

## Supplementary Information

### Magnetoelectric coupling effects on band alignments of the multiferroic In<sub>2</sub>Se<sub>3</sub>-CrI<sub>3</sub> trilayer heterostructures

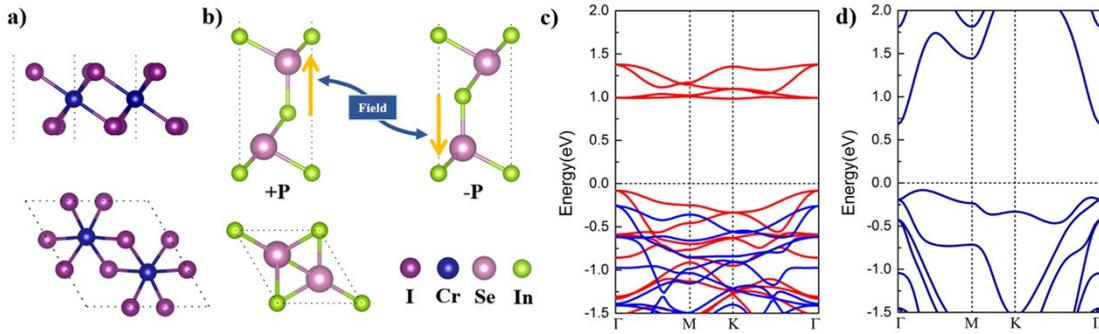
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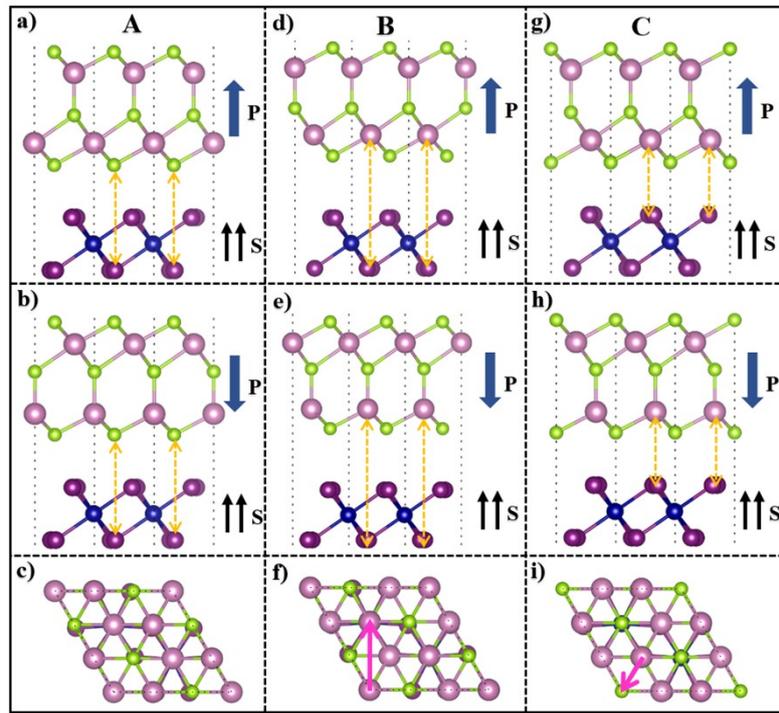
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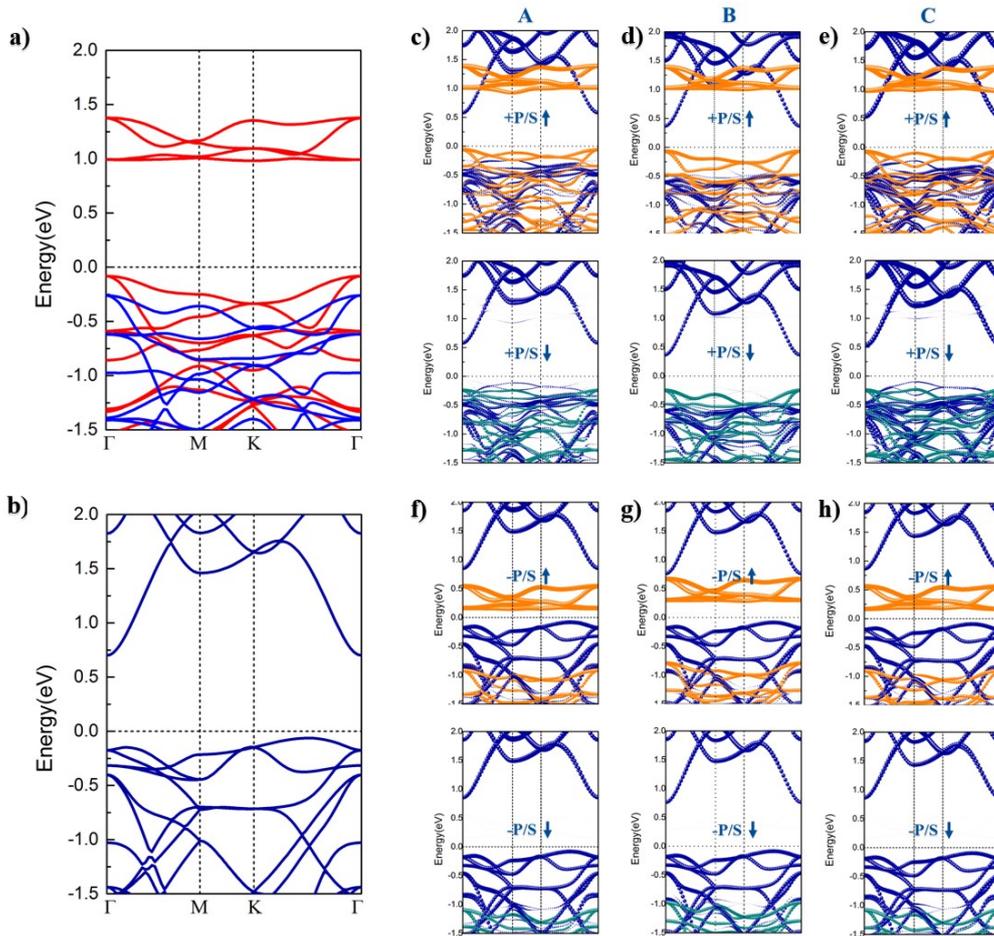
**Fig. S1** Side and top view of (a)  $\text{CrI}_3$  (b)  $\text{In}_2\text{Se}_3$  monolayers, where I, Cr, Se, and In atoms are denoted by purple, dark blue, pink, and light green balls, respectively. The diamond-shaped dotted line represents the primitive cell. The band structure of  $\text{CrI}_3$  and  $\text{In}_2\text{Se}_3$  are shown in (c) and (d). The red and blue lines in the Figure (c) represent the energy bands of the spin-up and spin-down channels.



