

Electronic supplementary information

A pyrrolopyridazinedione-based copolymer for fullerene-free organic solar cells

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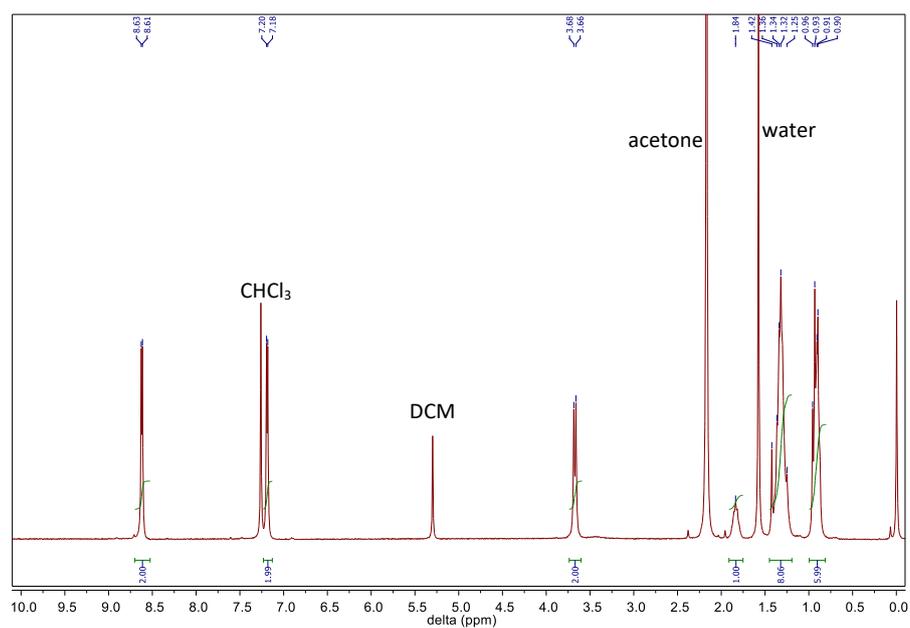


Fig. S1 $^1\text{H-NMR}$ spectrum of 1,4-bis(5-bromothiophen-2-yl)-6-(2-ethylhexyl)-5H-pyrrolo[3,4-d]pyridazine-5,7(6H)-dione (T-EHPPD-T monomer).

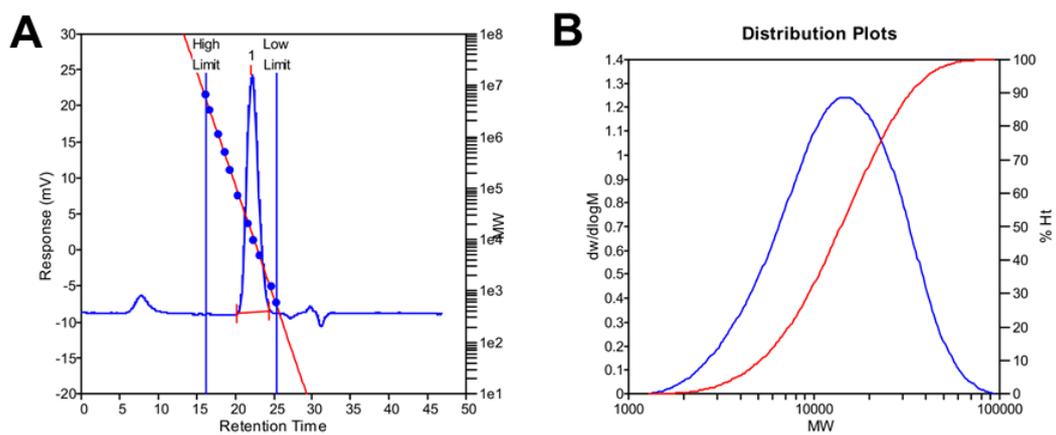


Fig. S2 (A) GPC chromatogram and (B) the corresponding distribution plots of T-EHPPD-T-EHBDT.

