# **Electronic Supplementary Information for**

## "Characterisation of the temperature-dependent $M_1$ to R phase

### transition in W-doped VO<sub>2</sub> nanorod aggregates by Rietveld

### refinement and theoretical modelling"

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(a)



(b)



(d)



(f)



(c)





**Fig. S1 (previous page)** Low-magnification TEM images of the samples fabricated under the optimised synthesis conditions as explained in the main manuscript: (a) pure  $VO_2$ ; (b)  $V_{0.995}W_{0.005}O_2$ ; and (c)  $V_{0.98}W_{0.02}O_2$ , and (d) high-magnification TEM image of a nano-rod in (c). The sample was hold at a room temperature for those four. (e) HAADF STEM image showing columns of strong contrast corresponding to heavier W atoms (highlighted with red circles) in a nanorod shown as (c). The inset shows the SAED pattern corresponding to the M<sub>1</sub> phase taken down <011>. (f) Energy dispersive spectroscopy spectrum of the region shown in (e) confirming presence of W atoms. The sample was kept at 100 K.





(a)





**Fig. S2** Lattice constant ratios  $a'_M/a^T_R^{end}$ ,  $b'_M/b^T_R^{end}$  and  $c'_M/c^T_R^{end}$  for (a) pure VO<sub>2</sub>, (b) V<sub>0.995</sub>W<sub>0.005</sub>O<sub>2</sub>, and (c) V<sub>0.98</sub>W<sub>0.02</sub>O<sub>2</sub>.



**Fig. S3 (a)** Crystal structure of  $M_1$ -type VO<sub>2</sub> with the unit cell indicated by solid lines, and (b) a 2 × 3 × 2 supercell used in calculations of W-doped VO<sub>2</sub>. Blue octahedra are centred on V atoms and the green octahedron on the W dopant atom.



**Fig. S4** Calculated partial DOSs of R-type  $V_{1-x}W_xO_2$  for (a) x = 0.0, (b) x = 0.0104, (c) x = 0.0208, and (d) x = 0.0416. DOSs per atom close to the Fermi level are shown as insets in c and d. The system remains metallic for all W contents examined.



**Fig. S5** Schematic diagram of the relation between the various parameters used in the BWC Mott IMT model.

**Table S1** Rietveld-refined structure parameters of R-type V<sub>0.98</sub>W<sub>0.02</sub>O<sub>2</sub> at 380 K in tetragonal space group  $P4_2/mnm$  with lattice constants a = b = 4.56787 Å, c = 2.86787 Å,  $\alpha = \beta = \gamma = 90^{\circ}$ .

Site	Coordinates	Occupancy
V1	0.000, 0.000, 0.000	0.98334
W1	0.000, 0.000, 0.000	0.01666
V2	2.284, 2.284, 1.434	0.98334
W2	2.284, 2.284, 2.284	0.01666
01	1.361, 1.361, 0.000	1.00000
02	3.207, 3.207, 0.000	1.00000
03	0.923, 3.645, 1.434	1.00000
04	3.645, 0.923, 1.434	1.00000

