

1 **Supporting Information:**

2

3

4 **Plasma-assisted oxidation of Cu(100) and Cu(111)**

5

6 Sebastian Kunze<sup>a,b</sup>, Liviu C. Tănase<sup>a</sup>, Mauricio J. Prieto<sup>a</sup>, Philipp Grosse<sup>a,b</sup>, Fabian Scholten<sup>a,b</sup>,

7 Lucas de Souza Caldas<sup>a</sup>, Dennis van Vörden<sup>b</sup>, Thomas Schmidt<sup>a,\*</sup>, Beatriz Roldan Cuenya<sup>a\*</sup>

8 <sup>b</sup> *Department of Interface Science, Fritz-Haber Institute of the Max Planck Society, 14195*

9 *Berlin, Germany*

10 <sup>b</sup> *Department of Physics, Ruhr-University Bochum, 44780 Bochum, Germany*

11 *\*E-Mail: [roldan@fhi-berlin.mpg.de](mailto:roldan@fhi-berlin.mpg.de), [schmidtt@fhi-berlin.mpg.de](mailto:schmidtt@fhi-berlin.mpg.de)*

12

13

14

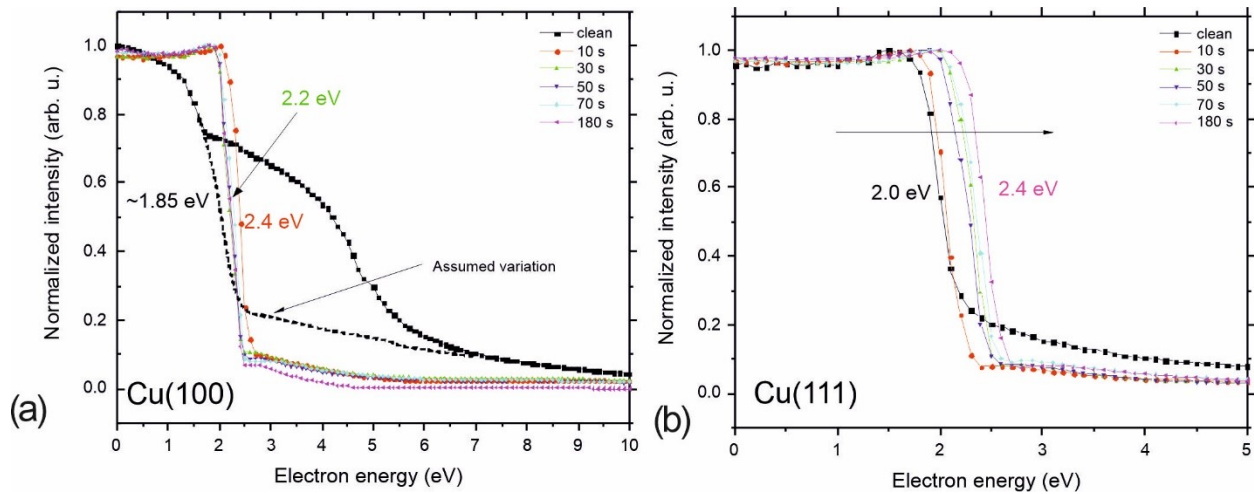
15

16

17

# 1 Work function measurements

2



3

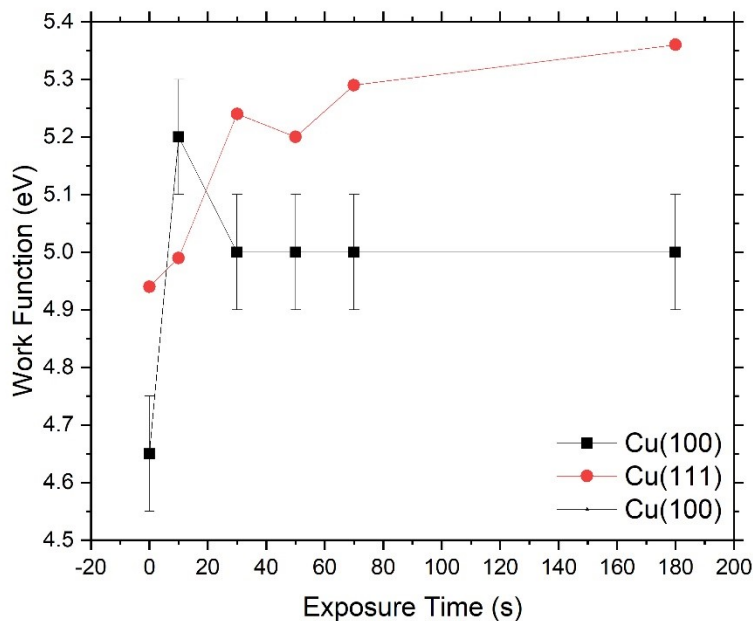
4 **Figure S1:** LEEM I-V curves showing the MEM-LEEM transition during the initial oxidation  
5 steps of a (a) Cu(100) and (b) Cu(111) single crystals. The plasma treatment was done at RT in  
6  $4 \times 10^{-4}$  mbar  $O_2$ .

