Supplemental material

Strain modulated ferromagnetic phase transitions in monolayer FeCl₂ through exchange competitions: the first-principle and Monte Carlo simulations

Ya Yang^a*, Peiyin Guo^b, and Yongsong Luo^a*

^aKey Laboratory of Microelectronics and Energy of Henan Province, Henan Joint International Research Laboratory of New Energy Storage Technology, School of Physics and Electronic Engineering, Xinyang Normal University, Xinyang 464000, P. R. China

^bAnalysis & Testing Center, Xinyang Normal University, Xinyang 464000, P. R. China



Fig.S1 The phonon dispersion of no-strain, biaxial $\pm 5\%$ monolayer FeCl₂. A triangle cell was adopted and no imaginary frequency appears.



Fig.S2 The total energy of structures of nonmagnetic monolayer $FeCl_2$ depend on the strain (armchair, zigzag and biaxial). The no-strain structure is the lowest energy stable state.



Fig.S3 Orbital projected energy bands of (a) -5% armchair, (b) +5% armchair, (c) -5% zigzag, (d) +5% zigzag, (e) -5% biaxial and (f) +5% biaxial strains. The insulator properties were found no changing by strains.



Fig.S4 The energy without U of four magnetic states dependence of (a) armchair strain , (b) zigzag strain and (c) biaxial strain. (d) The four magnetic states sketch



Fig.S5 (a) The magnetic exchange parameter J dependence of strains from calculations without U. Monte Carlo simulations of (b) armchair strain, (c) zigzag strain and (d) biaxial strain. The M-T curves of \pm 5% and 0% structure are exhibited.



Fig.S6 The magnetoanisotropy energy dependence of the magnetic moment angles with (a) 0%, (b) -5% armchair, (c) -5% zigzag, (d) -5% biaxial, (f) +5% armchair, (g) +5% zigzag and (h) +5% biaxial strain. The magnetic moment angles are identified as (e). The magnetic easy-axial and easy-plane are presented by dotted lines. All the energy are calculated without the Hubbard U.



Fig.S7 The Monte Carlo simulation of the monolayer $CrI_{3.}$