# Metal-to-metal communication during the spin state transition of a [2x2] Fe (II) metallogrid at equilibrium and out-of-equilibrium conditions 

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## Synthesis

The complex grid FE4 and its ligand were synthesized following the procedure reported in literature ${ }^{1}$. The crystallographic data for this paper can be obtained free of charge from The Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/data request/cif. Deposition numbers:

FE4 from 100 K to $390 \mathrm{~K}: 2126240$, 2127042-2127051, 2127052-2127057 and 2127059-2127064.
Time resolved: 2127311-2127319.

## Vibrational spectroscopy

The infrared spectra in the $2000-2500 \mathrm{~cm}^{-1}$ range of the original and desolvated $\mathbf{F E 4}$ grid are shown in Figure S6. The presence of acetonitrile in the original FE4 sample is confirmed by the appearance of two distinctive signals: $v(\mathrm{C} \equiv \mathrm{N})$ at $2249 \mathrm{~cm}^{-1}$ and a combined band of $v(\mathrm{C}-\mathrm{C})+\delta(\mathrm{CH} 3)$ at $2293 \mathrm{~cm}^{-1}{ }^{2-5}$. In contrast, the desolvated FE4 grid did not display the aforementioned bands, which indicate a complete desolvation.


Figure S1. FTIR spectra for original and desolvated FE4 grid.

## Characterization

Table S1. Crystal data and structure refinement parameters of FE4 at 100 K .

| Temperature | 100 K |
| :---: | :---: |
| Empirical Formula | C100 H74 B4 F16 Fe4 N26 |
| Crystal color/habit | Black/block |
| Crystal size (mm) | (0.090x0.120x0.050) |
| Crystallizing solvlent | Acetonitrile |
| Crystal system/ Space group | Monoclinic/ C2/c |
| $a$ (Å) | 30.680 (6) |
| $b$ (Å) | 12.780 (3) |
| $c$ (Å) | 26.810 (5) |
| $\alpha\left({ }^{\circ}\right)$ | 90 |
| $\beta\left({ }^{\circ}\right)$ | 120.76 (3) |
| $Y\left({ }^{\circ}\right)$ | 90 |
| Volume ( ${ }^{\text {² }}$ ) | 9107.0 (4) |
| Z/Z' | 4/0.5 |
| Molecular Weight | 2210.49 |
| Calculated density (g/cm3) | 1.625 |
| F(000) | 4496 |
| Radiation | Synchrotron ( $\lambda=0.61990$ A $)$ |
| $\theta$ range ( ${ }^{\circ}$ ) | 1.902/31.355 |
| Scan type | $\varphi$ |
| Measured reflections | 48052 |
| Unique reflections | 14709 |
| Observed reflections $[\|F\|>4 \sigma(F)]$ | 13786 |
| Final R (\%) | 5.03 |
| wR2 (\%) | 14.88 |
| Good-of-fit on $\mathrm{F}^{2}(\mathrm{~S})$ | 1.086 |
| $\Delta \rho \max \left(\mathrm{e} . \AA^{-3}\right)$ | 0.794 |
| $\Delta \rho \min \left(\mathrm{e} . \AA^{\circ}{ }^{-3}\right)$ | -0.696 |
| No of restrains/parameters | 90/751 |
| Data [\|F|>4 ${ }^{\text {c }}$ (F)]-to-parameter ratio | 18.35:1 |

Table S2. Selected bond lengths and bond angles for compound FE4 at 100K.