7

8

9 The LEEM I-V curves in **Fig. S1** allow the determination of the MEM-LEEM transition value,  
10 which represents the value at which the electron kinetic energy can overcome the potential energy  
11 of the surface and penetrate into the bulk. Due to this effect, the MEM-LEEM transition values  
12 are directly related with the surface. For calibration purposes, the transition energy of the clean  
13 surface is aligned to the known work function and used as reference for the oxygen-exposed  
14 surfaces.

15 Additionally, we determined the change in surface work function (WF) due to plasma-induced  
16 oxidation. Therefore, the image intensity has been recorded in dependence on the sample voltage  
17 (LEEM I-V) around the MEM-LEEM transition, e.g. from the voltage range where all electrons  
18 are reflected in front of the surface up to voltages where the electrons can overcome the surface  
19 potential and are partially reflected within the crystal. Considering that the MEM-LEEM transition  
20 can be correlated with the surface WF, one can track WF changes as function of the overall plasma  
21 treatment time. **Fig. S2** presents the WF variation as determined from the I-V curves represented  
22 in **Fig. S1** for both crystal orientations during the initial oxidation steps, using the WF of pristine  
23 Cu(100) at  $4.65 \text{ eV}^1$ , respectively  $4.94 \text{ eV}$  for Cu(111)<sup>2</sup> as references. It is important to note that  
24 the MEM-LEEM is not steep in the case of the clean Cu(100) crystal, due to the presence of a  
25 Bragg gap in the vicinity of the Fermi Level<sup>3, 4</sup>, which might be misleading for the correct  
26 determination of the MEM-LEEM transition. In the present case, we included in **Fig. S1(a)** an  
27 additional reference curve (black dotted) of the assumed intensity variation in the absence of the  
28 signal resulting from the electron reflection on the unoccupied states, situation in which one can  
29 infer a MEM-LEEM transition of  $1.85 \text{ eV}$  for the pristine Cu(100) crystal.



1

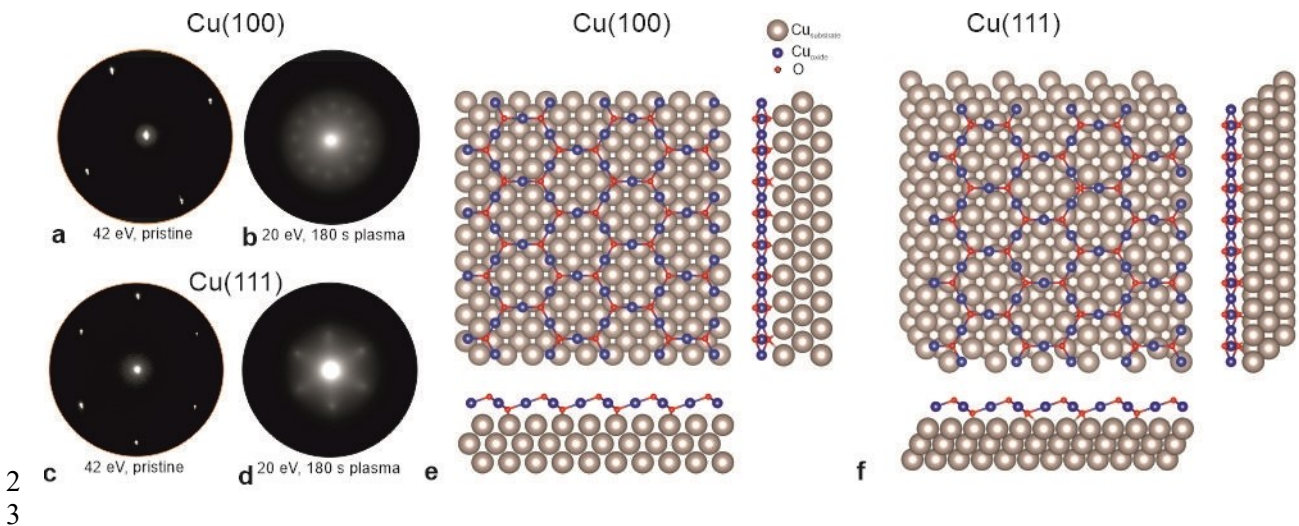
2 **Figure S2:** Work function measurements during the oxide films grown on Cu(100) (black) and  
 3 Cu(111) (red)

