# Supporting Information

# P-type molecular doping by charge transfer in halide perovskite

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MR (%)	Thickness (nm)
0	$500 \pm 100$
0.1	$750 \pm 300$
0.2	$450 \pm 100$
1	$500 \pm 100$
5	$400 \pm 100$

Fig. S1 Selected SEM cross-section images of MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films with (a) 0%, (b) 0.1%, (c) 0.2%, (d) 1% and (e) 5% MR (molar ratio) F4TCNQ deposited on glass. (f) Averaged thickness of the layers over 6 to 12 different measurements along each film. The scale bar is 3  $\mu$ m.



**Fig. S2** Time evolution of the conductivity  $\sigma$ , Hall density  $n_{Hall}$  and mobility  $\mu_{Hall}$  extracted in the dark and under illumination (grey region) for (a) undoped and (b) 1% MR (molar ratio) doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films. Equilibrium in the dark is reached after waiting approximately one hour. Light excitation is performed with a red laser (638 nm) with a light intensity of 0.04 W/cm<sup>2</sup>.



**Fig. S3** SEM images transformed into black and white images for pristine 0%, 0.1%, 0.2%, 1% and 5% MR (molar ratio) F4TCNQ-doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films. The table summarizes the proportion of white (perovskite) over total area *C<sub>white</sub>*.

### **Percolation threshold**

To determine if the percolation threshold is reached, we calculate the fraction of conducting (perovskite) over total area from SEM images and compare this value to the percolation threshold for an appropriate 2D system. The SEM image is transformed into a black and white image and the ratio of white over total pixels ( $C_{white}$ ) is calculated using ImageJ. The  $C_{white}$  values are summarized in Figure S3.

Our perovskite layers can be considered as nominally 2D systems, as our thin film cross-section shows that the entire film thickness encompasses one single grain (see Figure S1).<sup>1</sup> For 2D systems, we can expect the percolation threshold to be lower than 0.625, as grains are surrounded by concave pores.<sup>2,3</sup> The fractions of conducting/total area determined from SEM images are above this limit. Therefore, we can conclude that our samples should be above the percolation threshold, a statement confirmed by a conductivity significantly higher than the background glass conductivity expected if the percolation threshold is not met.

#### Impact of percolating network on Hall measurements

While the Hall coefficient is expected to increase as we get closer to the percolation threshold for 3D systems, it has been shown that the Hall coefficient remains constant even when close to the percolation threshold in 2D networks.<sup>4,5</sup> Therefore, as long as the percolation threshold is reached, the measured Hall coefficient should correspond to the Hall coefficient in the conducting medium only. This implies that Hall mobilities should not depend on the percolating network. Nevertheless, we expect the Hall mobility to vary with grain boundary density if grain boundaries are not benign.



Fig. S4 SEM images of  $MAPb_{0.5}Sn_{0.5}I_3$  with 1% (a ; c) and 5% (b ; d) MR (molar ratio) F4TCNQ dopant, using Trinity detectors with a 5000x magnification and collecting mostly secondary electrons (a ; b) and mostly backscattered electrons (c ; d).



**Fig. S5** TOF-SIMS 2D images of 1% MR (molar ratio) doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>J<sub>3</sub> (different sample from main text Fig. 3) exhibiting the total signal (a) and the F signal (b). The color scale for each image shows the intensity in counts/pixel. TOF-SIMS 3D tomography of the same sample showing the signal coming from MA (MA: methylammonium) and indicating the presence of the perovskite (c) and the F signal (d). Each 3D reconstruction is  $50 \times 50 \times 50 \times film$  thickness  $\mu m^3$ .



**Fig. S6** TOF-SIMS 2D images of undoped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> exhibiting the total signal (a) and the F signal (b). The color scale for each image shows the intensity in counts/pixel. TOF-SIMS 3D tomography of the same sample showing the signal coming from MA (MA: methylammonium) and indicating the presence of the perovskite (c) and the F signal (d). Each 3D reconstruction is  $50 \times 50 \times$  film thickness  $\mu$ m<sup>3</sup>. We observe the background signal for F.



