

## Electronic Supporting Information

### Solvent-free dynamic nuclear polarization enhancements in organically modified mesoporous silica

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Table S1 – Weights of Fmoc-Gly-OH and 3-Carboxy-PROXYL precursors used in the synthesis. The overall mass of the final materials (yield mass) is also given.

Sample	Weigh Fmoc-Gly-OH [mg]	Weigh 3-Carboxy-PROXYL [mg]	Yield [mg]
99.5Fmoc/0.5Prx	79	0.3	54
95Fmoc/5Prx	75	2.5	59
90Fmoc/10Prx	71	5	58
75Fmoc/25Prx	59	12	50
50Fmoc/50Prx	40	25	64

Table S2: Median pore diameters calculated by the BJH (left) and NLDFT (right) method of the prepared SBA-15 with APTES linker and functionalized mesoporous silica materials as obtained from BET analysis.

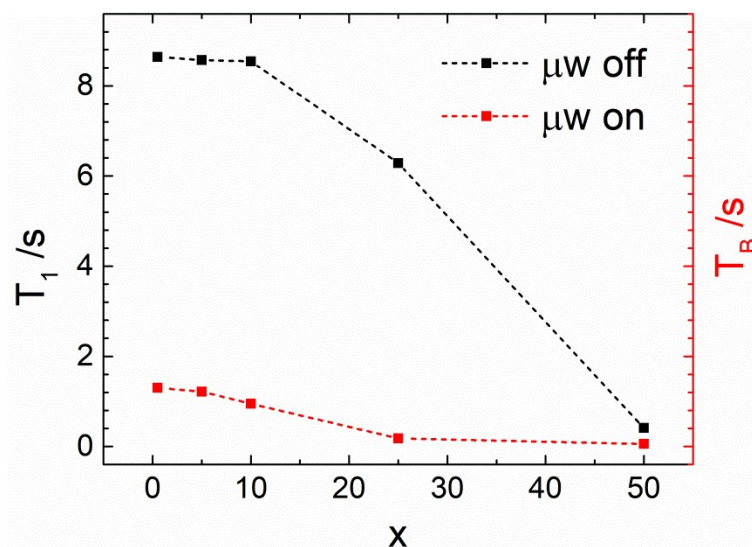
Sample	Pore diameter by BJH / nm	Pore diameter by NLDFT / nm
SBA-15 with APTES	5.5	7.9
99.5Fmoc/0.5Prx	5.3	6.6
95Fmoc/5Prx	5.3	7.0
90Fmoc/10Prx	5.4	6.6
75Fmoc/25Prx	5.3	6.6
50Fmoc/50Prx	5.2	6.6

Table S3. Pore volume (Gurvich at  $p/p^0$  0.95), specific surface area (BET), nitrogen loading and carbon/nitrogen ratio of the support material determined by EA.

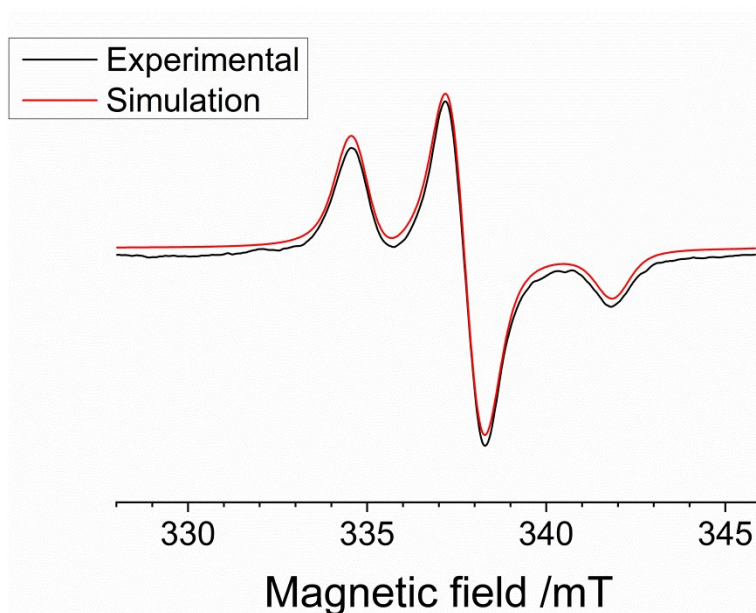
Sample	Pore Volume (BET) / $\text{cm}^3 \cdot \text{g}^{-1}$	Surface Area (BET) / $\text{m}^2 \cdot \text{g}^{-1}$	Nitrogen loading (EA) / $\text{mmol} \cdot \text{g}^{-1}$	C/N ratio (Theo.)	C/N ratio (Exp.)
SBA-15 with APTES	0.72	537	$1.78 \pm 0.3$	3.00	$3.13 \pm 0.5$

Table S4 - Elemental Analysis. Two measurements were performed for each sample.

