Electronic Supplementary Information:

From crystallographic data to solution structure of photoreceptors through Molecular Dynamics and Spectroscopy Simulations: the case of AppA BLUF domain.

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S1 Supplementary Tables

Table S1 Equilibration protocols for MD simulations

Structures	MD Simulations	System preparation/Heating/Equilibration					
		Residue at position 20	Temperature (K)	Thermodynamic ensemble	Time (ps)	Harmonic restraints	
Trp _{in}	MD1/MD2		0-50	NVT	50		
		Cys4	50-300	NPT	250	protein backbone	
			300	NPT	750		
	MD3/MD4	Cys4	0-100	NVT	50	protein backbone	
			100-315	NPT	250	protein backbone except for loops (resid 95-103,108-111)	
			315-300	NPT	2000		
			300	NPT	1000		
Met _{in}	MD1/MD2	Ser4	0-50	NVT	50	protein backbone	
			50-300	NPT	250		
			300	NPT	750		
	MD3/MD4-S20C	Cys4	0-50	NVT	50		
			50-300	NPT	250	protein backbone	
			300	NPT	750		

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Table S2 Probabilities of hydrogen bond formation during MD simulations. Hydrogen bonds were calculated using CPPTRAJ on snapshots sampled every 100 ps using a cutoff of 3.2 Å on the donor-acceptor distance and a cutoff of 135° on the donor-hydrogen-acceptor angle.

		Gln63(NH)- FMN(O4)	Gln63(NH)- FMN(N5)	Tyr21(OH)- Gln63(O)	Tyr21(O)- Gln63(NH)	Asn45(NH)- FMN(O4)	Asn45(O)- FMN(N3H)	Trp104(H)- Gln63(O)
Trp _{in}	MD1	14	29	69	7	62	57	7
	MD2	2	47	13	7	91	77	1
	MD3	23	19	86	4	60	38	5
	MD4	18	33	53	2	92	81	1
	Average	14	32	55	5	76	63	4
Met _{in}	MD1	14	44	96	0	72	92	0
	MD2	7	52	93	0	68	91	0
	MD3-S20C	4	48	56	0	59	91	0
	MD4-S20C	7	56	93	0	68	88	0
	Average	8	50	85	0	67	90	0

S2 Supplementary Figures



Fig. S1 Stability of Met_{in} trajectories over Trp_{in} ones. Distribution of RMSD values calculated between the Trp_{in} (magenta), Met_{in} (green) and Met_{in} -S20C (cyan) MD simulations, calculated from the respective crystal structures. All replicas are reported. RMSD values were calculated on backbone atoms coordinates of snapshots extracted from the production trajectory every 100 ps.



Fig. S2 Distributions of the distances detecting Gln 63 motion along the MD simulations. (a) Distribution of the distance between the oxygen of Gln63 and C6 of FMN. (b) Distribution of the difference between FMN(O4)-Gln63(O) distance denoted as d_{OO} and FMN(O4)-Gln63(N) distance denoted as d_{OO} . (c) Distribution of the distance between the oxygen of Gln63 and the hydroxy group of Tyr21. Trp_{in}, Met_{in}, Met_{in}-S20C simulations are shown in magenta, green and cyan, respectively. The distances obtained from the Trp_{in} (1YRX) (magenta) and Met_{in} (2IYG) (green) crystal structures are represented as dotted lines.



Fig. S3 Conformational dynamics of residues in the binding pocket during the Trp_{in} simulations (pink hues). (a) Time evolution of the N45(OD1)-FMN(H3) distance. (b) Time evolution of the N45(HD21)-FMN(O4) distance.



Fig. S4 Conformational dynamics of residues in the binding pocket during the Met_{in} (green hues) and Met_{in} -S20C simulations (cyan hues). (a) Time evolution of the N45(OD1)-FMN(H3) distance. (b) Time evolution of the N45(HD21)-FMN(O4) distance.



Fig. S5 Orientation of Gln63 along four Trp_{in} replicas as measured by the difference between FMN(O4)-Gln63(O) distance d_{OO} and FMN(O4)-Gln63(N) distance d_{ON} . The values corresponding to the two available crystal structures are shown as horizontal dotted lines (2IYG in green and 1YRX in pink).



Fig. S6 Effect of Trp104 on the chemical shifts of Asn45 and Gln63 amide protons. (a) Scatter plot of amide protons chemical shifts calculated with the full-QM cluster and with Trp104 excluded from the QM part. The dashed x = y line serves as a guide to the eye. (b) Distribution of chemical shifts for the Asn45 HD22 atom, with Trp104 included (red) or excluded (blue) from the QM part.



Fig. S7 Correlation between $C_4 = O_4$ QM/MM harmonic frequencies and hydrogen bond distance from Gln63 amino hydrogen to the FMN O₄ atom.