**Fig. S2** Side (the upper two panels) and top (bottom panels) view of  $\text{CrI}_3/\text{In}_2\text{Se}_3$  vdWHs in three stackings: A, B and C. Here, the dark blue arrow marked as P indicates the direction of FE polarization, and the black arrow labeled as S represents the spin orientation of Cr atoms. In Fig (f) and (i), the translational displacement of  $\text{In}_2\text{Se}_3$  relative to  $\text{CrI}_3$  are marked by magenta arrows.

For the vertical  $\text{CrI}_3/\text{In}_2\text{Se}_3$  multiferroic vdWHs, we constructed three

typical atomic stackings, i.e. A, B, and C. In view of two polarization states of  $\text{In}_2\text{Se}_3$ , each stacking corresponds to two structures, as shown in Fig. S1. The corresponding stabilities and structural parameters are depicted in Table S2. For A superposition, the lower I atoms are located directly below the bottom Se atoms of  $\text{In}_2\text{Se}_3$ , while the two Cr atoms are located directly below the bottom In atoms of  $\text{In}_2\text{Se}_3$ . We constructed the B and C stacking patterns by shifting the  $\text{In}_2\text{Se}_3$  of the A stacking along the directions marked by pink arrows.



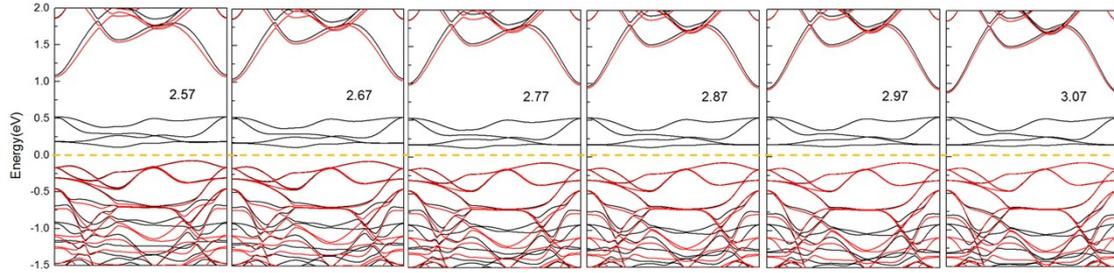
**Fig. S3** Spin-polarized band structures of (a) monolayer  $\text{CrI}_3$  (b)  $\sqrt{3}\times\sqrt{3}$   $\text{In}_2\text{Se}_3$ . (c-h)  $\text{CrI}_3/\text{In}_2\text{Se}_3$  bilayer vdWHs with A, B and C three stacking patterns. The red and blue lines in Fig. (a) represent the band structures of the spin-up and spin-down channels. In Fig. (c-h), orange and green lines represent the contribution of  $\text{CrI}_3$ , and dark blue lines represents the contribution of  $\text{In}_2\text{Se}_3$ .