**Fig. S7** (a) Current-voltage characteristics of undoped and 1% MR (molar ratio) doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>]<sub>3</sub> films processed under different conditions: with toluene as antisolvent (AS), with chlorobenzene (CB) as antisolvent and without antisolvent. SEM images of MAPb<sub>0.5</sub>Sn<sub>0.5</sub>]<sub>3</sub> films (b) undoped with toluene as antisolvent, (c) 1% MR doped with toluene as antisolvent and (d) 1% MR doped with CB as antisolvent. The scale bar is 4  $\mu$ m.



Fig. S8 Current at 5V measured over several days on encapsulated Hall bars for  $MAPb_{0.5}Sn_{0.5}I_3$  films with 0%, 1% and 5% MR (molar ratio) dopant concentration.



Fig. S9 XRD patterns of MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films with 0%, 0.1%, 0.2%, 1% and 5% MR (molar ratio) F4TCNQ. The Miller indices of the major peaks and associated with the tetragonal I4/mcm space group structure are indicated.



**Fig. S10** EDX spectra obtained on MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films with 0%, 0.1%, 0.2%, 1% and 5% F4TCNQ. The measurements are performed on three different spots and the average stoichiometry for Sn, I and F with respect to Pb are summarized in the table. We note that for 5% MR (molar ratio) F4TCNQ, we would expect a value of 0.2 for F, close to the value obtained. Nevertheless, the signal is weak for F in EDS spectra and it is difficult to quantify F content reliably.



**Fig. S11** Normalized absorption (a) and Tauc plot for direct bandgap (b) for MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films with 0%, 0.1%, 0.2%, 1% and 5% MR (molar ratio) F4TCNQ. The inset table in b) summarizes the bandgap ( $E_G$ ) extracted for each sample.

Table S1	Summary	of the	WF	values	obtained	for	undoped	and	1%	MR	(molar	ratio)	doped	MAPb <sub>0.5</sub> Sn <sub>0.5</sub> I <sub>3</sub>	using	UPS	and	KP-CPD
measurem	ents. High	ly order	ed p	yrolytic	graphite	(HOF	PG) is used	as re	efere	ence	to calib	rate th	e WF of	the Kelvin pro	be tip.			

Samples	WF, <sub>UPS</sub> (eV)	WF, <sub>KP</sub> (eV)
Undoped	4.51	4.48
1% doped	4.73	4.85



**Fig. S12** Hole mobility of MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> films for different F4TCNQ concentration in molar ratio (MR) obtained from AC Hall effect measurements. Several samples have been measured for each dopant concentration and the box charts indicate the maximum and minimum values (whiskers), the upper quartile and lower quartile (box) and the mean value (empty square).



**Fig. S13** Hall signal extraction on undoped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> using a PDL setup: a) Magnetic field B traces as the reference signal, and b) Fourier transform of the reference signal B. c) Transverse Hall signal R<sub>XY</sub> and d) Fourier transform of R<sub>XY</sub>. The dashed dotted lines correspond to the second and third harmonics of the AC signal B. The red filled region is the power spectral density (PSD) and the curve is the Fourier spectra of the signal. The former accentuates the periodic signal content as shown in the fundamental first harmonic component of the signal. e) Lock-in detection of the in-phase X (Hall signal) and out-of-phase Y signals over time where the Hall resistance R<sub>H</sub> is extracted. A lock-in time constant of 100s is used in this example. The single harmonic ac magnetic field along with lock-in detection allow the extraction of an R<sub>H</sub> value  $\sim 10^5 \Omega$  with a background resistance of  $\sim 10^9 \Omega$ .



Fig. S14 FTIR absorption spectra for pristine MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> and pure F4TCNQ deposited from CB solution on KBr substrates.



**Fig. S15** Multiple Gaussian peak fitting performed on the FTIR spectra of MAPb<sub>1-x</sub>Sn<sub>x</sub>l<sub>3</sub> with Pb/Sn ratio of 50/50 (x=0.5) (a), 75/25 (x=0.25) (b) and 100/0 (x=0) (c) doped with 5% MR (molar ratio) F4TCNQ. The center and area of each peak is given in the table (d).



