## **Supporting Information**

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## ANALYSES OF THE INTERFACE REGIONS BASED ON TRADITIONAL "LAYERING AND IN-PLANE ORDERING" FRAMEWORK

Following **Figure S1~S3** are the analytical results based on the framework of "layering and in-plane ordering". Also included are the results of hyperbolic tangent fitting of two representative order parameters,  $aq_6$  and Frenkel number (FN).



Figure S1. Three representative snapshots, the decaying oscillation of number density, and the slice-wise profiles of  $aq_6$  and FN and their hyperbolic tangent fittings of FCC-Al based on the framework of "layering and in-plane ordering".



Figure S2. Three representative snapshots, the decaying oscillation of number density, and the slice-wise profiles of  $aq_6$  and FN and their hyperbolic tangent fittings of BCC-Fe based on the framework of "layering and in-plane ordering".



Figure S3. Three representative snapshots, the decaying oscillation of number density, and the slice-wise profiles of  $aq_6$  and FN and their hyperbolic tangent fittings of HCP-Mg based on the framework of "layering and in-plane ordering".

## ANALYSES OF INTERFACES BASED ON CBS FRAMEWORK

Following **Figure S4~S6** are the analytical results based on the framework of CBS. The properties contain the average atom number  $N_a$ , the average density  $\rho$ ; BOO parameters (i.e.  $q_4$ ,  $q_6$ ,  $w_4$ , and  $w_6$ ); ABOO parameters (i.e.  $aq_4$ ,  $aq_6$ ,  $aw_4$  and  $aw_6$ ); coordinate number (CN); and Frenkel number (FN). Also exhibited are the fraction distributions of four atom types.





**Figure S4.** Three representative snapshots of FCC-Al with the CBS layer-wise profiles of atom number  $N_a$ ; the density  $\rho$ ; BOO parameters (i.e.  $q_4$ ,  $q_6$ ,  $w_4$ , and  $w_6$ ); ABOO parameters (i.e.  $aq_4$ ,  $aq_6$ ,  $aw_4$  and  $aw_6$ ); coordinate number (CN); and Frenkel number (FN). Also exhibited are the fraction distributions of four atom types.



Figure S5. Three representative snapshots of BCC-Fe with the CBS layer-wise profiles of atom number  $N_a$ ; the density  $\rho$ ; BOO parameters (i.e.  $q_4$ ,  $q_6$ ,  $w_4$ , and  $w_6$ );

ABOO parameters (i.e.  $aq_4$ ,  $aq_6$ ,  $aw_4$  and  $aw_6$ ); coordinate number (CN); and Frenkel number (FN). Also exhibited are the fraction distributions of four atom types.



**Figure S6.** Three representative snapshots of HCP-Mg with the CBS layer-wise profiles of atom number  $N_a$ ; the density  $\rho$ ; BOO parameters (i.e.  $q_4$ ,  $q_6$ ,  $w_4$ , and  $w_6$ ); ABOO parameters (i.e.  $aq_4$ ,  $aq_6$ ,  $aw_4$  and  $aw_6$ ); coordinate number (CN); and Frenkel number (FN). Also exhibited are the fraction distributions of four atom types.

## ESTIMATION OF INTERFACIAL FREE ENERGY (IFE)

IFE of a SL interface can be defined as the free energy change to form an interface:

$$\gamma = \frac{\Delta F}{\Delta A} \tag{S1}$$

Where,  $\gamma$  is IFE,  $\Delta F$  is the Helmholtz free energy excess, and  $\Delta A$  is the area of the formed interface. For  $\Delta F$ , we have:

$$\Delta F = \Delta p + \Delta K - T_{\rm m} \Delta S = \Delta p - T_{\rm m} \Delta S \tag{S2}$$