Fig. S3. shows the spin-polarized band structures of above bilayer

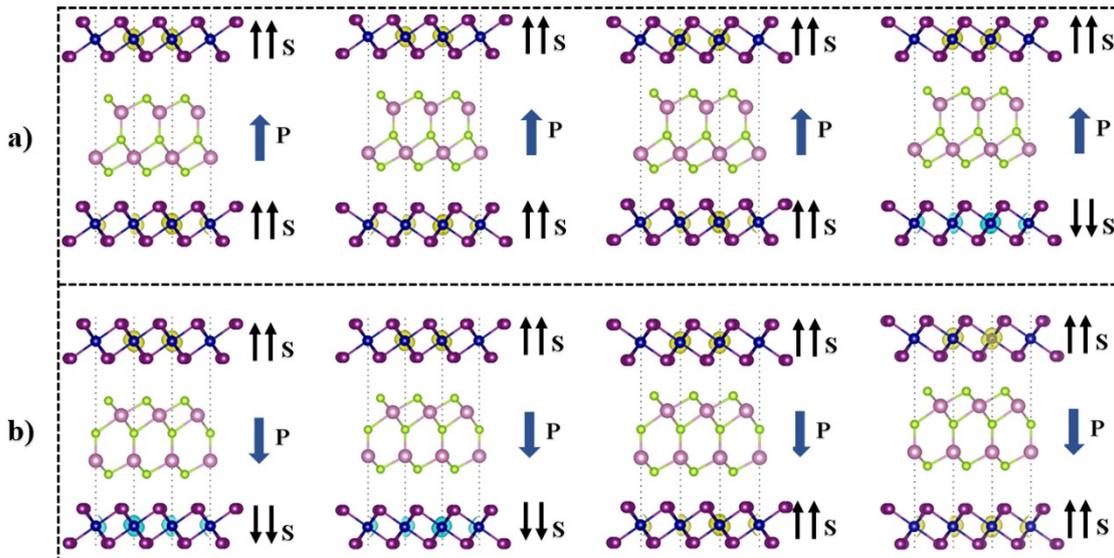
vdWHs. Notably, whether combined with +P or -P  $\text{In}_2\text{Se}_3$ ,  $\text{CrI}_3$  still retains the intrinsic semiconductor. As seen from the Figs. S3(c-e) or S3(f-h), different stacking patterns have little effect on band structures except for a slight shift of the band edges. For +P cases [Fig. S3(c-e)], the band structures of vdWHs can be regarded as a simple superposition of two monolayers. Thus both  $\text{CrI}_3$  and  $\text{In}_2\text{Se}_3$  maintain their intrinsic semiconductor characteristics well. Comparing the Figs. S3(c) and 3(f), it is observed that for the same atoms superposition, the band edge shifts downward when the polarization is reversed to -P case, resulting that the conduction band edge of  $\text{CrI}_3$  approach to the fermi level ( $E_f$ ).

Usually, a smaller interlayer distance within vdWH indicates a stronger vdW interlayer interaction. Table S2 reveals that the A configuration has the strongest interlayer interaction among the three stacking configurations. Moreover, as shown in Fig. S3, the conduction band minimum (CBM) of A stacking with -P polarization is the closest to the Fermi level ( $E_f$ ). Both the strongest interlayer interaction and the CBM closest to the  $E_f$  decide that the A stacking is the most promising to be regulated as a half-metal. Thus, in the following calculations, we conducted some further investigations for the A stacking.

Decreasing the interlayer distance  $d$  is a straightforward way to tune the features of 2D materials. The interlayer distance has a very limited influence on the electronic structure of the  $\text{CrI}_3$  monolayer inside  $\text{CrI}_3/(-\text{P})\text{In}_2\text{Se}_3$  vdWH [Fig. S3]. At this point, we focused on other methods to regulate the magnetoelectric coupling effect, including adding a FE or FM monolayer to the hetero-bilayer structures.

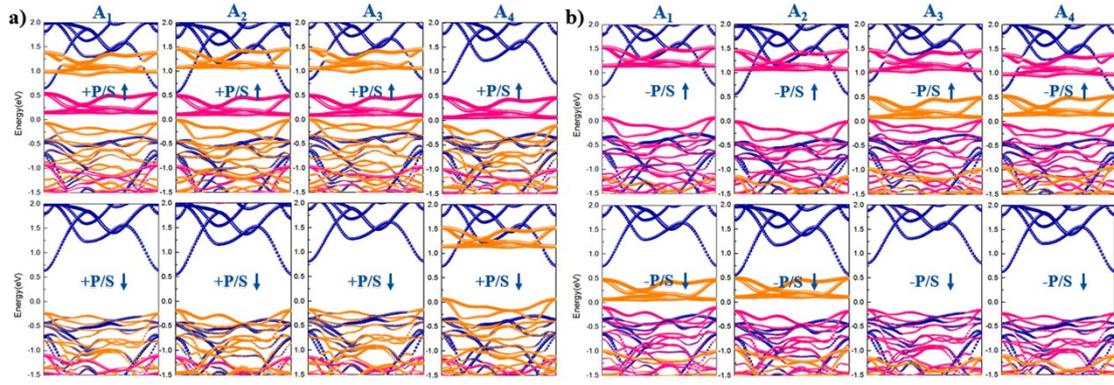


**Fig. S4** The effect of interlayer distance  $d$  on the band structures of  $\text{CrI}_3/\text{In}_2\text{Se}_3$  vdWHs for A stacking configuration, ranging from  $2.57\text{\AA}$  to  $3.07\text{\AA}$ .