4

5 For the two Cu surface orientations there is already a significant difference between the initial,  
 6 i.e. clean state, and the first plasma treatment of 10 s in  $4 \times 10^{-4}$  mbar  $O_2$ . The WF suffers a shift of  
 7 +0.55 eV on Cu(100), which can be correlated also with the formation of the  $c(2 \times 2)$   
 8 superstructure, as observed in LEED (see **Fig. 6**), while the (111) crystal does not show a  
 9 reasonable shift. This difference points out not only to the distinct intrinsic nature of the two  
 10 orientations, but could also indicate a different initial oxidation.

11 Upon further dosage, the WF of Cu(100) gets stabilized to 5.0 eV. In the case of Cu(111),  
 12 even though it does not show any further change after the first treatment, there is a gradual shift  
 13 of +0.4 eV up to a total treatment of 180 s at  $4 \times 10^{-4}$  mbar  $O_2$ , which translates to a WF value of  
 14 5.36 eV, with 0.36 eV higher than the one obtained for the Cu(100) surface. The work function of  
 15 the oxide surface is in good agreement with published values in the range between 5.3 and 5.5  
 16 eV<sup>5</sup>, but values of 4.7-5.5 eV have been also reported for CuO<sup>5</sup>. The comparison of the two curves  
 17 in **Fig. S2** proves that the gradual dosage of oxygen plasma on the two orientations with initially  
 18 different WFs will induce a different behaviour regarding the WF evolution. It seems that the WF  
 19 of Cu(100) is prone to stabilize after the initial oxidation, while that of Cu(111) shows a gradual  
 20 change during the first stages of the plasma oxidation. This finding could be correlated with the  
 21 fact that Cu(111) shows a gradual mixing of  $Cu_2O$  and CuO in the initial oxidation steps, while  
 22 Cu(100) shows a slower oxidation in the first stages. The fact that the 180 s treatment is  
 23 characterized by a WF difference of 0.36 eV between the two crystals is therefore not completely  
 24 surprising, considering that the WF value is sensitive to the influence of various parameters such  
 25 as step density, defects, adsorbates or surface reconstruction, which is in line with the other  
 26 differences reported in the current study.