**Fig. S16** FTIR absorption spectra for 5% MR (molar ratio) doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> over time. The first measurement (2 min) is performed 2 min after opening the sealed bag containing the sample (sealed in nitrogen environment). Each scan takes about 4 min to complete. The spectra are normalized at 2220 cm<sup>-1</sup>. It is important to note that the wavenumber chosen for normalization influences the observation of peak intensity evolution. Nevertheless, it is clear the F4TCNQ<sup>2</sup> contribution is decreasing over time.



**Fig. S17** XRD patterns of 5% MR (molar ratio) doped MAPb<sub>1-x</sub>Sn<sub>x</sub>I<sub>3</sub> with Pb/Sn ratio of 50/50 (x=0.5), 75/25 (x=0.25) and 100/0 (x=0) (a). EDX spectra obtained on MAPb<sub>0.75</sub>Sn<sub>0.25</sub>I<sub>3</sub> (b) and MAPbI<sub>3</sub> (c). The measurements are performed on three different spots and the average stoichiometry for Sn and I with respect to Pb are summarized in the table (d).



**Fig. S18** Picture of the MAPb<sub>x</sub>Sn<sub>1-x</sub>I<sub>3</sub> precursor solutions taken immediately after filtration: undoped Pb/Sn 50/50 (a), 5% MR (molar ratio) doped Pb/Sn 50/50 (b), 5% doped Pb/Sn 75/25 (c) and 5% doped Pb/Sn 100/0 (d). Picture of MAI:F4TCNQ solution in DMF:DMSO (80:20 in volume) (e). Of note, the color does not change significantly with time over a period of a few weeks when stored in nitrogen.



**Fig. S19** Carrier-resolved photo Hall H- $\sigma$  plot with an excitation wavelength of 638 nm for undoped (a) and 1% MR (molar ratio) doped (b) MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub>. A maximum light intensity  $I_{max} = 30 \ mW cm^{-2}$  is used.

#### <u>CRPH analysis with $\Delta p \geq \Delta n$ </u>

In the case where  $\Delta p \neq \Delta n$  ( $\Delta p$  and  $\Delta n$  being the photogenerated hole and electron densities), the conductivity for a p-type semiconductor under illumination should be written as:

$$\sigma = q(\mu_p(p_0 + \Delta p) + \mu_n \Delta n), \tag{S1}$$

with q the elementary charge,  $\mu_p$  and  $\mu_n$  the hole and electron mobilities and  $p_0$  the background carrier density.  $\Delta p$  and  $\Delta n$  are related to the hole  $\tau_p$  and electron  $\tau_n$  recombination lifetimes (which are different when  $\Delta p \neq \Delta n$ ) according to:

$$\frac{\Delta n}{\tau_n} = \frac{\Delta p}{\tau_p}.$$
(S2)

Taking into account the parameter  $\beta=rac{\mu_n}{\mu_p}$ , equation (S1) can be written:

$$\sigma = q\mu_p \left( p_0 + \Delta p \left( 1 + \beta \frac{\tau_n}{\tau_p} \right) \right).$$
(S3)

From equation (S3), we can deduce  $\Delta p$  for a p-type material when  $\Delta p \neq \Delta n$ 

$$\Delta p = \frac{\sigma(1-\beta) - q p_0 \Delta \mu}{q \left(1 + \beta \frac{\tau_n}{\tau_p}\right) \Delta \mu}.$$
(54)

Note that  $\Delta \mu = \mu_p - \mu_n$ . In the standard CRPH analysis, considering  $\Delta p = \Delta n$ ,  $\Delta p_{CRPH}$  is given by:

$$\Delta p_{CRPH} = \frac{\sigma(1-\beta^2) - qp_0\Delta\mu(1+\beta)}{q\Delta\mu(1+\beta)^2}.$$
(S5)

