Yu, P. R. Pudasaini, A. W. Ghosh, Z. Liu, C. M. Rouleau, B. G. Sumpter, M. F. Chisholm, W. Zhou, P. D. Rack, D. B. Geohegan and K. Xiao, *J. Am. Chem. Soc.*, 2017, **139**, 14090-14097.
- [12] Y. Shao, S. Song, X. Wu, J. Qi, H. Lu, C. Liu, S. Zhu, Z. Liu, J. Wang, D. Shi, S. Du, Y. Wang and H.-J. Gao, *Appl. Phys. Lett.*, 2017, **111**, 113107.
- [13] B. Zhao, W. Dang, Y. Liu, B. Li, J. Li, J. Luo, Z. Zhang, R. Wu, H. Ma, G. Sun, Y. Huang, X. D. Duan and X. F. Duan, *J. Am. Chem. Soc.*, 2018, **140**, 14217-14223.
- [14] E. Li, R.-Z. Zhang, H. Li, C. Liu, G. Li, J.-O. Wang, T. Qian, H. Ding, Y.-Y. Zhang, S.-X. Du, X. Lin and H.-J. Gao, *Chin. Phys. B*, 2018, **27**, 086804.