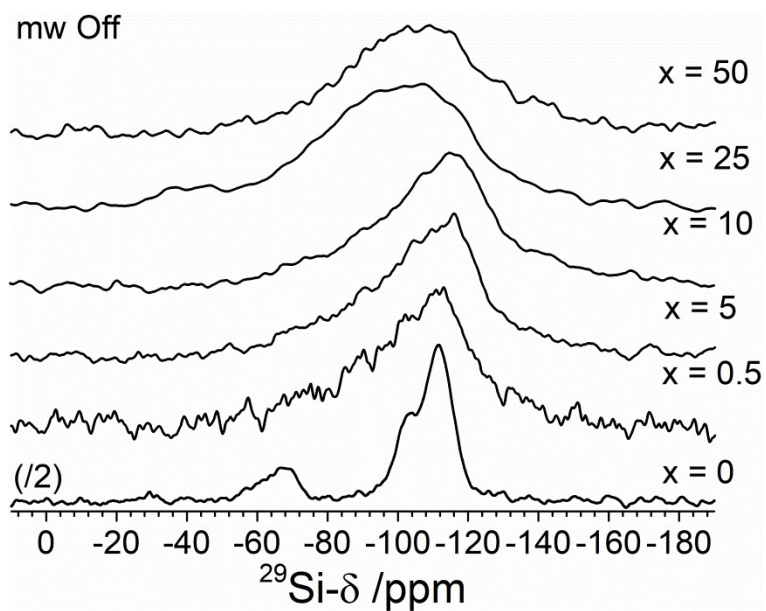
Sample	C/N ratio	Carbon (%)	Hydrogen (%)	Nitrogen (%)
Amine functionalized SBA-15	2.835	8.018	3.11	2.829
Amine functionalized SBA-15	2.842	7.728	3.399	2.719
X 0.5	3.611	17.24	3.013	4.775
X 0.5	3.65	17.1	2.859	4.685
X 5	3.408	17.05	3.271	5.004
X 5	3.29	16.71	3.682	5.079
X 10	3.68	16.59	3.101	4.509
X 10	3.511	16.39	3.108	4.667
X 25	3.092	15.48	3.291	5.006
X 25	3.112	15.43	3.349	4.958
X 50	3.266	15.51	3.024	4.747
X 50	3.268	14.96	3.182	4.578