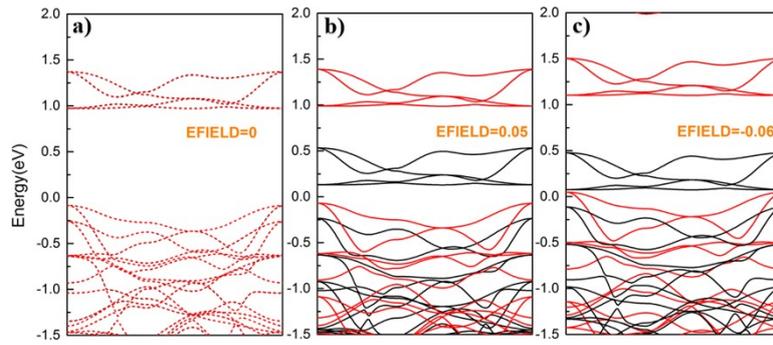


**Fig. S5** The structure of  $\text{CrI}_3/\text{In}_2\text{Se}_3/\text{CrI}_3$  vdWHs with  $A_1$ ,  $A_2$ ,  $A_3$  and  $A_4$  stackings for (a) +P and (b) -P case, where the yellow (blue) color around Cr atoms means the spin-up (spin-dw) charge distribution. The black and dark blue arrow arrows indicate the spin orientation of Cr atoms and the direction of ferroelectric polarization.

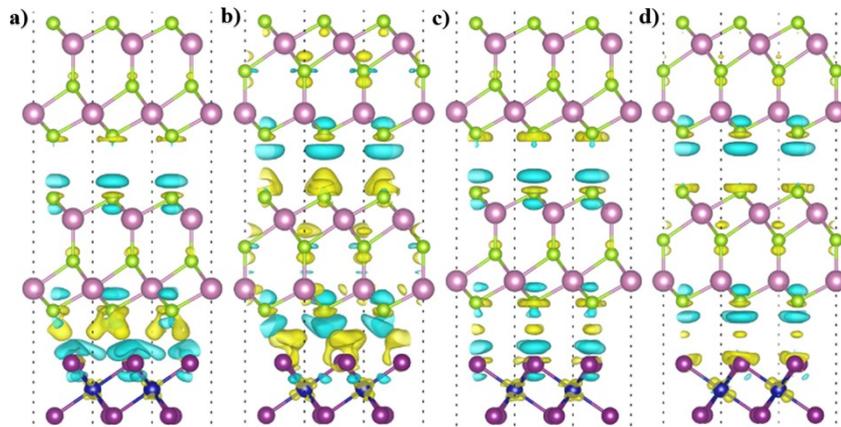
The  $A_1$  pattern was constructed by inserting  $\text{In}_2\text{Se}_3$  layer between bilayer  $\text{CrI}_3$  layers. As for the  $A_2$  stacking, the lower  $\text{CrI}_3$  layer was flipped vertically along the diagonal direction with respect to  $A_1$  stacking. For  $A_3$  superposition, flipping engineering was applied to the upper  $\text{CrI}_3$  layer. And for  $A_4$  alignment, both  $\text{CrI}_3$  layers were flipped vertically.



**Fig. S6** The spin- and layer-resolved band structures of  $\text{CrI}_3/\text{In}_2\text{Se}_3/\text{CrI}_3$  vdWHs with (a) +P case (b) -P case, where the  $A_1$ ,  $A_2$ ,  $A_3$  and  $A_4$  configurations represent the structures without flipping and applying flipping, respectively. The orange and magenta lines represent the contributions of different  $\text{CrI}_3$  layers with the majority spin ( $S\uparrow$ ) and minority spin ( $S\downarrow$ ) channels, while dark blue lines denote the contributions of  $\text{In}_2\text{Se}_3$ .