## 1 Supplementary LEED images



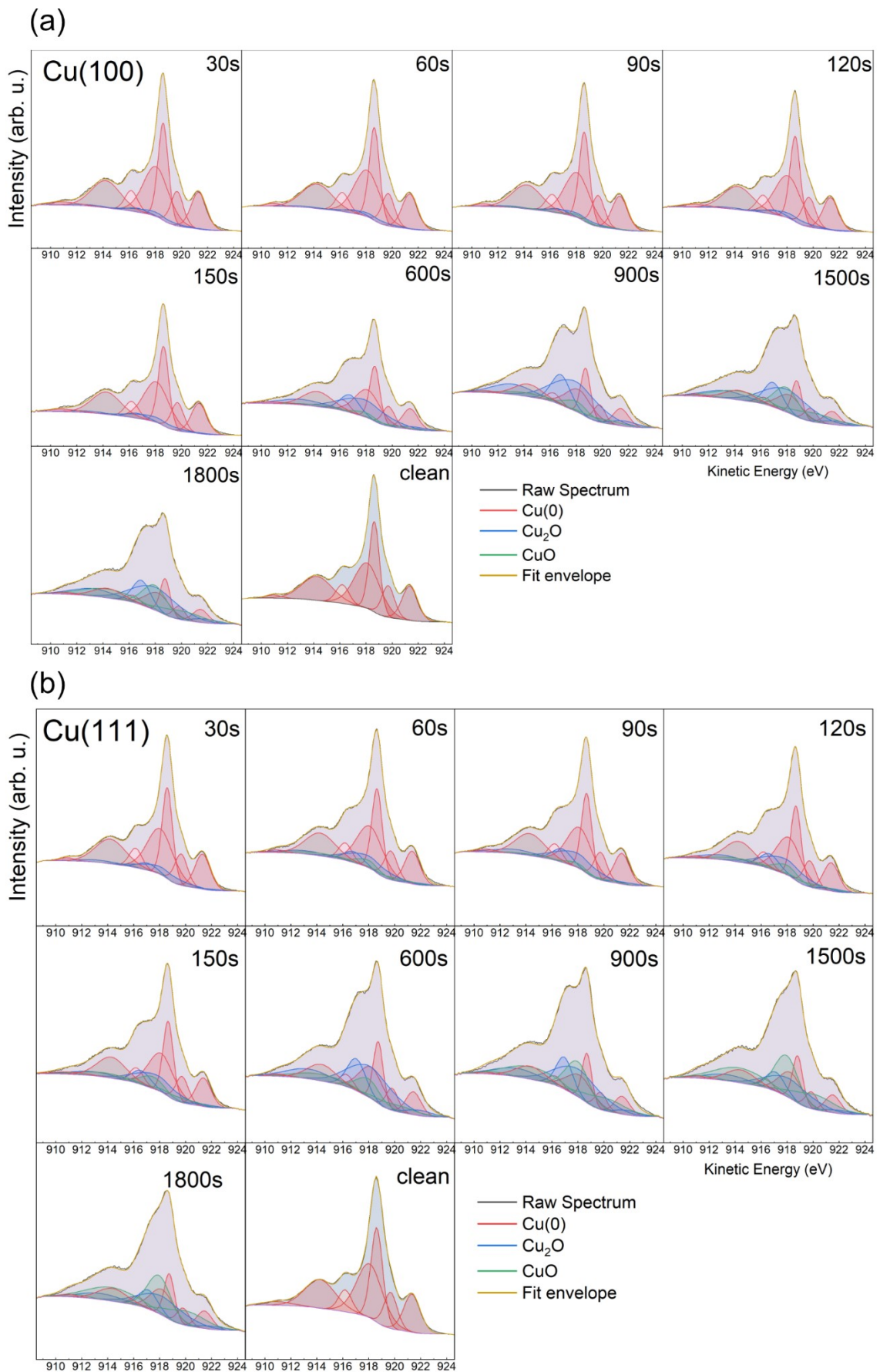
4 **Figure S3:** LEED images recorded with 42 eV on Cu(100) for the pristine sample (a), and with  
5 20 eV after 180 s treatment (b), respectively on Cu(111) in the same conditions (c-d). Fig. (b) and  
6 (d) show a better contrast on the quasi (2×2) spots on the plasma treated samples. Nevertheless,  
7 due to the low energy involved, the size of the Ewald sphere is decreased and the LEED image  
8 does not contain the fundamental spots corresponding to the (1×1) structure. For the same reason,  
9 the LEED images measured with 20 eV on the pristine crystal (not shown) display only the (00)  
10 spot. (e-f) Structural models displaying the surface growth of Cu<sub>2</sub>O(111) on top of the two Cu  
11 model surfaces; for simplicity only one layer is shown.

## 13 X-ray Photoelectron Spectroscopy (XPS)

14 In order to quantify the composition of the plasma treated surfaces, we performed peak fitting  
15 of the Cu LMM spectra on a Shirley background based on the procedure and constraints outlined  
16 in *Biesinger, M. C. (2017)*.<sup>6</sup> The obtained fit curves and Cu LMM data are shown in **Fig S4**. The  
17 binding energies were references to the most intense LMM peak corresponding to metallic Cu.  
18 We used the seven most intense peaks of metallic Cu, and four peaks for Cu<sub>2</sub>O and CuO each.  
19 The fitting parameters for the fitting model are listed in table S1. The peak shapes are modelled  
20 with mixed Gaussian-Lorentzian (GL) line shapes with a 30:70 weight. One should note that  
21 quantification of copper in different oxidation states is quite challenging, and that the overall trend  
22 of the oxidation process of the two differently-oriented surfaces is the aspect of interest here.

24

25



1 **Figure S4:** Cu LMM spectra and fitted components of (a) Cu(100) and (b) Cu(111) after a room temperature  $O_2$ -plasma treatment ( $3 \times 10^{-5}$  mbar  $O_2$ ) for the times indicated and the clean surfaces.



1 **Table S1:** Fitting parameters for the fitting model of the Cu LMM deconvolution shown in Fig.  
 2 S4. The parameters of position, position constraints ( $\Delta$  Position) in respect to the previous peak  
 3 position, full width half maximum (FWHM) and FWHM variation ( $\Delta$ FWHM ) are shown.

