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Sliding dynamics of multi-rings on semiflexible polymer in poly[*n*]catenanes

Received 00th 20xx, Accepted 00th 20xx

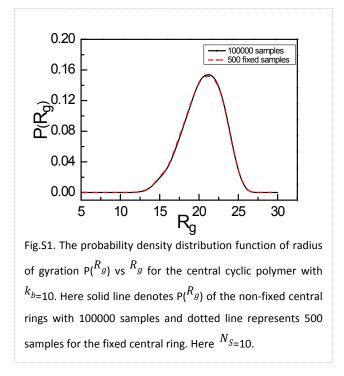
DOI: 10.1039/x0xx00000x

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S1. The probability density distribution function P(S) of central cyclic polymer

In our simulations, we first calculated the probability density distribution function of radius of gyration $P_0(^{R_g})$ with different bending energy of central rings in equilibrated state, and made them as the criterion for sampling the fixed rings. We choose 500 equilibrated conformations at random and then also calculate the probability density distribution function of radius of gyration $P(^{R_g})$.

If the relative difference between $P_0(^{R_g})$ and $P(^{R_g})$ is less than 5% for different R_g , i.e, $|P_0(^{R_g})-P(^{R_g})|/P_0(^{R_g})<5\%$ for $0<^{R_g}<^{R_g,max}$, these 500 initial central ring polymers can be fixed and the diffusion coefficients of sliding dynamics are averaged in those 500 simulation runs. Figure S1 shows that the distribution function $P(^{R_g})$ of our 500 fixed samples as a function of R_g for the ring with $^{k_b}=10$.



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S2. Persistence length $({}^{l_{p}})$ and bond vector correlation function $({}^{C_{1}})$ of central ring polymer

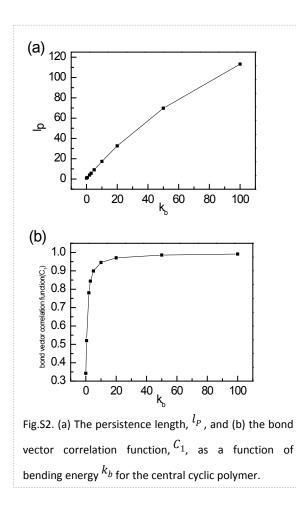
Persistence length is an important parameter to describe the conformations of semiflexible polymer. We calculated the persistence length of central cyclic polymer according to the exponential decay of orientation correlation in a chain $^{1-3}$

$$\langle u(r) \cdot u(0) \rangle = e^{-r/l_p} \tag{S1}$$

Where $\mathbf{u}(r)$ and $\mathbf{u}(0)$ are the unit bond vectors which are separated by a contour distance r along the chain, and the brackets <> represent an average of overall starting position and chain configurations. The bond vector correlation function C_1 is defined as⁴

 $C_1 = < b_i \cdot b_{i+1} >$

(S2)



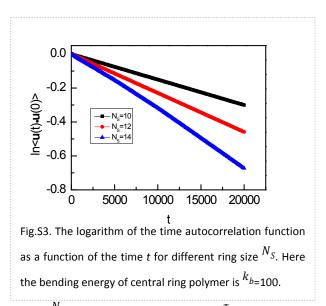
 b_i and b_{i+1} are the bond vectors of i-th and (i+1)-th beads of the central ring polymer. In fact, $C_1 = \langle \cos \theta' \rangle$, here θ' is the angle between two adjacent bond vectors of central cyclic polymer. Fig. S2 shows the results for (a) persistence length l_p and (b) bond vector correlation function C_1 , as a function of the bending energy k_b of the central ring polymer.

S3. The rotational diffusion of rigid central ring polymer

In order to analysis the rotational diffusion of the central cyclic polymer, we computed the time autocorrelation function $\langle \vec{u}(t) \cdot \vec{u}(0) \rangle$ of a unit vector $\vec{u}(t)$ directed along the longest chain dimension. Here $\vec{u}(t)$ is taken to be along the diameter vector $\vec{u}_d(t) = \vec{u}_{N_L/2}(t)$ for the ring polymer⁵. We calculated the time autocorrelation function for the stiff central cyclic polymer with k_b =100, and the results are shown in Figure S3. Through a Kohlrausch-Williams-Watts (KWW) function of the form⁵:

$$\langle \vec{u}(t) \cdot \vec{u}(0) \rangle = exp^{[m]} \left[-\left(\frac{t}{\tau_{KWW}}\right)^{\beta} \right]$$
(S3)

where τ_{KWW} and β are the characteristic relaxation time and stretching parameter, respectively, we can calculate the relaxation time τ_{KWW} which is estimated from the reciprocal of the slope of corresponding curve in Figure S3. According to Table S1, with the



ring size ${}^{N_{S}}$ increasing, the relaxation time ${}^{ au_{KWW}}$ decreases, which