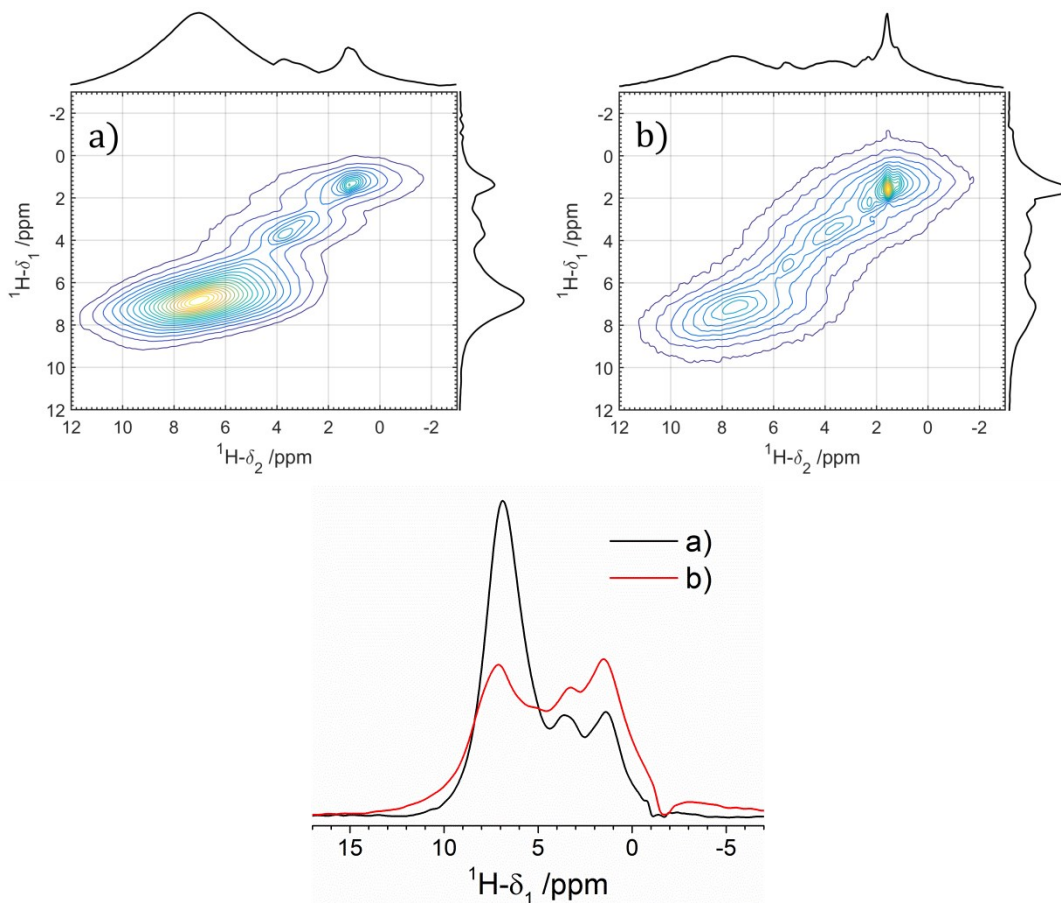
**Figure S1** –  $^1\text{H}$  spin-lattice relaxation times ( $T_1$ ,  $\mu\text{W}$  off) and DNP build-up times ( $T_B$ ,  $\mu\text{W}$  on) as a function of composition parameter  $x$  for the investigated samples. Values were obtained from unconstrained least-square fittings of saturation-recovery curves by using stretched exponential functions.



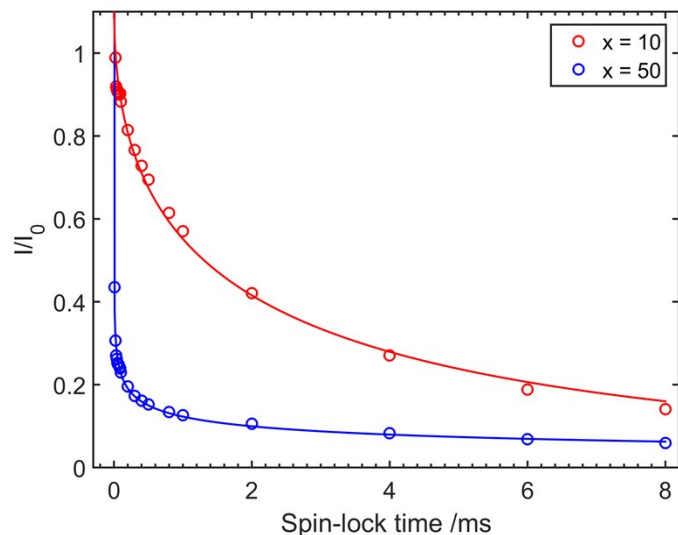
**Figure S2** – Experimental (black curve, bottom) and simulated (red curve, top) cw EPR spectra for the sample with  $x = 10$  mol%. The experimental spectrum was measured at 100 K. The simulation was performed using the EasySpin<sup>1</sup> package for the Matlab<sup>®</sup> software. The simulation uses the “pepper” function from EasySpin,<sup>1</sup> which considers a solid state powder. The system consists of an electron spin with rhombic  $g$  and  $^{14}\text{N}$ -hyperfine tensors ( $^{15}\text{N}$  was also considered in natural abundance). The parameters were optimized by comparing experimental and simulated spectra, using a combination of least-square fitting and genetic algorithms and taking the parameters reported by Dhanasekaran et al.<sup>2</sup> as initial guess. The final parameters for the simulation are  $g_{xx} = 2.0088$ ,  $g_{yy} = 2.0057$ ,  $g_{zz} = 2.0021$ ,  $A_{xx} = 12 \pm 1$  MHz and  $A_{yy} = 14 \pm 1$  MHz and  $A_{zz} = 102 \pm 1$  MHz. The simulation also considers a Gauss-Lorentz (0.6-0.4) isotropic line broadening of 28 MHz (peak-to-peak).