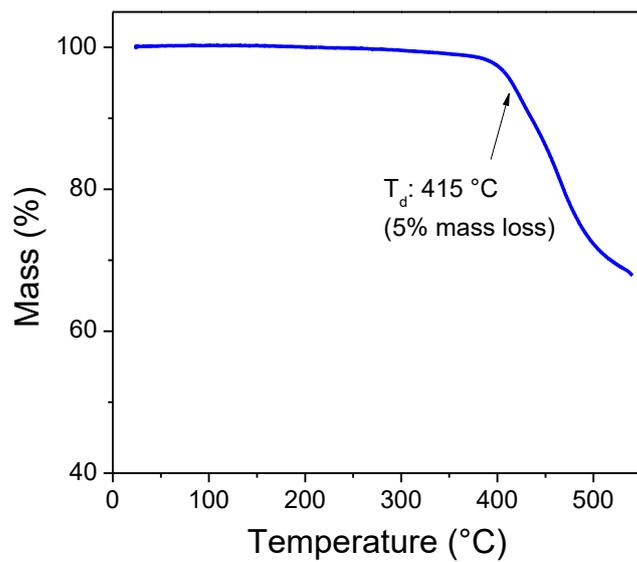


Fig. S3 Thermogravimetric analysis of the polymer T-EHPPD-T-EHBDT.

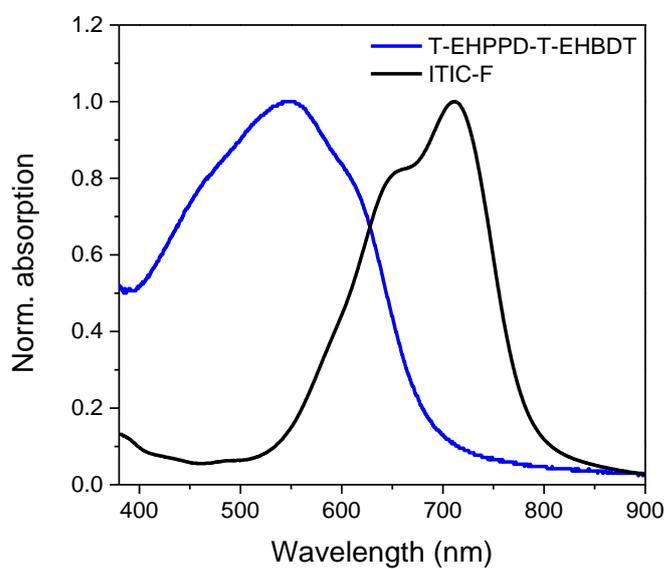


Fig. S4 Normalized UV-Vis absorption spectra of T-EHPPD-T-EHBDT and ITIC-F thin films.

Table S1 Characteristic parameters of the solar cells prepared with different layer thicknesses (D: donor, A: acceptor, L: absorber layer thickness). The average solar cell parameters have been calculated from the best four devices

D:A weight ratio	L (nm)	V_{oc} (V)	J_{sc} (mA/cm ²)	FF	PCE (%)
1:1	50 ± 5	0.74 ± 0.04	15.6 ± 0.8	0.49 ± 0.03	5.7 ± 0.4
	60 ± 5	0.78 ± 0.04	16.1 ± 0.2	0.58 ± 0.03	7.2 ± 0.4
	90 ± 10	0.78 ± 0.04	16.4 ± 0.7	0.46 ± 0.02	5.9 ± 0.4
	140 ± 10	0.74 ± 0.04	14.5 ± 0.6	0.39 ± 0.02	4.2 ± 0.3
1:1.5	50 ± 5	0.75 ± 0.04	16.2 ± 0.8	0.47 ± 0.03	5.7 ± 0.3
	60 ± 5	0.74 ± 0.04	15.7 ± 0.8	0.50 ± 0.03	5.9 ± 0.3
	95 ± 10	0.77 ± 0.04	15.8 ± 0.7	0.45 ± 0.03	5.5 ± 0.4
	140 ± 10	0.78 ± 0.04	12.4 ± 0.6	0.44 ± 0.02	4.1 ± 0.2