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indicates that the central cyclic polymer shows a faster rotational

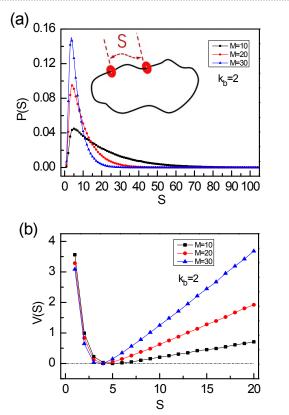


Fig. S4. (a)Probability distribution P(S) of two neighboring rings for different ring numbers M (b) Potential of mean force V(S) for small rings with $k_{b}=2$ for different M. Inset figure defines the path S between two neighboring rings.

diffusion for large N_s . These results also prove that the friction between the ring and central polymer is weakened as increasing small ring size.

Table S1. The relaxation time τ_{KWW} of rigid central cyclic polymer for different small ring sizes N_S .

N _S	10	12	14
$ au_{KWW}$	66500	43500	29700

S4. Potential of the mean force (PMF) V(S)

The probability distribution P(S) for rings along the central cyclic polymer is presented in Fig. S4(a), where *S* represents the path between two neighbouring rings, see the inset figure in Fig. S4(a).

Here the rings don't penetrate each other during the whole simulation because there is the same size for all small rings. As the small rings in poly[*n*]catenane are marked in number order initially, *S* is defined as

$$S_m = \begin{cases} i_{m+1} - i_m, & i_{m+1} - i_m \ge 0\\ N_L + (i_{m+1} - i_m), & i_{m+1} - i_m < 0 \end{cases}$$
(S4)

Where i_m is the index number of the monomer of the polymer closest to the *m*-th ring, and $i_{M+1} = i_1$. Meanwhile, the peak value of *P*(S) is enhanced with increasing the ring number from *M*=10 to 20, and 30. The ring-ring potential of the mean force (PMF) can be calculated by $V(S) = -k_BTlnP(S)$, ⁶⁻⁸ which represents the ring-ring potential as a function of *S*. Here, we take the minima of the PMF as the reference states.

S5. Effects of the threading ring polymer on static properties of the central cyclic polymer

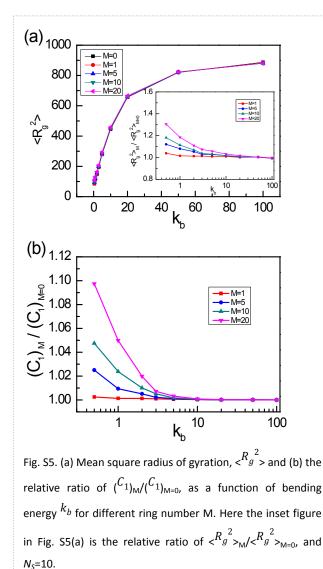
The static properties of the central cyclic polymer are evaluated for various numbers of the threading rings M. Figs. S5(a) and (b) show the mean-square radius of gyration, $\langle R_g^2 \rangle$, and the relative ratio of bond vector correlation function of the central ring polymer, $({}^{C_1})_{\rm M=0}$, for various bending energies k_b with different *M*, respectively. Here, the mean-square radius of gyration $\langle R_g^2 \rangle$ is

$$<\!R_g^2\!>=\!\frac{1}{N_L}\!<\!\sum_{i=1}^{N_L}\!(r_i\!-\!r_{cm})^2\!>$$
 defined as

, where r_{cm} denotes

the center-of-mass of the central polymer, N_L is the number of monomers in polymer. It is shown that the presence of the threading rings does not greatly affect the statistics properties of the central cyclic chain except for the flexible central ring polymer, i.e., $k_b < 2$.

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