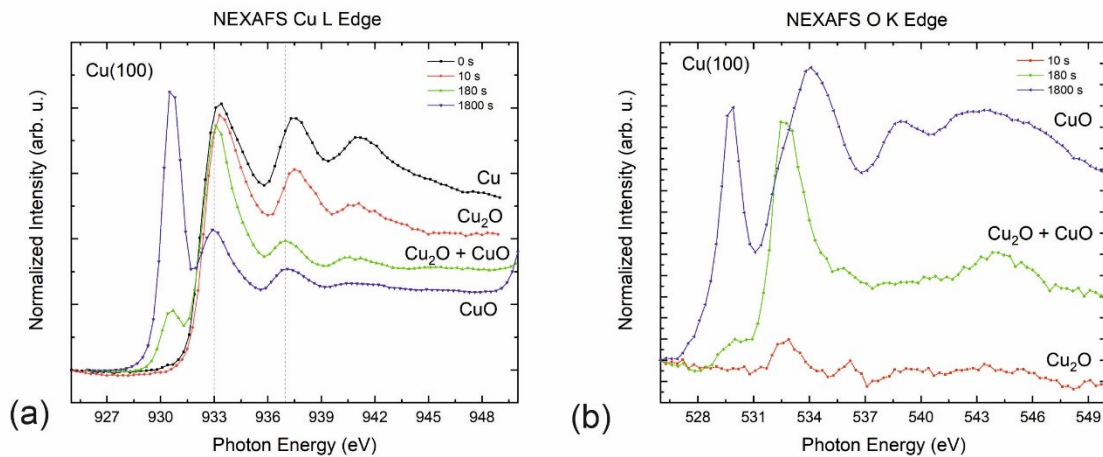
Component	Position (eV)	$\Delta$ Position (eV)	FWHM	$\Delta$ FWHM
Cu(0)	916.24		1.2	0.2
Cu(0)	921.39	5.15	1.4	0.3
Cu(0)	919.74	1.65	1.0	0.1
Cu(0)	918.68	1.06	0.9	0.25
Cu(0)	918.13	0.43	2.3	0.2
Cu(0)	914.30	3.95	2.5	0.4
Cu(0)	910.98	3.32	1.6	0.3
Cu <sub>2</sub> O	921.64		2.3	0.4
Cu <sub>2</sub> O	917.85	3.79	4.2	0.4
Cu <sub>2</sub> O	916.76	1.09	1.6	0.3
Cu <sub>2</sub> O	913.09	3.67	3.9	0.4
CuO	920.02		3.8	0.2
CuO	917.94	2.08	2.1	0.2
CuO	914.30	3.64	4.5	0.3
CuO	911.45	2.85	2.4	0.25

4

5

## 6 Cu L-edge and O K-edge NEXAFS

7



8

9 **Figure S5.** NEXAFS spectra measured at different stages of plasma treatment. (a) Cu L-edge,  
 10 respectively (b) O K-edge NEXAFS spectra of the initial metallic surface and after 10 s, 180 s and  
 11 1800 s of plasma oxidation at  $4 \times 10^{-4}$  mbar of O<sub>2</sub>, proving characteristic fingerprints of different  
 12 copper oxide species. In the case of image (b), the spectrum measured for the clean sample was  
 13 used for the normalization of the files, and is therefore not shown.

14

15

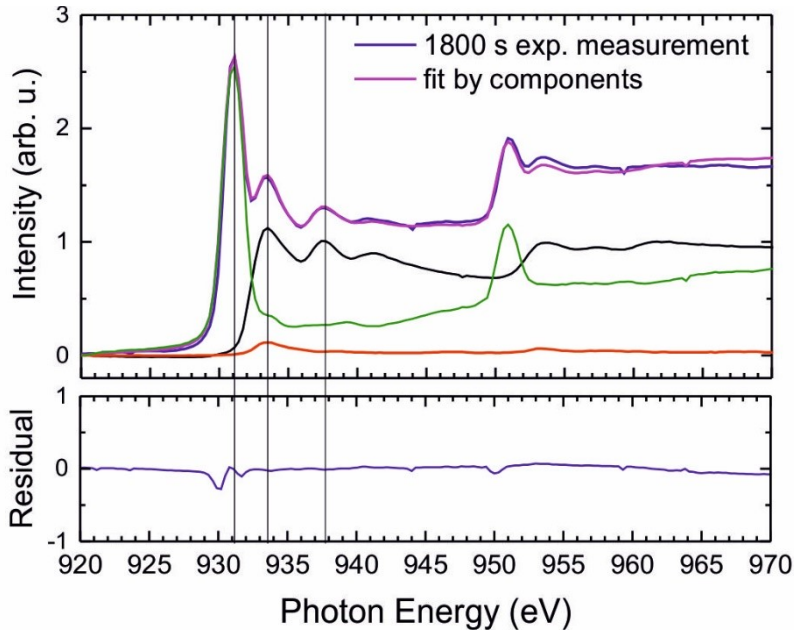
## 1 Composition analysis of NEXAFS spectra