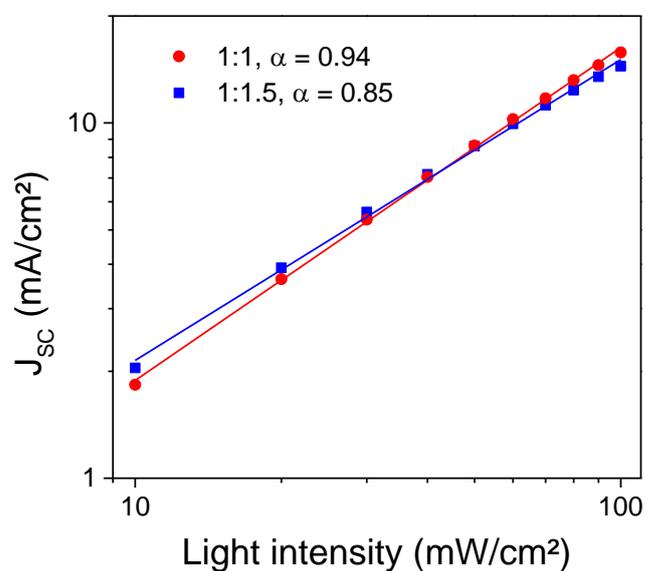


Fig. S5 Dependence of the J_{sc} on the light intensity.

Determination of the SCLC charge carrier mobility of the blend films

In addition to the OFET-based mobility measurements, an analysis of the dark current in the organic solar cells was performed to obtain the SCLC charge carrier mobility of the blend films. The j-V characteristics used for the analysis are presented in Fig. S6.

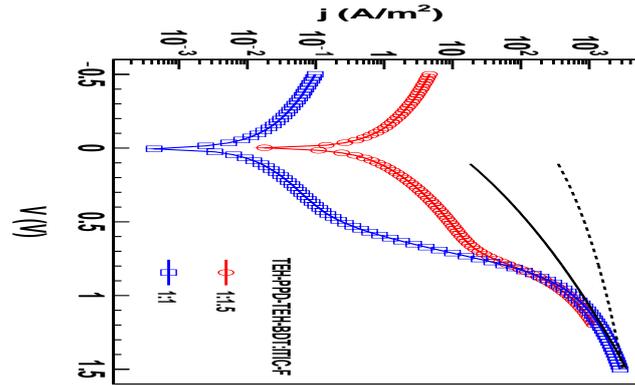


Fig. S6 j-V characteristics measured in the dark of solar cells with polymer:ITIC-F weight ratios of 1:1 and 1:1.5.

Since the anode (ZnO/ITIC-F) and cathode (T-EHPPD-T-EHBDT/MnO₃) interfaces represent a low injection barrier for electrons and holes, respectively, the observed j-V dependence is limited by the bimolecular recombination. Due to the low film thickness ($L = 60$ nm), the transit time of injected carriers is presumably shorter than the dielectric relaxation time. Therefore, the current density can be described as:^{1,2}

$$j = 2\varepsilon_r\varepsilon_0 \frac{(V - V_{bi})^2}{L^3} \sqrt{\frac{\beta_L}{\beta}} \mu_h \mu_e$$

where (ε_0) is the vacuum permittivity and ε_r is a relative permittivity taken as 3. The externally applied bias voltage (V) is reduced by the built-in voltage (V_{bi}), which we found to be -0.60 ± 0.02 V in this case. β and β_L are the bimolecular and Langevin recombination rates, respectively. However, in the case when one of the injection barriers is limiting carrier injection, j is not limited by bimolecular recombination. Therefore, the observed j-V dependence can be interpreted in terms of the space-charge limited current (SCLC) model of Murgatroyd and Gill:^{3,4}

$$j = \frac{9}{8} \varepsilon_r \varepsilon_0 \frac{(V - V_{bi})^2}{L^3} \mu(V) \quad \mu(V) = \mu_0 \exp\left(0.891\gamma \sqrt{\frac{V - V_{bi}}{L}}\right)$$