| Bond lengths ( $\AA$ ) |  | Bond angles ( ${ }^{\circ}$ ) |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Fe(A)-N2 | 1.9030 (12) | N2-Fe(A)-N5 | 172.81(5) | N7-Fe(B)-N10 | 100.60(5) |
| $\mathrm{Fe}(\mathrm{A})$ - N 5 | $1.9036(12)$ | $\mathrm{N} 2-\mathrm{Fe}(\mathrm{A})-\mathrm{N} 00$ | 80.53(6) | $\mathrm{N} 7-\mathrm{Fe}(\mathrm{B})-\mathrm{N} 11$ | 129.94(5) |
| $\mathrm{Fe}(\mathrm{A})$-N00 | $1.9893(14)$ | $\mathrm{N} 5-\mathrm{Fe}(\mathrm{A})$ - N 00 | 95.49(6) | N10-Fe(B)-N11 | 76.23(5) |
| $\mathrm{Fe}(\mathrm{A})$-N3 | $1.9898(13)$ | N2-Fe(A)-N3 | 80.08(6) | N7-Fe(B)-N8 | 76.59(6) |
| $\mathrm{Fe}(\mathrm{A})$-N4 | $1.9949(12)$ | $\mathrm{N} 5-\mathrm{Fe}(\mathrm{A})$ - N 3 | 103.96(6) | N10-Fe(B)-N8 | 128.36(5) |
| Fe(A)-N6 | 1.9953(11) | $\mathrm{N} 00-\mathrm{Fe}(\mathrm{A})-\mathrm{N} 3$ | 160.55(5) | N11-Fe(B)-N8 | 143.82(6) |
| $\mathrm{Fe}(\mathrm{B})$-N7 | $2.1203(14)$ | $\mathrm{N} 2-\mathrm{Fe}(\mathrm{A})$ - N 4 | 94.07(5) | N7-Fe(B)-N12 | 90.50(5) |
| $\mathrm{Fe}(\mathrm{B})$-N10 | 2.1327 (13) | $\mathrm{N} 5-\mathrm{Fe}(\mathrm{A})-\mathrm{N} 4$ | 80.37(5) | N10-Fe(B)-N12 | 146.11(5) |
| $\mathrm{Fe}(\mathrm{B})$-N11 | $2.1331(13)$ | $\mathrm{N} 00-\mathrm{Fe}(\mathrm{A})-\mathrm{N} 4$ | 96.60(5) | N11-Fe(B)-N12 | 72.20(5) |
| $\mathrm{Fe}(\mathrm{B})$-N8 | $2.1413(13)$ | N3-Fe(A)-N4 | 86.31(5) | $\mathrm{N} 8-\mathrm{Fe}(\mathrm{B})-\mathrm{N} 12$ | 85.27(5) |
| $\mathrm{Fe}(\mathrm{B})$-N12 | $2.2799(14)$ | $\mathrm{N} 2-\mathrm{Fe}(\mathrm{A})-\mathrm{N} 6$ | 105.09(5) | N7-Fe(B)-N9 | 147.01(5) |
| $\mathrm{Fe}(\mathrm{B})$-N9 | $2.3278(17)$ | N5-Fe(A)-N6 | 80.30(5) | N10-Fe(B)-N9 | 92.69(6) |
|  |  | N00-Fe(A)-N6 | 84.10(5) | N11-Fe(B)-N9 | 82.44(5) |
|  |  | N3-Fe(A)-N6 | 99.48(5) | N8-Fe(B)-N9 | 71.53(6) |
|  |  | N4-Fe(A)-N6 | 160.64(5) | N12-Fe(B)-N9 | 95.07(6) |



Figure S2. Asymmetric unit of the FE4 grid at 100K. Hydrogen were omitted for clarity.



Figure S3. Thermal ellipsoid plot ( $50 \%$ of probability) for FE4 at 100K.


Figure S4. Packing of molecules down crystallographic ' $a$ ' axis. Hydrogen atoms are omitted for clarity.


Figure S5. Packing of molecules down crystallographic ' $b$ ' axis. Hydrogen atoms are omitted for clarity.


Figure S6. Packing of molecules down crystallographic ' $c$ ' axis. Hydrogen atoms are omitted for clarity.


Figure S7. Thermal ellipsoid plot (50\% of probability) for FE4 for the -5 ns data.


Figure S8. Thermal ellipsoid plot (50\% of probability) for FE4 for the -2 ns data.


Figure S9. Thermal ellipsoid plot (50\% of probability) for FE4 for the -200ps data.


Figure S10. Thermal ellipsoid plot (50\% of probability) for FE4 for the 100ps data.


Figure S11. Thermal ellipsoid plot (50\% of probability) for FE4 for the 200ps data.


Figure S12. Thermal ellipsoid plot (50\% of probability) for FE4 for the 500ps data.


Figure S13. Thermal ellipsoid plot (50\% of probability) for FE4 for the 1 ns data.


Figure S14. Thermal ellipsoid plot (50\% of probability) for FE4 for the 2 ns data.


Figure S15. Thermal ellipsoid plot (50\% of probability) for FE4 for the 5 ns data.

## Structural details

Table S3. Structural data that provides a basic description of the SCO in FE4

 $\zeta=\left(\mathrm{Fe}-\mathrm{N}_{\mathrm{i}}\right)-<\mathrm{Fe}-\mathrm{N}>9$. Information about $\mathrm{S}(\mathrm{Oh})$ and $\mathrm{S}(\mathrm{itp})$ is provided in the structural analysis section.

## Structural Analysis

The Octadist program ${ }^{10}$ was used to determine the $\left.<\mathrm{Fe}-\mathrm{N}\right\rangle$ bond length and the angular distortion parameter that describe the octahedral coordination environment of the metal centres in the FE4 grid. The angular distortion parameter, $\Theta$, is the sum of the deviations from $60^{\circ}$ of the twenty-four $\mathrm{N}-\mathrm{Fe}-\mathrm{N}$ angles, six per pseudo three-fold axis, measured on a projection of opposite triangular faces of the $\left\{\mathrm{FeN}_{6}\right\}$ octahedron, orientated by superimposing the face centroids (Figure S16). ${ }^{8,11}$


Figure S16. Environment of $\mathrm{Fe}^{\mathrm{II}}$ ions and definition of the $\theta$ angle and the angular distortion parameter $(\Theta) .{ }^{8,11}$