2

3 The analysis of the Cu L-edge NEXAFS spectra implies that the initial state of the crystal is  
 4 metallic and that the further spectra are a linear combination of characteristic spectra of metallic  
 5 copper ( $\text{Cu}^0$ ),  $\text{Cu}_2\text{O}$  ( $\text{Cu}^+$ ) and  $\text{CuO}$  ( $\text{Cu}^{2+}$ ) at a certain exposure time  $t$ :

$$6 \quad I(t) = \sum_{i=1}^3 \alpha_i(t) I_i(\text{Cu}_i) = \alpha_1 I_1(\text{Cu}^0) + \alpha_2 I_2(\text{Cu}^+) + \alpha_3 I_3(\text{Cu}^{2+}) \quad (1),$$

7 where  $\alpha_i$  represents a normalized weighting factor at a certain time, such that  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ .  
 8 They correspond to the “composition” of  $\text{Cu}^0$ ,  $\text{Cu}^+$  and  $\text{Cu}^{2+}$  in the spectra shown in **Fig. 8 (a)** and  
 9 **(c)**. The individual component spectra have been extracted from the actual data, as following: (i)  
 10 the  $\text{Cu}^0$  curve corresponds to the NEXAFS spectrum recorded on the pristine crystal; (ii) based on  
 11 equation (1) the  $\text{Cu}^+$  curve is produced as a difference between the NEXAFS spectrum measured  
 12 at 30 s and the weighted pristine one, whereas the weighting factor  $\alpha_2 = 1 - \alpha_1$  is chosen in a  
 13 way that the residual oscillations at 938 eV and 941.5 eV (characteristic for metallic copper) are  
 14 cancelled out; (iii) correspondingly, the  $\text{Cu}^{2+}$  curve is produced as a difference between the  
 15 NEXAFS spectrum measured at 1800 s and the sum of weighted  $\text{Cu}^0$  and  $\text{Cu}^+$  curves with  
 16 weighting factors  $\alpha_1$ ,  $\alpha_2$  and  $\alpha_3 = 1 - \alpha_1 - \alpha_2$  in a way that the characteristic peak at 931 eV is  
 17 well described.



**Figure S6:** Example of the fit components of a Cu L-edge NEXAFS. The experimental curve (blue) is fitted by a sum of the NEXAFS components of  $\text{Cu}^0$ ,  $\text{Cu}^+$  and  $\text{Cu}^{2+}$  (black, red and green curves, respectively). The residual in the bottom exhibits the quality of the fit, whereas the emphasis is put on the range of the first three maxima. The data shows the NEXAFS measurement of the Cu(111) surface after an overall 30 min treatment with oxygen plasma.

18

19

20

## 1 **Model of Copper Oxide Formation for Cu(100)**

2

3 The model presented in the main text assumes that Cu<sub>2</sub>O grows after the initial oxidation steps  
4 until a certain thickness  $\delta$  is reached, above which the growth is continued for a  $\Delta$  thickness as  
5 CuO.

6 Starting with the assumption that each Cu atom from the depth  $z$  contributes with an intensity  
7  $I_0 e^{-\frac{z}{\Gamma}}$ , it yields the following terms for the three contribution:

$$8 \quad I_{Cu^{2+}} = n_{2+} I_0 \left(1 - e^{-\frac{\Delta}{\Gamma}}\right) \quad (2)$$

$$9 \quad I_{Cu^+} = n_+ I_0 \left(1 - e^{-\frac{\delta}{\Gamma}}\right) e^{-\frac{\Delta}{\Gamma}} \quad (3)$$

$$10 \quad I_{Cu^0} = n_0 I_0 e^{-\frac{(\delta + \Delta)}{\Gamma}} \quad (4),$$

11 where  $\Gamma$  represents the mean free path length of the electrons at the detected energy which is  
12 assumed to be the same for all three copper species.  $n_0$ ,  $n_+$  and  $n_{2+}$  are atomic concentrations  
13 of copper in the metallic, the Cu<sub>2</sub>O and CuO phase. The intensity  $I_0$  is assumed to be identical for  
14 all Cu atoms.