Because the kinetic energy  $\Delta K = 0$  in equilibrium coexisting system (with uniform temperature), IFE is contributed by potential energy (PE)  $\Delta p$  and entropy (EN)  $T_{\rm m}\Delta S$ 

terms:  

$$\gamma = \frac{\sum(\langle \Delta f_i \rangle \cdot \langle n_i \rangle)}{\Delta A} = \frac{\sum((\langle \Delta p_i \rangle - \langle \Delta s_i \rangle \cdot T_m) \cdot \langle n_i \rangle)}{\Delta A} = \gamma_{\rm PE} + \gamma_{\rm EN}, \quad \text{where,}$$

$$\gamma_{\rm PE} = \frac{\sum(\langle \Delta p_i \rangle \cdot \langle n_i \rangle)}{\Delta A} \text{ is the contribution of PE; } \gamma_{\rm EN} = \frac{\sum(\langle \Delta s_i \rangle \cdot T_m \cdot \langle n_i \rangle)}{\Delta A} \text{ is the contribution of EN; } \Delta f_i = f_i - f_{\rm B}, \quad \Delta p_i = p_i - p_{\rm B}, \text{ and } \Delta s_i = s_i - s_{\rm B} \text{ represent the differences, between } i \text{ layer and the corresponding bulk phases, of atomic Helmholtz}$$
free energy, potential energy, and entropy, respectively;  $n_i$  is the atom number of  $i$  layer; the subscript  $i$  represents CBS layer L*i*, S*i*, IS or IL; and <...> means ensemble average.

So, according to the phase equilibrium, the entropy difference between BS and BL is the melting entropy:

$$\Delta s_{\rm m} = \Delta s_{\rm BS \to BL} = \frac{p_{\rm BL} - p_{\rm BS}}{T_{\rm m}}$$
(S3)

Where,  $p_{BL}$  and  $p_{BS}$  are potential energies of bulk solid and bulk liquid, respectively. **Table S1** lists the PEs of bulk phases and corresponding melting entropies.

**Table S1.** The potential energies (PEs) and resulted melting entropies. The meltingpoints are listed in **Table 2** in main text.

	   	FCC-A	Al I	1	BCC-F	e	HCP-Mg			
	$p_{\rm BL}$	$p_{\rm BS}$	$\Delta s_{\rm m}$	$p_{\rm BL}$	$p_{\rm BS}$	$\Delta s_{\rm m}$	$p_{_{ m BL}}$	$p_{\rm BS}$	$\Delta s_{\rm m}$	
	eV/atom	eV/atom	meV/(atom·K)	eV/atom	eV/atom	meV/(atom·K)	eV/atom	eV/atom	meV/(atom·K)	
[001] [0001]	-3.1519	-3.2588	0.1143	-3.6718	-3.8454	0.0951	-1.2863	-1.3794	0.1016	
[011] [ <b>1100</b> ]	-3.1531	-3.2578	0.1113	-3.6738	-3.8417	0.0914	-1.2856	-1.3793	0.1024	
[111] [ <sup>1210</sup> ]	-3.1530	-3.2566	0.1113	-3.6769	-3.8366	0.0869	-1.2869	-1.3793	0.1009	

If we pre-set the entropy of BL as zero, the relative entropy of BS would be the negative melting entropy. As we know, the Frenkel number (FN) is used to distinguish solid and liquid atoms based on the symmetry of bonding environment coming originally from the 6-fold rotational symmetry. So, it calibrates to some extent the symmetry of the bonding configuration of an atom and subsequently its phase state. Therefore, we calibrate the relative entropies of IS and IL layers as proportional to FN for the intermediate interfacial layers as follow:

$$\Delta s_{\rm IS} = \Delta s_{\rm BS} \times \frac{\rm FN(IS) - FN(BL)}{\rm FN(BS) - FN(BL)}; \Delta s_{\rm IL} = \Delta s_{\rm BS} \times \frac{\rm FN(IL) - FN(BL)}{\rm FN(BS) - FN(BL)}$$
(S4)

Then, the entropy differences of IS and IL can be estimated as listed in Table S2.

 Table S2. The ensemble-averaged FNs and the estimated entropy differences of four

 types of atoms relative to the pre-set zero entropy of BL, meV/(atom·K).