		Composition			
Phase	Parameter	$VO_2 V_{0.995}W_{0.005}O_2 V_{0.98}W_{0.02}O_2$			
R	$T_{end}$ (K)	355	340	250	
(tetragonal)	$a_{\rm R}$ (Å) (= $a_{\rm M} = a_{\rm M}'$ )	4.553 <sub>1</sub>	4.556 <sub>6</sub>	4.562 <sub>2</sub>	
	$b_{\rm R}$ (Å) (= $b_{\rm M}$ = $b_{\rm M}'$ )	4.553 <sub>1</sub>	4.556 <sub>6</sub>	4.562 <sub>2</sub>	
	$c_{\rm R}$ (Å) (= $c_{\rm M}$ = $c_{\rm M}$ )	2.851 <sub>6</sub>	2.8542	2.86439	
	$V_{\rm R}$ (Å <sup>3</sup> ) (= $V_{\rm M} = V_{\rm M}'$ )	59.11 <sub>5</sub>	59.26 <sub>1</sub>	59.61 <sub>8</sub>	
$M_1$	$T_{on}\left(\mathrm{K} ight)$	340	280	150	
(monoclinic)	$a_{\rm M}'(a_{\rm M})$ (Å)	4.851 <sub>1</sub> (5.756 <sub>6</sub> )	4.8511 (5.7540)	4.864 <sub>2</sub> (5.759 <sub>1</sub> )	
	$a_{\rm M}'^{-a} {\overset{ena}{R} \atop R} ({\rm \AA})$	0.2980	0.2944	0.3019	
	$\xi_{a}(T_{on})$ (%)	6.54	6.46	6.62	
	$b_{\rm M}' (= b_{\rm M}) ({\rm \AA})$	4.528 <sub>2</sub> (4.5282)	4.530 <sub>8</sub> (4.530 <sub>8</sub> )	4.5582 (4.5582)	
	$b_{\rm M}-b_{\rm R}^{ena}$ (Å)	$-0.024_{9}$	$-0.025_{8}$	$-0.004_{0}$	
	$\xi_{b}(T_{on})$ (%)	-0.55	-0.57	-0.09	
	$c_{\rm M}'(c_{\rm M})$ (Å)	2.2844 (5.3839)	2.288 <sub>8</sub> (5.383 <sub>4</sub> )	2.2967 (5.3799)	
	$c_{\rm M}'^{-c} \stackrel{ena}{R} ({\rm \AA})$	$-0.567_{1}$	$-0.565_{3}$	$-0.567_{2}$	
	$\xi_{c}(T_{on})$ (%)	-19.89	-19.73	-19.80	
	$V_{\rm M}'$ (Å <sup>3</sup> )	50.182	50.307	50.92 <sub>2</sub>	
	$V_{\rm M}'/V_{\rm R}^{ena}$	0.84887	0.84892	0.8542 <sub>8</sub>	
	$T_{-\infty}\left(\mathrm{K} ight)$	100	100	100	
	$a_{\rm M}'(a_{\rm M})({\rm \AA})$	4.8394 (5.7409)	$4.845_4 (5.746_1)$	4.864 <sub>8</sub> (5.760 <sub>5</sub> )	
	$a_{\rm M}'^{-a} {}^{ena}_{R} ({\rm \AA})$	0.286318	0.288804	0.302588	
	$\xi_{a}(T_{-\infty})$ (%)	6.2883912	6.338077	6.63250	
	$b_{\rm M}'(b_{\rm M})({\rm \AA})$	4.5254 (4.5254)	4.5329 (4.5329)	4.5582 (4.5582)	
	$b_{\rm M}' - b_{R}^{enu}({\rm \AA})$	-0.02768	-0.02728	-0.00407	
	$\xi_{b}(T_{-\infty})$ (%)	-0.60793478	-0.59868	-0.0892	
	$c_{\rm M}'(c_{\rm M})({\rm \AA})$	$2.290_5 (5.378_8)$	$2.292_0$ (5.380 <sub>6</sub> )	2.2943 (5.3793)	
	$c_{\rm M} - c_{\rm R} R^{-\alpha} ({\rm \AA})$	-0.561085	-0.562244	-0.569559	
	$\xi_{\rm c}(T_{-\infty})$ (%)	-19.6762	-19.6987	-24.8246	
	$V_{\rm M}'(V_{\rm M})({\rm \AA}^3)$	50.162 (117.79)	50.307 (118.09)	50.87 <sub>6</sub> (119.2 <sub>8</sub> )	
	$V_{\rm M}'/V_{R}^{enu}$	0.84866	0.84879	0.85350	

**Table S2** Structural parameters for the monoclinic (M1) and tetragonal (R) phases ofundoped and W-doped VO2 nanorods at different temperatures.

Note: Values in parentheses are for the monoclinic coordinate system.

**Table S3**: Volumetric ratios  $V_{\rm M}'/V_{\rm R}(T_{end})$  and the corresponding volumetric differentials,  $\Omega$ , for nanorods with three different W dopant levels spanning the corresponding temperature ranges  $T_{on} \le T \le T_{end}$ : (a) pure VO<sub>2</sub>, (b) V<sub>0.995</sub>W<sub>0.005</sub>O<sub>2</sub>, and V<sub>0.98</sub>W<sub>0.02</sub>O<sub>2</sub>.