15 The normalization of the intensity eq. (2)-(4) with the total intensity

$$16 \quad I_{total} = I_{Cu^0} + I_{Cu^+} + I_{Cu^{2+}} \quad (5)$$

17 yields in the three normalized intensities which describe the intensity composition of **Figure 9(a)**  
18 and **(d)** and is used in formula eq. 1.

$$19 \quad \alpha_{Cu^{2+}} = \frac{I_{Cu^{2+}}}{I_{total}} \quad (6)$$

$$20 \quad \alpha_{Cu^+} = \frac{I_{Cu^+}}{I_{total}} \quad (7)$$

$$21 \quad \alpha_{Cu^0} = \frac{I_{Cu^0}}{I_{total}} \quad (8)$$

22 It is to be noted that the sum of the three components gives:

$$23 \quad \alpha_{Cu^0} + \alpha_{Cu^+} + \alpha_{Cu^{2+}} = 1 \quad (9)$$

24

## 25 **Correlation between film thickness and plasma exposure time**

26



1 The growth of the Cu<sub>2</sub>O film can be described as a linear increase between 0 to thickness  $\delta$   
2 within 30 s (at  $4 \times 10^{-4}$  mbar O<sub>2</sub>), the value above which the thickness remains constant. The CuO  
3 starts to grow on top from  $t = 30$  s on. The time dependence of the  $\Delta$  thickness of this growing  
4 CuO film might be described by two models:

- 5 (1) Simple linear model:  $\Delta(t) = R \times t$ , with  $R$  is the constant oxidation rate of CuO  
6 (2) Damping model, in which the oxidation rate is damped by the thickness of the film

7 In the both cases, the damping of the composition in **Fig. 9(a)** and **(d)** is fitted over the full  
8 exposure time of 1800 s (at  $p(\text{O}_2) = 4 \times 10^{-4}$  mbar). We make the remark that it was not yet  
9 considered that the copper layers expand in volume when the oxygen is incorporated into copper,  
10 which means that the oxide grows into the bulk, but also increases the height (see next subsection).

11 The damping model considers that the film grows with dosage according to the following rate  
12 equation:

$$\frac{d\Delta}{dt} = R e^{-\frac{\Delta}{\Lambda}}, \quad (10)$$

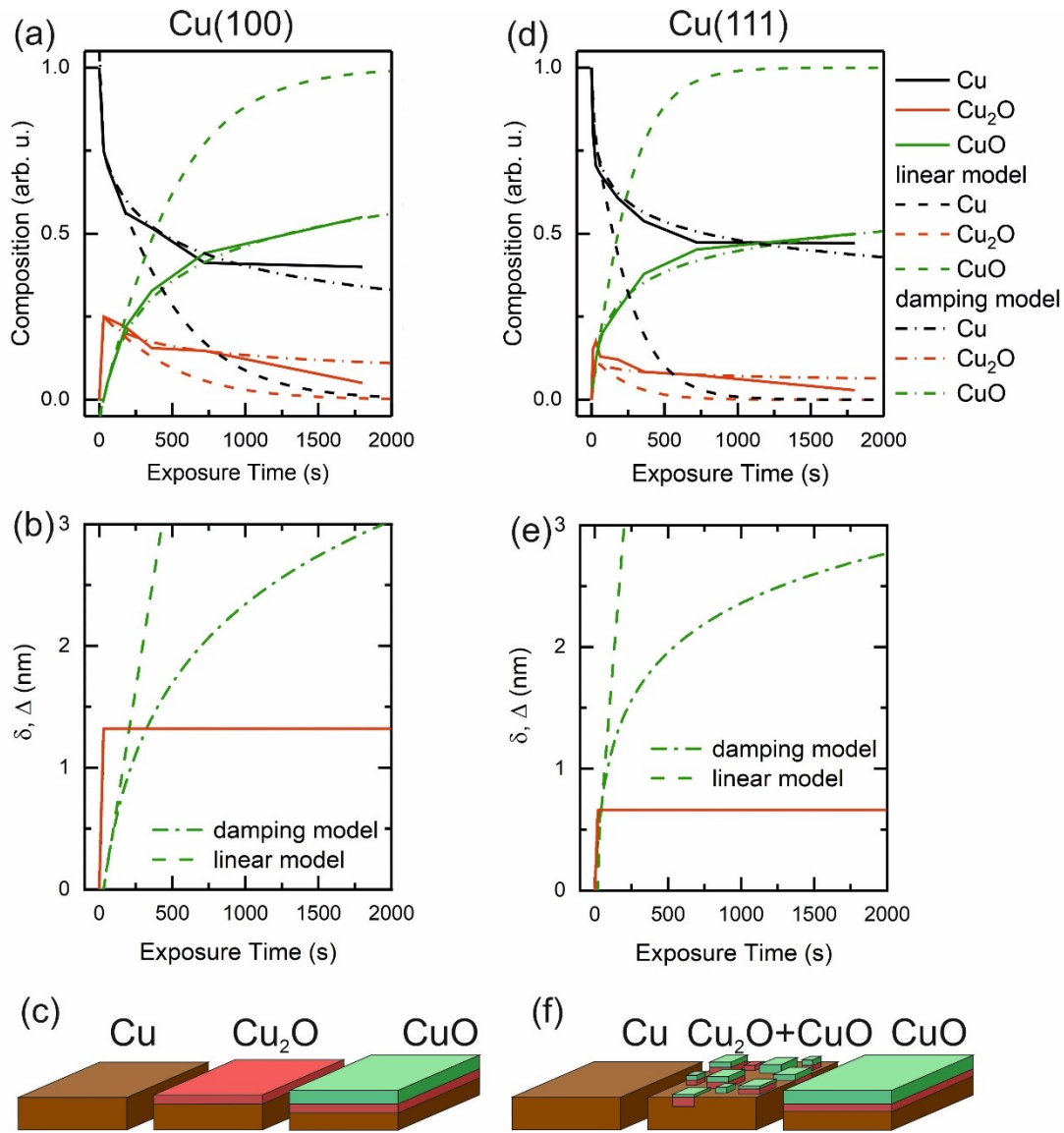
14 where  $R$  is the initial oxidation rate, *i.e.* the oxidized film thickness per time. Nevertheless, the  
15 oxidation is hindered when the film gets thicker. The damping of the process is described by the  
16 exponential factor with  $\Lambda$  being the effective oxidation length. Considering the starting condition  
17  $\Delta(t = 0 \text{ s}) = 0$  nm, *i.e.* no oxide film in the beginning, the solution of the rate equation is:

