

Electronic Supporting Information

**Photo-isomerization of the isolated photoactive yellow protein
chromophore: What comes before the primary step?**

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Experimental Details

Time-resolved photoelectron imaging of mass-selected *para*-coumaric ketone (*p*CK) was performed using an instrument that has been described previously.¹⁻³ Deprotonated anions of *p*CK, *p*CK⁻, were produced by negative-mode electrospray ionisation of a ~1 mM solution of *p*CK in methanol. Ions were transferred into vacuum through a stainless steel capillary and guided through 3 differentially pumped vacuum regions using a series of ring-electrode RF ion guides. Ions were trapped at the end of these guides and extracted collinearly into a time-of-flight mass-spectrometer at 100 Hz. At the focus of the mass-spectrometer (1.3 m from extraction), the mass-selected ion packet was irradiated with femtosecond pulses at the centre of a perpendicular velocity-map imaging photoelectron spectrometer.² Emitted electrons were detected in the spectrometer using a position sensitive detector (dual multichannel plate coupled to phosphor screen with output monitored using a CMOS camera). Images were analysed using polar onion peeling⁴ to yield the central slice through the 3D electron distribution, which offered the photoelectron spectrum and photoelectron angular distributions. The spectra were calibrated using the well-known photoelectron spectrum of iodide and the spectral resolution as ~5% of the electron kinetic energy (eKE). The raw data is available from at <https://doi.org/10.5281/zenodo.3750967>

Femtosecond pulses were derived from a commercial femtosecond laser system (2.5 W, 800 nm, 35 fs pulses, 1 kHz). Part of the output was used directly as probe pulses (~450 μJ pulse⁻¹). Pump pulses were produced by 4th harmonic generation of the idler output of an optical parametric amplifier, yielding ~5 μJ pulse⁻¹ at 444 nm. Probe pulses were delayed relative to the pump pulses using a motorized optical delay line and recombined collinearly using a dichroic mirror before being directed into the instrument unfocussed (beam diameter ~ 3 mm). Pump and probe polarizations were set parallel to the detector.

Computational Details

Energetic minima corresponding to the critical points on the $p\text{CK}^- \text{S}_0$ and S_1 potential energy surfaces were optimized using $\omega\text{B97X-D/aug-cc-pVDZ}$.⁵ Specifically, Density Functional Theory (DFT) was used to optimize the Franck-Condon minimum of the S_0 state, $\text{FC}(\text{S}_0)$, while linear response time-dependent DFT (LR-TDDFT) with the Tamm-Danoff Approximation (TDA)⁶ was used to optimize the S_1 planar minimum, $\text{S}_{1,\text{PM}}$. The nature of each stationary point located was confirmed by frequency calculations. All optimization and frequency calculations were carried out using Gaussian 09 (rev D.01).⁷

The vertical excitation energies of the S_0 , S_1 and D_0 states were recalculated using multi-state extended multi-configurational quasi-degenerate perturbation theory (MS-XMCQDPT2).⁸ An (aug)-cc-pVTZ basis set⁹ was used, in which the augmented function was affixed to only the oxygen atoms. Inclusion of the D_0 state in the active space (and therefore detachment to the D_0 plus free electron continuum was) was achieved through addition of an orbital characterised by an extremely diffuse p-function ($\alpha = 1\text{E}-10$), transitions to this state were included in the 6-state averaging procedure to mimic detachment to the continuum. Following previously established protocols,¹⁰⁻¹² a DFT/PBE0-based one-electron Fock-type matrix was used to obtain energies of the MCSCF semicanonical orbitals used in perturbation theory. All MS-XMCQDPT calculations were carried out using the Firefly quantum chemistry package.¹³

Dyson Orbitals

The Dyson orbitals were calculated using the equation-of-motion coupled cluster singles and doubles method for electronically excited and ionisation potential energies (EOM-EE/IP-CCSD),^{14,15} using the Pople 6-31+G** basis set.^{16,17} The Dyson orbitals were

then used to model the trends in PADs as a function of changing energy of the photodetached electron. The EOM-EE/IP-CCSD calculations were carried out using the QChem 5.2 computational package¹⁸ and the PADs were modelled using ezDyson v4.¹⁹



Figure S1: Dyson orbitals for S₁ in FC geometry (left) and S_{1,PM} geometry (right).

Energies are calculated for the EOM-EE/IP-CCSD Dyson orbital calculation but, given the reduced basis set used, these are of lower accuracy as those calculated with the XMCQDPT2 level of theory. The EOM-EE/IP-CCSD energies were calculated to be 2.8657/2.6407 eV and 2.8041/2.5220 eV, for the FC and S_{1,PM} geometry, respectively.

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