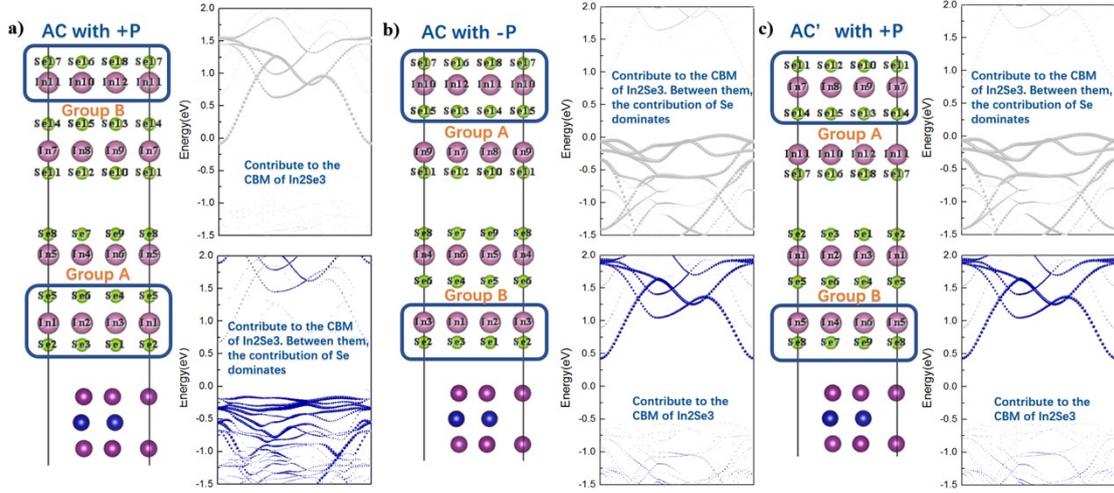


**Fig. S7** The band structures of bilayer  $\text{CrI}_3$  under certain E-fields (a)  $0\text{eV}/\text{\AA}$ , (b)  $0.05\text{eV}/\text{\AA}$ , (c)  $-0.06\text{eV}/\text{\AA}$ .



**Fig. S8** Differential charge density of  $\text{CrI}_3/2\text{L-In}_2\text{Se}_3$  trilayer vdWHs with (a)  $A_1$  stacking for +P case, (b)  $A_1$  stacking for -P case, (c)  $A_{II}$  stacking for +P case and (d)  $A_{II}$  stacking for -P case.

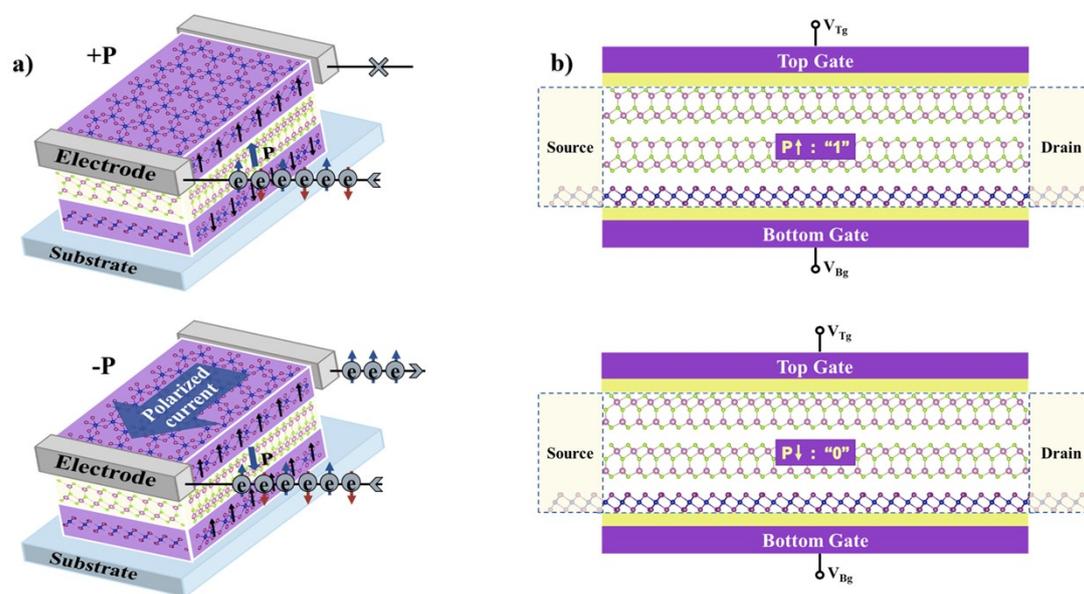
$A_{II}$  stacking for  $-P$  case.