**Figure S3** –  $^{29}\text{Si}$  MAS NMR spectra for the investigated samples ( $x = 0.5$  to 50) and for a sample without carboxy-PROXYL radical ( $x = 0$ ). The spectral line width increases for increasing radical concentration due to electron-nucleus hyperfine interaction.



**Figure S4** – Top:  $^1\text{H}$ – $^1\text{H}$  MAS FSLG spectrum of a) mesoporous SBA-15 functionalized with APTES and Fmoc-Glycine (diamagnetic sample) and b) mesoporous SBA-15 functionalized with APTES, Fmoc-Glycine and Carboxy-PROXYL radicals (paramagnetic sample,  $x = 50$ ). Skyline projections are also displayed. Bottom: comparison of the projections taken from the sum over the indirectly detected ( $\delta_1$ ) dimension for both spectra shown on top. 2D spectra were measured with a z-filtered  $^1\text{H}$ – $^1\text{H}$  correlation pulse sequence with Frequency Switched Lee–Goldberg decoupling during  $t_1$  evolution.<sup>3</sup> The experiments were performed using  $^1\text{H}$  carrier frequency of 400 MHz, spinning rate of 10 kHz and z-filter delay of 1  $\mu\text{s}$ . In this experiment a MAS spectrum is detected in the direct domain ( $\delta_2$ ) while FSLG decoupled signal is achieved in the indirect domain ( $\delta_1$ ). Due to the FSLG decoupling the  $^1\text{H}$ – $^1\text{H}$  dipolar coupling is averaged out and the line broadening caused by  $^1\text{H}$  spin-spin relaxation mechanism is suppressed. The result is a resolution enhancement in the  $\delta_1$  projection and the detection of the  $\delta_2$  projections similar to the 1D spectrum. On the other hand, the electron-nucleus hyperfine interaction is not averaged out by FSLG decoupling leading to fast relaxation of areas containing the radicals. As a result, only components outside of the effective paramagnetic shell will survive seen as relative enhancement of the peak at 1.2 ppm in the projections, accompanied by a higher resolved  $\delta_1$  projections. This peak belongs most probably to mobile and isolated water molecules on the surface,<sup>4</sup> far away from the radicals.



**Figure S5** –  $^1\text{H}$  magnetization as a function of the spin-lock time for samples  $x = 10$  and  $x = 50$ , measured using a  $90^\circ$ -spinlock pulse sequence, nuclei are excited by a  $90^\circ$  pulse and then spin-locked by a continuous wave pulse (with a phase shift of  $\pi/2$  relative to the excitation pulse) for a time  $\tau$ .<sup>5</sup> The experimental points are normalized to the extrapolated equilibrium magnetization. The evolution of the magnetization curves gives information about the  $^1\text{H}$  spin lattice relaxation in the rotating frame ( $T_1^\rho$ ). Solid curves are exponential fit to the data, using a stretched exponential function of the form  $I = I_0 \exp\left(-\left(t/T_1^\rho\right)^\beta\right)$  (see parameter values in Table S5). Due to the distribution of distances between H species and the unpaired electrons in the sample, a continuous distribution of relaxation times is expected. Clearly, for the sample with  $x = 50$  the effective  $^1\text{H}$  spin relaxation is much faster than for the sample with  $x = 10$ . For this reason,  $^1\text{H}$ - $^{13}\text{C}$  CP transfer could not be achieved for sample  $x = 50$  (see main text).

**Table S5:** Fitting parameters to a stretched exponential decay function,  $I = I_0 \exp\left(-\left(t/T_1^\rho\right)^\beta\right)$ , adjusted on the data shown in Figure S5.

Sample	$T_1^\rho$ (ms)	$\beta$
$x = 10$	2.1	0.5
$x = 50$	$2 \times 10^{-3}$	0.1

## References

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