VO <sub>2</sub>				
<i>T</i> (K)	$V_{\rm M}'/V_{\rm R}(T_{end})$	Ω (%)		
<i>T</i> <sub>on</sub> 340	0.8489	-15.11		
<i>T</i> <sub>1</sub> 345	0.9955	-0.4515		
$T_2$ 350	0.9995	-0.0530		
T <sub>end</sub> 355	1.00	0.00		

(a)

(b)

$V_{0.995}W_{0.005}O_2$				
<i>T</i> (K)	$V_{\rm M}'/V_{\rm R}(T_{end})$	Ω (%)		
<i>T</i> <sub>on</sub> 280	0.8489	-15.10		
<i>T</i> <sub>1</sub> 290	0.8494	-15.06		
$T_2$ 300	0.8499	-15.0 <sub>1</sub>		
<i>T</i> <sub>3</sub> 310	0.8526	-14.73		
<i>T</i> <sub>4</sub> 320	0.8599	-14.01		
<i>T</i> <sub>5</sub> 330	0.8712	-12.88		
T <sub>end</sub> 340	1.00	0.00		

$V_{0.98}W_{0.02}O_2$				
<i>T</i> (K)	$V_{\rm M}'/V_{\rm R}(T_{end})$	$\Omega\left(\% ight)$		
<i>T</i> <sub>on</sub> 150	0.8543	-14.57		
$\frac{T_I}{200}$	0.8573	-14.26		
$\begin{array}{c} T_2 \\ 210 \end{array}$	0.8579	-14.21		
T <sub>end</sub> 250	1.00	0.00		

(c)

D. (	Dopant content, x			
Parameter	0.0	0.0104	0.0208	0.0416
Upper Hubbard band width, $W(eV)$	2.345	2.489	2.650	2.828
Band splitting, $U_{eff}$ (eV)	3.349	3.219	3.121	3.023
$W/U_{eff}$	0.700	0.758	0.849	0.935
Lower Hubbard band width, $W_{low}$ (eV)	0.442	0.452	0.529	0.637
Charge-transfer energy, $\Delta$ (eV)	5.66	5.71	5.83	5.99
Band gap, $E_{gap}$ (eV)	0.510	0.382	0.217	0.076

**Table S4**: Hubbard-band parameters for V-3d orbitals extracted from DOSs of the M1phase and used in the BWC-Mott IMT model.

Notes:  $U_{eff}$  is the difference between the weighted average energy of the lower Hubbard band and that of the upper Hubbard band;  $\Delta$  is the gap between the bottom of the O 2p band and the bottom of the upper Hubbard band (conduction band minimum).

Dopant content. x	$T_{on}(\mathbf{K})$	W/U <sub>eff</sub>	1-W/U <sub>eff</sub>	Ω (%)	$(1-W/U_{eff})/\Omega$	$\Omega/(1-W/U_{eff})$
0.0	340	0.700	0.300	-0.1511	-1.9855	-0.2159
0.005	280	0 729	0 271	-01510	-1 7947	-0 2073
0.000		0.125	0.271	0.1010		0.2070
0.020	150	0.849	0.151	-0.1457	-1.0360	-0.9649

**Table S5** Ratios between  $1-W/U_{eff}$  and volumetric differential,  $\Omega$ , for different W contents, *x*.

Note:  $1-W/U_{eff}$  corresponds to the energy ratio from the apex of the hypothetical phase boundary between insulator and metal state in an f vs  $W/U_{eff}$  diagram.

#### **Conventions and notation used in the analysis**

#### (i) Lattice constants

The unit cell of the M<sub>1</sub> phase is monoclinic and can be described using the four lattice constants  $a_M$ ,  $b_M$ ,  $c_M$  and  $\beta$  (see Fig. S1), which we label collectively as  $\Lambda_M$ . The R phase is tetragonal and can be described using the three lattice constants  $a_R$ ,  $b_R$  and  $c_R$ , where  $a_R = b_R$ , represented collectively by  $\Lambda_R$  To compare the M<sub>1</sub> and R phases directly, we also calculated orthogonalised lattice parameters for the M<sub>1</sub> phase, as shown in Fig. S1, denoted by  $a'_M$ ,  $b'_M$  and  $c'_M$ , and collectively as  $\Lambda'_M$ .  $a'_M$ ,  $b'_M$ , and  $c'_M$ are related to  $a_M$ ,  $b_M$ ,  $c_M$  and  $\beta$  according to  $a'_M = a_M cos^{[M]}(\beta - 90)$ ,  $b'_M = b_M$  and  $c'_M = c_M - a_M sin^{[M]}(\beta - 90)$ .