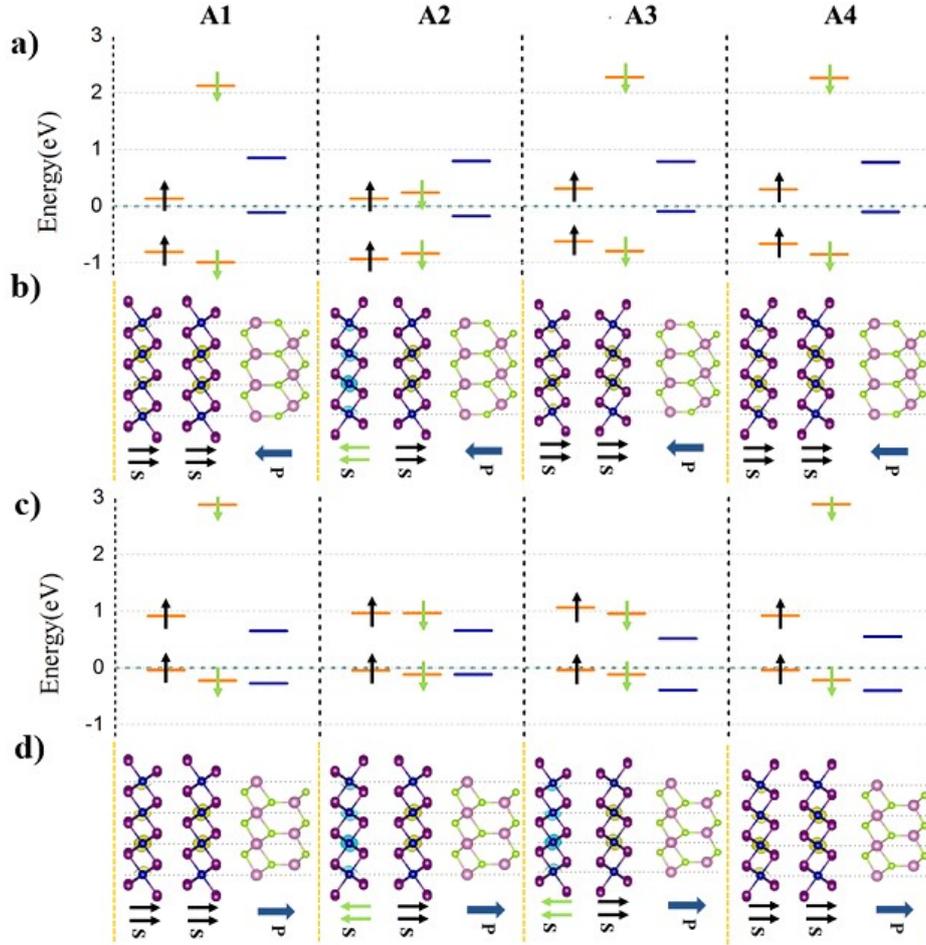
**Fig. S9** Contribution of atoms to CBM and VBM for (a)  $A_I$  stacking with  $+P$  polarization (b)  $A_I$  stacking with  $-P$  polarization (c)  $A_{II}$  stacking with  $+P$  polarization.

In order to clarify the internal mechanism of FM effect on FE, the contribution of atoms to CBM and VBM are drawn. From Fig. S9(a-c), we find that the VBM of the  $In_2Se_3$  part is contributed by a set of Se/In/Se atomic layers (marked as group A) with a sandwich structure, of which the main contributions are come from the two Se layers. However, the  $In_2Se_3$  part of the CBM is contributed by a set of two layers of Se/In atomic layers (labeled as group B). Whether the  $In_2Se_3$  band gap is opened or not depends on the B group atoms. For the  $A_I$  superposition with  $-P$  configuration [Fig. S9(b)], the B group (consisting of 1-3 Se atoms and 1-3 In atoms) is so close to  $CrI_3$  that the band gap of  $In_2Se_3$  is opened under the influence of the magnetic proximity effect. On the contrary, for the  $A_I$  stacking with  $+P$  polarization [Fig. S9(a)], the B group atoms (consisting of 16-18 Se atoms and 10-12 In atoms) are far away from the influence of FM  $CrI_3$ , which tends to maintain the properties of the bilayer  $In_2Se_3$ . In order to verify this result, we reversed the polarizations of the two layers of  $In_2Se_3$  separately in  $A_I$  stacking with  $+P$  case, i.e., the  $A_{II}$  superposition with  $+P$  configuration [Fig. S9(c)]. And then the new 7-9 Se atoms and 4-6 In atoms

compose a new group B, which is closer to  $\text{CrI}_3$ . As we conjectured, the band gap of  $\text{In}_2\text{Se}_3$  is opened. For the case of  $\text{CrI}_3$  flipped vertically, the principle is the same to the above. However, owing to the mirror asymmetry, the influence of the flipped  $\text{CrI}_3$  on the ferroelectricity is smaller than that before, so that the band gap of  $\text{In}_2\text{Se}_3$  opened by  $\text{CrI}_3$  becomes smaller.