Therefore, if we multiply equation (S4) by  $\frac{1+\beta}{1+\beta}$ , we can write:

$$\Delta p(\Delta p \neq \Delta n) = \Delta p_{CRPH} \frac{(1+\beta)}{\left(1+\beta \frac{\tau_n}{\tau_p}\right)}.$$
(S6)

Of note, if  $\Delta p = \Delta n$ , we obtain  $\Delta p = \Delta p_{CRPH}$  from equation (S6), as expected. Moreover, if  $\beta \to 0$ , we also obtain  $\Delta p = \Delta p_{CRPH}$  (even if  $\Delta p \neq \Delta n$ ). Therefore, when there is a large imbalance of mobility, we obtain the lifetime of the majority type carriers, here holes. For a given value of  $\beta$ , we find that:

$$\Delta p(\Delta p \neq \Delta n) \xrightarrow[\tau_n \to 0]{\tau_n} \Delta p_{CRPH} \times (1 + \beta), \text{ and}$$
(S7)

$$\Delta p(\Delta p \neq \Delta n) \xrightarrow[\tau_p \to 1]{\tau_n} \Delta p_{CRPH}$$
(S8)

From equations (S7) and (S8), we know that  $\Delta p$  is bonded between  $\Delta p \left(\frac{\tau_n}{\tau_p} = 1\right)$  and  $\Delta p \left(\frac{\tau_n}{\tau_p} = 0\right)$ , if the function  $\Delta p$  (equation (S6)) monotonously decreases with  $\frac{\tau_n}{\tau_p}$  between 0 and 1. This condition is verified with the determination of  $\frac{d\Delta p}{d\frac{\tau_n}{\tau_p}}$  for  $0 \leq \frac{\tau_n}{\tau_p} \leq 1$ :

$$\frac{d\Delta p}{d\frac{\tau_n}{\tau_p}} = -\frac{\Delta p_{CRPH}(1+\beta)\beta}{\left(1+\beta\frac{\tau_n}{\tau_p}\right)^2}.$$
(59)

Since  $\beta \ge 0$ , we can show that  $\frac{d\Delta p}{d\frac{\tau_n}{\tau_p}} \le 0$  and, therefore, monotonously decreases with  $\frac{\tau_n}{\tau_p}$ . Finally, we can write that, for  $\Delta p \ge \Delta n$ ,

$$\Delta p_{CRPH} \le \Delta p \le \Delta p_{CRPH} (1+\beta). \tag{S10}$$

Since  $\tau = \Delta p/G$ , we can write, for  $\Delta p \ge \Delta n$ ,

$$\tau_{CRPH} \le \tau_p \le \tau_{CRPH} (1+\beta), \tag{S11}$$

where  $\tau_{CRPH}$  is the carrier recombination lifetime calculated with the standard CRPH model considering  $\Delta p = \Delta n$ , and  $\tau_p$  is the true hole recombination lifetime for  $\Delta p \ge \Delta n$ .



**Fig. S20** Sheet resistance  $R_s$  as a function of generation rate *G* for two undoped and two 1% MR (molar ratio) doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub> samples. The grey area highlights the resistance decrease upon photocarrier generation.



**Figure S21.** Carrier-resolved photo Hall *H-o* plot for a second sample of undoped (a) and 1% MR (molar ratio) doped (b) MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub>. A maximum light intensity  $I_{max} = 30 \ mW \ cm^{-2}$  is used.



**Figure S22.** (a) Hole  $\mu_p$  and electron  $\mu_n$  mobilities as a function of generation rate *G* extracted from CRPH measurements for a second sample of 0% and 1% MR (molar ratio) doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub>. (b) Hole lifetime  $\tau_p$  and diffusion length  $L_{d,p}$  for 0% and 1% doped MAPb<sub>0.5</sub>Sn<sub>0.5</sub>I<sub>3</sub>. A linear fit at high light intensity provides the power law  $\tau_p \propto G^{\gamma-1}$ . The grey area highlights the data points obtained at the highest light intensity  $I_{max} = 30 \ mW \ cm^{-2}$  (approximaterly a third of one sun intensity).

## References

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