For comparison, continuous shape measurements (CShM) were also used to characterise the relative deviation of the metal coordination spheres in FE4 from ideal polyhedra described by a particular point symmetry group (Table S2). ${ }^{12}$ Mathematically, CShM of the coordination polyhedron Q with the geometric centre $\vec{q}_{0}$ relative to an ideal polyhedron P is expressed as:

$$
\begin{equation*}
S_{Q}=\min \left[\frac{\sum_{i=1}^{N}\left|\vec{q}_{i}-\vec{p}_{i}\right|^{2}}{\sum_{i=1}^{N}\left|\vec{q}_{i}-\vec{q}_{0}\right|^{2}}\right] \times 100 \tag{S1}
\end{equation*}
$$

where $\vec{q}_{i}$ and $\vec{p}_{i}$ are the position vectors for atoms of two polyhedral. CShM relative to an ideal octahedron $\left(\mathrm{S}(\mathrm{Oh})\right.$ ) and an ideal trigonal prism ( $\mathrm{S}(\mathrm{itp})$ ) were calculated using the SHAPE program ${ }^{13}$. The calculation of $S(O h)$ and $S(i t p)$ were performed for all crystallographic-symmetry independent metal atoms.

It is well-known that the $\left\{\mathrm{FeN}_{6}\right\}$ coordination sphere of $\mathrm{LS} \mathrm{Fe}^{\text {II }}$ ions is more regular, i.e., closer to an ideal octahedron. Therefore, the $\mathrm{S}(\mathrm{Oh})$ parameter is small and closer to 0 , while the parameter $\mathrm{S}(\mathrm{itp}) \gg$ 0 . Contrary, $\mathrm{Fe}^{\mathrm{II}}$ ions in the HS state are characterised by a more irregular structure with structural parameters $\mathrm{S}(\mathrm{Oh}) \gg 0$ and S (itp) closer to zero.

## Temperature difference

The well-known Wilson plot ${ }^{14}$ is frequently used to estimate the scale factor, $\mathrm{k}_{\mathrm{s}}$, and the overall isotropic temperature factor (B) of a data set. This plot can be obtained by statistical comparison of the observed intensities ( $\mathrm{I}_{\mathrm{obs}}$ ) with the average squared structure factor equals $\sum_{i=1}^{M} f_{i}^{2}$ for each $\sin \theta / \lambda$ range according to:

$$
\begin{equation*}
\ln \left(\frac{I_{o b s}}{\sum_{i=1}^{M} f_{i}^{2}}\right)=\ln \left(k_{s}\right)-2 B(\sin \theta / \lambda)^{2} \tag{S2}
\end{equation*}
$$

Where M is the number of atoms and $f_{\mathrm{i}}$ is the scattering factor of the ith atom.
The change in the overall isotropic temperature factor $(\Delta \mathrm{B})$ between two data sets at different temperature can be estimated from modified Wilson plots. The plots are obtained by a scale-factor refinement of the low temperature data (e.g. 100K) with the high-temperature data (e.g. 290K) structural model and plotting the $\ln \left(\mathrm{I}^{100 \mathrm{~K}} / \mathrm{I}^{290 \mathrm{~K}}\right)$. The slope of the dependence of $\ln \left(\mathrm{I}^{100 \mathrm{~K}} / \mathrm{I}^{290 \mathrm{~K}}\right)$ with $(\sin \theta / \lambda)^{2}$ gives the overall increase of isotropic atomic motion, $\Delta B$ (equation $S 3$ ), which is associated with temperature difference between the data sets. ${ }^{15}$

$$
\begin{equation*}
\ln \left(\frac{I^{100 K}}{I^{290 K}}\right)=-2 \Delta B^{100 K-290 K}(\sin \theta / \lambda)^{2} \tag{S3}
\end{equation*}
$$

The greater the temperature difference between data sets the greater the value of $\Delta \mathrm{B}$, that means a more negative slope of the plot. This feature is used to analyse the temperature increase during the photocrystallographic experiments as explain below.

During the photo-crystallographic experiments, the energy deposited by the laser pulse largely exceed the energy necessary for the LS to HS transition, which results in some heat diffusion and global warming. A modified Wilson plot, known as photo-Wilson plot (Fig. S17), is then used to estimate the lase-induced temperature increase due to heat dissipation, in a similar way as described above for the temperature-Wilson plots. From the photo-Wilson plot is possible to calculate the variation of the isotropic temperature factor between the data sets at time " t " and the reference data at time $\mathrm{t}<0\left(\Delta \mathrm{~B}^{\mathrm{t}}\right.$ $\left.{ }^{\ll 0}\right)^{15}$ :

$$
\begin{equation*}
\ln \left(\frac{I^{d t}}{I^{d t<0}}\right)=-2 \Delta B^{d t-d t<0}(\sin \theta / \lambda)^{2} \tag{S4}
\end{equation*}
$$

where, $\mathrm{I}^{\mathrm{dt}}$ and $\mathrm{I}^{\mathrm{dt}<0}$ are the observed intensities of data sets at delay time "dt" and the reference data at delay time $\mathrm{dt}<0$. Note that both the thermal-Wilson plot and the photo-Wilson plot were brought to the same scale before calculating the value of $\Delta \mathrm{B}$. For more information about this topic, the reader is referred to the literature ${ }^{15}$.
a)

c)

e)

b)

d)

f)


Figure S17. PhotoWilson plots of FE4

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