**Fig. S10** Two prototypes of spin-field effect transistors (a) atom-thick multiferroic memory devices based on the  $\text{CrI}_3/\text{In}_2\text{Se}_3/\text{CrI}_3$  vdWH (b) spin switch based on the  $\text{CrI}_3/2\text{L-In}_2\text{Se}_3$  vdWH



**Fig. S11** Band alignments of  $\text{In}_2\text{Se}_3/2\text{L-CrI}_3$  vdWHs with  $A_1$ ,  $A_2$ ,  $A_3$  and  $A_4$  stackings for (a) +P polarization and (c) -P case, where the orange and dark blue line represent the contribution of  $\text{CrI}_3$  and  $\text{In}_2\text{Se}_3$ . And the black and green arrows in (a) and (c) indicate the spin-up channel and the spin-down channel energy level. The  $A_1$ ,  $A_2$ ,  $A_3$  and  $A_4$  configurations represent the structures without flipping and applying flipping, respectively. The corresponding structure are shown in (b) and (d), where the yellow and blue color around Cr atoms mean the spin-up and spin-down charge distribution. The black and green arrows in (b) and (d) indicate the spin orientation of Cr atoms, and the dark blue arrow represents the direction of FE polarization.

For the vertical  $\text{In}_2\text{Se}_3/2\text{L-CrI}_3$  multiferroic vdWHs, we can construct four typical atomic stackings and place bialyer  $\text{CrI}_3$  on the side of  $\text{In}_2\text{Se}_3$ , marked as  $A_1$ ,  $A_2$ ,  $A_3$  and  $A_4$ , shown in Fig. 12(b). For +P case with  $A_3$  stacking, the conduction bands near the  $E_f$  belong to the spin-down channel associated with the upper  $\text{CrI}_3$  layer. On the contrary, for the -P case, the

conduction bands near the  $E_f$  belong to the spin-up channel associated with the lower  $\text{CrI}_3$  layer. The result provides the field-dependent spin filtering in which one type of spin current can be switched to another type of spin current through an electric field-induced transition from +P to -P configuration.

**Table S1** The lattice constant  $\mathbf{a}$ , interlayer distance  $\mathbf{d}$ , binding energy  $\mathbf{E}_b$ , exchange energy  $\mathbf{E}_{\text{ex}}$  and magnetic ground state **MGS** for  $\text{CrI}_3/\text{In}_2\text{Se}_3$  vdWHs with A, B and C stacking patterns, where +P and -P stand for upward polarization and downward polarization, respectively.

Stackings	P	$\mathbf{a}(\text{\AA})$	$\mathbf{d}$	$\mathbf{E}_b(\text{eV}/\text{\AA}^2)$	$\mathbf{E}_{\text{ex}}(\text{meV})$	<b>MGS</b>
A	+P	7.096	3.141	-1.676522	64.029	FM
	-P	7.095	3.171	-1.675411	63.346	FM
B	+P	7.096	3.878	-1.666743	64.116	FM
	-P	7.095	3.889	-1.666096	63.492	FM
C	+P	7.096	3.155	-1.677653	64.285	FM
	-P	7.095	3.179	-1.676339	63.135	FM

To estimate the stability of the 2D multiferroic vdWHs, the binding energy ( $E_b$ ) is calculated by the following expression:

$$E_b = (E_{\text{HS}} - E_{\text{In}_2\text{Se}_3} - E_{\text{CrI}_3}) / A$$

where  $E_{\text{HS}}$ ,  $E_{\text{In}_2\text{Se}_3}$ ,  $E_{\text{CrI}_3}$  represent the total energies of the multiferroic vdWHs, isolated  $\text{CrI}_3$  and  $\text{In}_2\text{Se}_3$  monolayer,  $A$  is the interfacial area, as shown in Table S1. According to this definition, the more negative the  $E_b$  is, the more stable the structure can be in terms of energy. Such negative  $E_b$  proves that all stacking configurations are stable. It can be seen from Table S1 that the interlayer distances of +P and -P cases are highly close for the same stacked configuration. Taking the A configuration as example, the interlayer distances along the z direction for  $\text{CrI}_3 / (+\text{P})\text{In}_2\text{Se}_3$  and  $\text{CrI}_3 / (-\text{P})\text{In}_2\text{Se}_3$  vdWHs are 3.14 and 3.17  $\text{\AA}$ , respectively. Such interlayer distances illustrate the nature of the van der Waals combination between

CrI<sub>3</sub> and In<sub>2</sub>Se<sub>3</sub>. The interlayer interaction energy difference  $E_{\text{bis}}$  ( $\approx 1.15$  meVÅ<sup>-2</sup>) between the CrI<sub>3</sub>/(+P)In<sub>2</sub>Se<sub>3</sub> and CrI<sub>3</sub>/(-P)In<sub>2</sub>Se<sub>3</sub> vdWHs reveals the bistable nature of the heterojunction switch based on FE In<sub>2</sub>Se<sub>3</sub>.