#### (ii) IMT temperatures

We used two different procedures to identify the onset temperature,  $T_{on}$ , for the IMT: (a) the temperature at which the R (220) peak first appeared on the shoulder of the M<sub>1</sub> (022) peak in the XRD pattern, and (b) the temperature at which a noticeable change in the gradients of plots of orthogonalised lattice parameters of the M<sub>1</sub> phase on was observed. The endset temperature,  $T_{end}$ , was defined as the temperature where  $a_M' = b_M'$  (=  $a_R = b_R$ ) and  $\beta = 90^\circ$ .

The conventional phase change temperature,  $T_c$ , was taken as the mean of the onset and endset temperatures of the IMT, i.e.,  $T_c = (T_{on} + T_{end})/2$ . For temperature  $T_{-\infty}$  we took the lowest temperature below  $T_{on}$  accessible with our equipment, *viz.*, 100 K. This temperature was

sufficiently low for thermal expansion effects on the lattice parameters of the  $M_1$  phase to be ignored.

### (iii) Axial ratios

Ratios between the lattice constants in each axial direction of the M<sub>1</sub> phase at a given temperature and the lattice constants of the R phase at  $T_{end}$ ,  $\xi(T)$ , were calculated as

$$\xi_i(T) = \frac{\Delta \dot{i}_M(T)}{\dot{i}_R(T_{end})} \equiv \frac{\Delta \dot{i}_M(T)}{\dot{i}_M(T_{end})},$$
 (Eq. S1)

where i = a, b and c, and  $\Delta i_M(T)$  is the difference in (orthogonal) lattice parameter *i* between the M<sub>1</sub> phase at temperature *T* and that of the R phase at  $T_{end}$ .

### (iv) Volumetric differentials

The volumetric differential,  $\Omega$ , was defined as the difference in unit-cell volumes of the M<sub>1</sub> phase (using an orthogonalised basis) at temperature *T*,  $V_M'(T)$ , and the R phase at the endset temperature,  $V_R(T_{end})$ , divided by  $V_R(T_{end})$ , i.e.,

$$\Omega(T) = \frac{V'_{M}(T) - V_{R}(T_{end})}{V_{R}(T_{end})} = \frac{V'_{M}(T)}{V_{R}(T_{end})} - 1$$
(Eq. S2)

The unit-cell volume ratio  $V_M'(T)/V_R(T_{end})$  can be written as

$$\frac{V'_{M}(T)}{V_{R}(T_{end})} = \frac{\prod_{i=a}^{c} (i_{M}^{i}(T))}{\prod_{i=a}^{c} (i_{R}(T_{end}))} = \frac{a'_{M}(T)b'_{M}(T)c'_{M}(T)}{a_{R}(T_{end})b_{R}(T_{end})c_{R}(T_{end})}$$

$$(=\frac{a'_{M}(T)b'_{M}(T)c'_{M}(T)}{a_{R}(T_{end})b_{R}(T_{end})c_{R}(T_{end})})$$

$$=\frac{a'_{M}(T)}{a_{R}(T_{end})} \times \frac{b'_{M}(T)}{b_{R}(T_{end})} \times \frac{c'_{M}(T)}{c_{R}(T_{end})}$$

$$=L_{a}(T)L_{b}(T)L_{c}(T)$$

$$=\prod_{i=a}^{c} L_{i}(T)$$
(Eq. S3)

where  $L_i$  = is the ratio between the orthogonalised lattice constant of the M<sub>1</sub> phase and corresponding lattice constant of the R phase at  $T_{end}$  for each crystallographic axis *i*.

A negative  $\Omega$  indicates that the M<sub>1</sub> phase is compressed relative to the R phase, whereas a positive  $\Omega$  indicates that it is expanded.