which is an extended version of the Mott-Gurney law to account for the electric field dependence of the mobility ($\mu(V)$). In the model, the mobility is represented as a product of zero-field mobility (μ_0) and the exponential term including the parameter γ . In Fig. S7, we show j as a function of externally applied bias voltage (V) by red open circles. j is linearly increasing with V below 0.2 V. The corresponding slope of the j - V characteristics (dashed blue line and right axis) is approximately 1. In contrast, above 0.2 V, j exhibits a rapid increase, which is a typical trap-filling effect. The slope in this voltage region reaches a value between 16 and 17 in Fig. S7A (blend ratio = 1:1) and a value between 10 and 11 in Fig. S7B (blend ratio = 1:1.5). For higher bias, the slope drops towards 2, which is a fingerprint of the Mott-Gurney law. Therefore, the model was used to estimate V_{bi} , μ_0 and γ in the bias above 1 V.⁵ The resulting parameters are presented in Table S2.

Table S2 Parameters obtained based on the above described model

blend ratio	V_{bi} (V)	μ_0 (cm^2/Vs)	γ ($\text{m}^{1/2}/\text{V}^{1/2}$)
1:1	-0.587	2.9×10^{-8}	1.4×10^{-3}
1:1.5	-0.602	5.6×10^{-10}	2.2×10^{-3}

The parameters display a significant variation, which results from the analysis in a relatively narrow voltage range. Therefore, we calculated the SCLC mobility in the blend layer at a bias voltage of 1.2 V. The corresponding mobility is $2.8 \times 10^{-5} \text{ cm}^2/\text{Vs}$ and $2.7 \times 10^{-5} \text{ cm}^2/\text{Vs}$ for the 1:1 and the 1:1.5 blend, respectively. In case of double-injection, the geometric average of mobility $\sqrt{\mu_h \mu_e}$ is found to be $1.6 \times 10^{-5} \text{ cm}^2/\text{Vs}$ multiplied by a factor $\sqrt{\beta/\beta_L}$. Regarding the latter, it was found to be a fraction of unity in fullerene-based solar cells.⁶ However, recent kinetic Monte Carlo simulations predict that the recombination rate decreases with time.⁷ Hence, it is significantly higher for non-thermalized injected carriers. In our case with reduced layer thickness, the injected carriers are presumably non-thermalized. Therefore, we assume that the recombination rate is of the order of Langevin type and hence the geometric average of the mobility is in the order of $10^{-5} \text{ cm}^2/\text{Vs}$.

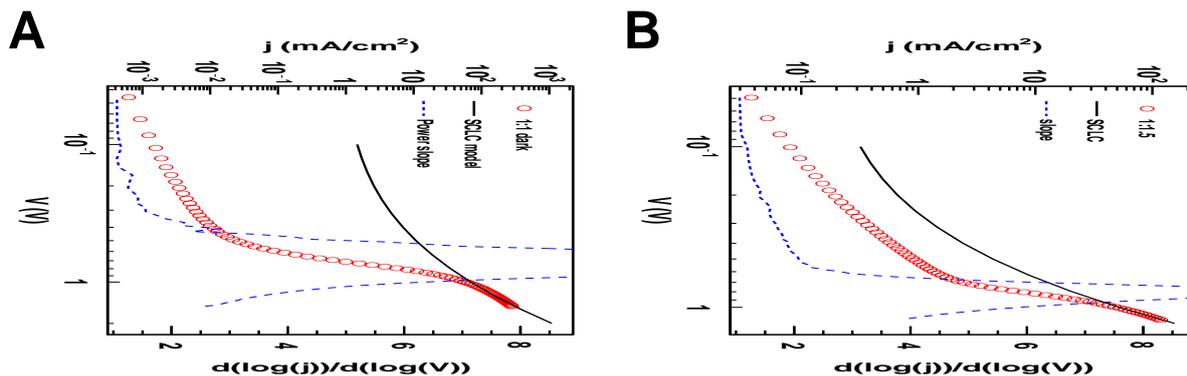


Fig. S7 Hole-only current density as a function of applied bias voltage of the blend films in a double-logarithmic plot (A: 1:1 blend ratio; B: 1:1.5 blend ratio). Measurements (open red circles) and SCLC model (black solid line) obtained with a least-square-method in the range above 1 V. The power slope of the measurements is presented with a dashed blue line and the corresponding scale on the right axis.

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