9 * Corresponding author

13 As shown in Fig. S1(a), ${ }^{d i v}$ is defined as the length of the fringe produced by a single sawtooth

$$
\begin{equation*}
{ }_{16} L_{d i v}=\frac{h_{s t}}{\sin \theta_{s t}} \tag{S1}
\end{equation*}
$$

18 Fig. $\mathrm{S} 1(\mathrm{~b})$ is the expanded view of Fig. S1(a). The distance between the two fringes is $k \lambda_{\theta}-h_{\theta}$, 19 where the blue line represents the resulting divergent fringe. The angle at which the divergent

## Supplementary Information

## Microfluidic acoustic sawtooth metasurfaces for patterning

 and separation using traveling surface acoustic wavesMingxin Xu, Peter VS Lee, David Collins*

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## Supplementary Note 1: Analytical formal and derivation for non-integer $k$ values

 element (equals to the length of the hypotenuse of each triangle element), written as fringe increases relative to $\theta_{\text {st }}$ is defined as ${ }^{\text {inc }}$, which is calculated as:21

27 and (S2) with $\theta_{\text {div }}^{+}$, this becomes
28
$\theta_{\text {div }}^{+}=\theta_{s t}+\theta_{i n c}^{+}=\theta_{s t}+\tan ^{-1}\left(\frac{\left(k \lambda_{\theta}-h_{\theta}\right) \sin \theta_{s t}}{h_{s t}}\right)$
30

31 Conversely, for $\theta_{d i v}$ slightly smaller than $\theta_{\text {st }}$, this expression incorporating $\theta^{-}{ }^{-}$(superscript
32 " - " for $\theta_{\text {div }}$ is slightly smaller than $\theta_{\text {st }}$ ) is expressed as:
33

34



42 FIG. S2. Analytical modelling results, plotting $\theta_{\text {div }}$ as a function of $\theta_{\text {st }}$ and $\sigma$, overlaid with red
43 lines corresponding to $k=1, k=2$ and $k=3$.
${ }_{50} \frac{1}{\lambda_{\theta}} \sin \left(\theta_{t}\right)=\frac{1 d \varphi(y)}{2 \pi d y}$,

$$
\begin{equation*}
58 d y=\frac{h_{s t}}{k \tan \left(\theta_{s t}\right)} \tag{S6}
\end{equation*}
$$

60 Substituting equation (S6) into equation (S5) with $d \varphi(y)=2 \pi$, and rewriting equation (S5)
${ }_{63} \frac{1}{\lambda_{\theta}} \sin \left(\theta_{d i v}\right)=\frac{k \tan \left(\theta_{s t}\right)}{h_{s t}}$,
64 (equal to $\theta_{d i v}$ ) as shown in Fig. S3. $d y$ is the length along the y axis for which the equivalent phase shift $(d \varphi(y))$ is equal to $2 \pi$.

As shown in Fig. S3, $d y=w_{s t} / k$, where $w_{s t}=\mathrm{h}_{s t} / \tan \left(\theta_{s t}\right)$. Therefore, $d y$ is expressed as: yields the generalized Snell's law for our sawtooth metasurfaces:

## Supplementary Note 2: Deriving generalized Snell's law for the proposed metasurfaces.

For acoustic waves that incident perpendicular to the metasurface (the incidence angle is 0 ), the generalized Snell's law of refraction is written as [1]:

Where the fringe spacing $\lambda_{\theta}$ is described in equation (1), $\theta_{t}$ is the equivalent refraction angle


66 FIG. S3. Schematic diagram of deriving generalized Snell's law. For $k=1$ and $k=2$.


70 71


0.0003 W

0.01 W

0.001 W

0.03 W

0.003 W

0.1 W

85 FIG. S7. Input power effects. Experimental results with $\lambda_{S A W}=100 \mu m$ and $k=1$, for the input power of the transducer is (a) 0.0003 W , (b) 0.001 W , (c) 0.003 W , (d) 0.01 W , (e) 0.03 W , and 87 (f) 0.1 W . Scale bars is $200 \mu \mathrm{~m}$.

88


89

90 FIG. S8. Experimental results demonstrating filling without bubble formation at low perfusion 91 flow rates ( $\sim 7 \mu \mathrm{~L} / \mathrm{min}$ ).


92

## 98

FIG. S9. Examples and acoustic pressure field simulations of sawtooth-like metasurfaces and other metasurfaces with $\lambda_{S A W}=100 \mu m$. (a)-(e) Sawtooth-like metasurfaces. The pattern generated by each sawtooth-like element has the same length at the dotted line to form spatially continuous patterns, where the spacing of the pattern at a given ${ }{ }_{\text {st }}$ can be calculated by equation (1). (f)-(i) Other types of metasurfaces. Scale bar is $100 \mu m$.

## Reference

1. Tian Z, Shen C, Li J, Reit E, Gu Y, Fu H, et al. Programmable Acoustic Metasurfaces. Adv Funct Mater. 2019 Mar;29(13):1808489.