	Al						Fe		Mg				
	ENI		Δs				Δs		EN	$\Delta s$			
	FN	[001]	[011]	[111]	FN	[001]	[011]	[111]		[0001]	[1100]	[1210]	
BS	11.94	-0.1143	-0.1113	-0.1113	12.25	-0.0951	-0.0914	-0.0869	11.87	-0.1016	-0.1024	-0.1009	
IS	10.31	-0.0956	-0.0931	-0.0931	11.02	-0.0838	-0.0805	-0.0766	10.85	-0.0910	-0.0917	-0.0904	
IL	3.52	-0.0178	-0.0173	-0.0173	4.27	-0.0218	-0.0209	-0.0199	4.52	-0.0252	-0.0254	-0.0251	
BL	1.97	0	0	0	1.90	0	0	0	2.09	0	0	0	

In above table, relative to the pre-set zero of entropy of BL, the rests of BS, IS, and IL are all negative. As for the entropy changes of BS $\rightarrow$ IS, and BL $\rightarrow$ IL listed in **Table S3**, the calculation equations are:

$$\Delta s_{\rm BS \to IS} = \Delta s_{\rm IS} - \Delta s_{\rm BS}; \Delta s_{\rm BL \to IL} = \Delta s_{\rm IL} - \Delta s_{\rm BL} = \Delta s_{\rm IL}$$
(S5)

All  $\Delta s_{BS \to ISS}$  are positive and all  $\Delta s_{BL \to ILS}$  are negative, meaning, IS atoms are more disordered than their counterpart BSs and ILs are more ordered than their counterpart BLs.

	Al				Fe		Mg			
	[001]	[011]	[111]	[001]	[011]	[111]	[0001]	[1100]	[1210]	
$\Delta s_{\mathrm{BS} \rightarrow \mathrm{IS}}$	0.0187	0.0182	0.0182	0.0113	0.0109	0.0103	0.0106	0.0107	0.0105	
$\Delta s_{\mathrm{BL} \rightarrow \mathrm{IL}}$	-0.0178	-0.0173	-0.0173	-0.0218	-0.0209	-0.0199	-0.0252	-0.0254	-0.0251	

**Table S3.** Entropy changes of BS $\rightarrow$ IS, and BL $\rightarrow$ IL, meV/(atom K).

Then, the PE and EN contributions to IFE can be calculated based on the ensemble averages:

$$\gamma_{\rm PE} = \frac{\sum (\langle \Delta p_i \rangle \cdot \langle n_i \rangle)}{\Delta A}; \gamma_{\rm EN} = \frac{\sum (\langle \Delta s_i \rangle \cdot T_{\rm m} \cdot \langle n_i \rangle)}{\Delta A}$$
(S6)

**Table S4** lists all data needed to calculate  $\gamma_{PE}$ . As shown in it, for liquid part,  $\langle p_{L1} \rangle_s$ and  $\langle p_{IL} \rangle_s$  of Al and Mg, as well as  $\langle p_{L1} \rangle_s$  of Fe, almost equal to corresponding  $\langle p_{BL} \rangle_s$ , and only  $\langle p_{IL} \rangle_s$  of Fe have a little drop relative to corresponding  $\langle p_{BL} \rangle_s$ ; for solid part, all  $\langle p_{S2} \rangle_s$  equal to corresponding  $\langle p_{BS} \rangle_s$ , and finally  $\langle p_{IS} \rangle_s$  and  $\langle p_{S1} \rangle_s$  have notable raises relative to corresponding  $\langle p_{BS} \rangle_s$ , especially  $\langle p_{IS} \rangle_s$ .