**Table S2** The average electrons of the Cr-3d orbital in the free-standing monolayer CrI<sub>3</sub> and bilayer vdWHs

Stackings	P	t2g		eg		total
		Spin up	Spin down	Spin up	Spin down	
CrI <sub>3</sub>		2.042500	0.283000	1.714600	0.0858	4.126
A	+P	2.048600	0.273900	1.711600	0.0829	4.117
	-P	2.047500	0.274700	1.711700	0.083	4.1168
B	+P	2.048300	0.273900	1.711400	0.0831	4.1167
	-P	2.046600	0.274600	1.712000	0.0831	4.1163
C	+P	2.046100	0.276000	1.711600	0.0832	4.1169
	-P	2.044500	0.276400	1.711200	0.0844	4.1165

**Table S3** The average electrons of the Cr-3d orbital in CrI<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub>/CrI<sub>3</sub> vdWHs with A<sub>1</sub>, A<sub>2</sub>, A<sub>3</sub>, and A<sub>4</sub> stackings.

P	stacking	atom	t2g		eg		total	MGS
			Spin up	Spin down	Spin up	Spin down		
+P	A <sub>1</sub>	Cr <sub>1</sub>	2.0369	0.2817	1.7122	0.0831	4.1139	FM
		Cr <sub>2</sub>	2.0382	0.2821	1.7121	0.0844	4.1168	
		Cr <sub>3</sub>	2.0384	0.2796	1.7127	0.0833	4.1140	
		Cr <sub>4</sub>	2.0372	0.2806	1.7126	0.0823	4.1128	
	A <sub>2</sub>	Cr <sub>1</sub>	2.0398	0.2815	1.7126	0.0840	4.1179	FM
		Cr <sub>2</sub>	2.0391	0.2807	1.7107	0.0844	4.1150	
		Cr <sub>3</sub>	2.0383	0.2791	1.7120	0.0838	4.1131	
		Cr <sub>4</sub>	2.0374	0.2804	1.7136	0.0834	4.1149	
	A <sub>3</sub>	Cr <sub>1</sub>	2.0382	0.2823	1.7117	0.0836	4.1158	FM
		Cr <sub>2</sub>	2.0382	0.2808	1.7116	0.0843	4.1149	
		Cr <sub>3</sub>	2.0369	0.2827	1.7122	0.0840	4.1158	
		Cr <sub>4</sub>	2.0349	0.2831	1.7130	0.0839	4.1148	
	A <sub>4</sub>	Cr <sub>1</sub>	0.2784	2.0392	0.0823	1.7131	4.1129	AFM
		Cr <sub>2</sub>	0.2790	2.0377	0.0825	1.7125	4.1116	

		Cr <sub>3</sub>	2.0375	0.2823	1.7110	0.0835	4.1143		
		Cr <sub>4</sub>	2.0380	0.2818	1.7130	0.0831	4.1160		
	A <sub>1</sub>	Cr <sub>1</sub>	0.2820	2.0374	0.0835	1.7131	4.1160	AFM	
		Cr <sub>2</sub>	0.2813	2.0378	0.0843	1.7124	4.1159		
		Cr <sub>3</sub>	2.0395	0.2792	1.7134	0.0824	4.1145		
		Cr <sub>4</sub>	2.0387	0.2782	1.7131	0.0828	4.1128		
	A <sub>2</sub>	Cr <sub>1</sub>	0.2794	2.0391	0.0828	1.7133	4.1146	AFM	
		Cr <sub>2</sub>	0.2798	2.0392	0.0829	1.7130	4.1148		
		Cr <sub>3</sub>	2.0394	0.2802	1.7124	0.0831	4.1151		
-P		Cr <sub>4</sub>	2.0388	0.2804	1.7132	0.0835	4.1160		
		Cr <sub>1</sub>	2.0356	0.2814	1.7117	0.0839	4.1125		
		A <sub>3</sub>	Cr <sub>2</sub>	2.0361	0.2807	1.7124	0.0841	4.1133	FM
			Cr <sub>3</sub>	2.0372	0.2811	1.7136	0.0834	4.1153	
			Cr <sub>4</sub>	2.0357	0.2801	1.7133	0.0829	4.1121	
		A <sub>4</sub>	Cr <sub>1</sub>	2.0378	0.2807	1.7135	0.0827	4.1147	FM
			Cr <sub>2</sub>	2.0384	0.2808	1.7121	0.0831	4.1144	
			Cr <sub>3</sub>	2.0360	0.2806	1.7130	0.0829	4.1124	
			Cr <sub>4</sub>	2.0377	0.2803	1.7132	0.0838	4.1150	

**Table S4.** The overview of the applicable devices.

structures	+P	-P	devices
			Spin switch
CrI <sub>3</sub> /In <sub>2</sub> Se <sub>3</sub> /CrI <sub>3</sub>	semiconductor	half-metal	Atom-thin multiferroic memory device
			Spin switch
CrI <sub>3</sub> /2L-In <sub>2</sub> Se <sub>3</sub>	half-metal	semiconductor	Atom-thin multiferroic memory device
			Field-dependent spin filtering
CrI <sub>3</sub> /2L-CrI <sub>3</sub>	semiconductor (S↓)	semiconductor (S↑)	