Fig. S8 Distributions of the distances detecting Gln63 motion along the Trp_{in} (magenta) and Met_{in} (green) QM/MMPol MD simulations. (a) Distribution of the distance between the oxygen of Gln63 and C6 of FMN. (b) Distribution of the difference between FMN(O4)-Gln63(O) distance denoted as d_{OO} and FMN(O4)-Gln63(N) distance denoted as d_{ON} . (c) Distribution of the distance between the oxygen of Gln63 and the hydroxyl group of Tyr21. The atom numbering is reported in Fig.3



Fig. S9 Mode resolution with SOBI independent component analysis on one of the QM/MM MDs. The power spectra of the $C_2=O_2$ (black) and $C_4=O_4$ (red) bond lengths are shown as solid lines, and offset by 1200 on the *y* axis for clarity. The power spectra of the SOBI coordinates with the largest contribution from $C_2=O_2$ (black) or $C_4=O_4$ (red) are shown as dashed lines.



Fig. S10 Comparison between $C_2=O_2$ and $C_4=O_4$ harmonic IR frequencies calculated with b3lyp-D3 and wb97XD functionals. Frequencies were calculated on ten QM/MM optimized structures corresponding to Met_{in} conformation. For all calculations we used the 6-31G(d) basis set.

S3 QM/MMPol simulations: computational details

QM/MMPol simulations were carried out using an interface¹⁻³ between the molecular dynamics package Tinker^{4,5} and the development version of Gaussian^{6,7} which implements an embedding based on the AMOEBA force field.^{8,9} For these simulations, initial configurations were extracted from the classical molecular dynamics trajectories: ten of them from the Trp_{in} dynamics and ten of them from the Met_{in} dynamics. For each of the selected frames, we selected a shell containing only the protein and solvent residues within 30 Å from the FMN residue to be used for the QM/MMPol simulations. The FMN isoalloxazine ring was treated using a QM method, whereas the remaining part of the system was described using the AMOEBA force field. As QM level of theory we chose the DFT ω b96XD/6–31G(d)). This functional shows a similar trend (see figure S10 to b3lyp-D3 for C₂O₂ and C₄O₄ ONIOM QM/MM harmonic frequencies calculated on Met_{in}-conformation structures. Non periodic boundary conditions were included by freezing all the residues beyond 22 Å from the FMN, moreover, to speed up the calculation, we set to zero all the AMOEBA polarizabilities of the frozen residues. To generate suitable starting condition for the molecular dynamics simulations, on each of the structures, we first performed a QM/MMPol geometry optimization using the limited-memory Broyden-Fletcher-Goldfarb-Shanno algorithm.¹⁰ Then, QM/MMPol simulations in the NVT esemble were classically propagated for 10 ps, using an integration step of 0.5 fs and the Velocity Verlet algorithm. A constant temperature of 300 K was kept using the Bussi thermostat.¹¹

S4 Frequency extraction from QM/MM MD simulations

Vibrational spectra can be calculated from MD simulations via autocorrelation functions. Power spectra of atomic velocities, *i.e.* the Fourier transform of the velocity autocorrelation function, contain all the vibrational frequencies of the system ¹². While for simple systems the velocity power spectra can be analyzed and decomposed in terms of effective normal modes¹³, our QM/MM simulations cannot be easily analyzed in this way. Instead, we used internal coordinates (bond lengths and angles) to compute power spectra that are directly related to specific normal modes. Example of power spectra for the C₂=O₂ and C₄=O₄ bond lengths are shown in Figure S9 (solid lines). Even though the power spectra are quite localized in frequency around 1800 cm⁻¹, a smaller contribution appears at ~1450 cm⁻¹, due to the mixing of C=O modes with ring modes. In order to obtain effective normal modes, we employed linear combinations of internal coordinates.

Assuming that each internal coordinate evolves as a combination of independent signals (normal modes), we can use a SOBI, a signal-processing algorithm, to disentangle the contributions from various normal modes. The independent components s(t) can be written as linear combinations of the internal coordinates $\mathbf{x}(t)$:

$$\mathbf{s}(t) = A\mathbf{x}(t)$$

The SOBI algorithm relies on the property that independent signals have negligible cross-correlation functions at all time delays. Therefore, independent components should satisfy the relation

$$\langle s_i(t)s_j(t+\tau)\rangle = 0$$
 $i \neq j$

for different time delays τ . By defining the time-lagged correlation matrix,

$$R_{\tau} = \langle \mathbf{x}(t)\mathbf{x}(t+\tau)^T \rangle$$

the above condition is equivalent to imposing the diagonality of the matrix $AR_{\tau}A^{T}$. In practice, one needs to find A such that $AR_{\tau}A^{T}$ is approximately diagonal for several values of τ . We use a joint diagonalization algorithm based on Jacobi angles¹⁴ in order to find the best transformation A. Correlation matrices were built at lag times 50, 100, ..., 500 fs and then averaged over all QM/MM MD trajectories. In this way, a single common transformation matrix was obtained for all trajectories. This ensures that exactly the same normal-mode coordinates are used to construct the power spectra.

The power spectra of the SOBI-transformed coordinates are shown at the bottom of Figure S9 (dashed lines) for an example trajectory. As it can be seen, the SOBI algorithm is able to resolve the two frequencies, and to remove the small contribution at \sim 1450 cm⁻¹.

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