	<b>I able S4.</b> Potential energy contribution to IFE from														
		   	$\langle p_i \rangle_{, eV/atom}$									AA 80	γ <sub>PE,</sub>		
		$\langle p_{_{ m BL}} \rangle$	$\left< p_{\rm L1} \right>$	$\langle p_{_{\rm IL}} \rangle$	$\langle p_{\rm IS} \rangle$	$\left< p_{\rm S1} \right>$	$\left< p_{\rm S2} \right>$	$\left< p_{\rm BS} \right>$	$\langle n_{\rm L1}^{} \rangle$	$\langle n_{_{\rm IL}} \rangle$	$\langle n_{\rm IS} \rangle$	$\langle n_{_{\rm S1}} \rangle$	$\langle n_{_{ m S2}} \rangle$	Δ.1, A <sup>2</sup>	mJ/m <sup>2</sup>
	[001]	-3.67	-3.68	-3.70	-3.79	-3.82	-3.84	-3.85	5100.34	6142.52	6351.39	5694.04	5172.92	20971.95	253.82
Fe	[011]	-3.67	-3.68	-3.70	-3.79	-3.82	-3.84	-3.84	5521.00	6561.07	6777.92	6114.35	5527.80	23036.15	233.03
	[111]	-3.68	-3.68	-3.70	-3.79	-3.82	-3.83	-3.84	6736.39	8003.89	8278.01	7462.68	6755.28	28149.86	192.57
Al	[001]	-3.15	-3.15	-3.15	-3.21	-3.25	-3.26	-3.26	6291.93	6796.57	6771.98	6207.89	5773.70	35282.78	175.52
	[011]	-3.15	-3.15	-3.15	-3.22	-3.25	-3.26	-3.26	6552.77	7054.90	7095.88	6634.57	6317.28	37284.77	161.58
	[111]	-3.15	-3.15	-3.15	-3.22	-3.25	-3.25	-3.26	5682.54	6106.61	6146.67	5718.28	5399.88	32676.98	154.57
Mg	[0001]	-1.29	-1.28	-1.29	-1.35	-1.37	-1.38	-1.38	5705.03	6577.66	6680.21	5900.19	5398.49	37701.10	114.61
	[1100]	-1.29	-1.28	-1.29	-1.35	-1.38	-1.38	-1.38	9080.47	10285.43	10399.07	9272.30	8655.71	33444.15	153.36
	[1210]	-1.29	-1.28	-1.29	-1.35	-1.38	-1.38	-1.38	5142.38	5810.67	5897.31	5256.43	4895.23	33497.13	91.04

**Table S4.** Potential energy contribution to IFE from  $\gamma_{\rm PE} = \frac{\sum (\langle \Delta p_i \rangle \cdot \langle n_i \rangle)}{\Delta A}$ .

As for  $\gamma_{\text{EN}}$ , the FNs, as shown in **Figure S4~S6**, of all L*i* and S*i* layers, except IL and IS, have their own Bulk values. Only IL and IS deviate from the bulk references. Moreover, the positive  $\Delta s_{\text{BS}\to\text{IS}}$  compensates largely the negative  $\Delta s_{\text{BL}\to\text{IL}}$  as shown in **Table S3**, which reduced mostly the contribution of entropy to IFE.  $\Delta s_{BS \to IS}$  and  $\Delta s_{BL \to IL}$  in **Table S3**,  $\langle n_i \rangle$  and  $\Delta A$  in **Table S4**, and  $T_m$  in **Table 2** of main text provide all data for the estimation of  $\gamma_{EN}$ . **Table S5** lists the final results of  $\gamma_{PE}$ ,  $\gamma_{EN}$ , and  $\gamma$ .

Table S5. The calculated PE and EN contributions to IFE,  $mJ/m^2$ .

		$\gamma_{\mathrm{PE}}$			$\gamma_{\rm EN}$		$\gamma = \gamma_{\rm PE} + \gamma_{\rm EN}$			
Element	[100] or	[110] or	[111] or	[100] or	[110] or	[111] or	[100] or	[110] or	[111] or	
	[0001]	[1100]	[1210]	[0001]	[1100]	[1210]	[0001]	[1100]	[1210]	
BCC-Fe	253.82	233.03	192.57	86.65	81.13	76.87	340.47	314.16	269.44	
FCC-Al	175.52	161.58	154.57	-2.68	-3.41	-2.84	172.84	158.17	151.73	
HCP-Mg	144.61	153.36	91.04	37.27	65.92	36.57	181.88	219.28	127.61	