$$\Delta(t) = \Lambda \ln \left( 1 + \frac{R}{\Lambda} t \right) \quad (11)$$

19 For  $Rt \ll \Lambda$  or small values of  $t$ , this formula can be approximated as  $\Delta(t) \cong Rt$ , what is just the  
20 expected initial linear behaviour.

21 **Figure S7** shows clearly that fitting with the damping model describes the time dependence  
22 of the NEXAFS components much better than the linear model.

23



1  
2

3 **Figure S7:** (a) and (d) Fitting of time dependence of the NEXAFS components with the linear  
4 (dashed lines) and the damping models (dashed dotted lines) for plasma-treated Cu(100) and  
5 Cu(111), left and right column respectively. The plasma treatment was done at  $4 \times 10^{-4}$  mbar  $O_2$ .  
6 (b) and (e) display the thickness of the growing CuO (red) and Cu<sub>2</sub>O (green) film. A schematic  
7 model for the film growth is given in the bottom. The fitting parameters are listed in **Table S2**.

8

9 **Table S2:** Parameters extracted from the fitting of the composition evolution of Cu, Cu<sub>2</sub>O and  
10 CuO species as function of plasma exposure time at room temperature at  $4 \times 10^{-4}$  mbar of oxygen  
11 for Cu(100) and Cu(111). Fits are displayed in **Fig. S7**.

	<b>Linear model:</b> $\Delta(t) = Rt$	<b>Damping model:</b> $\Delta(t) = \Lambda \ln \left( 1 + \frac{R}{\Lambda} t \right)$
<b>Cu(100)</b>		
Mean free path length <sup>7,8</sup> is set to	$\Gamma = 3$ nm	$\Gamma = 3$ nm

Thickness of buried Cu <sub>2</sub> O layer	$\delta = 1.3 \text{ nm}$	$\delta = 1.3 \text{ nm}$
Initial oxidation rate of CuO	$R = 0.0075 \text{ nm/s}$	$R = 0.009 \text{ nm/s}$
Effective oxidation length (damping)	N.A.; $\Lambda = \infty$	$\Lambda = 1.05 \text{ nm}$
<b>Cu(111)</b>		
Mean free path length is set to	$\Gamma = 3 \text{ nm}$	$\Gamma = 3 \text{ nm}$
Thickness of buried Cu <sub>2</sub> O layer	$d = 0.66 \text{ nm}$	$d = 0.66 \text{ nm}$
Initial oxidation rate of CuO	$R = 0.015 \text{ nm/s}$	$R = 0.03 \text{ nm/s}$
Effective oxidation length (damping)	N.A.; $\Lambda = \infty$	$\Lambda = 0.6 \text{ nm}$

1

2 In the literature, a wide range of IMFP values can be found depending on the material and also  
3 on the experimental or theoretical model. Therefore, for simplicity, we set one value for the IMFP,  
4 independently of the material, but from the kinetic energy. Based on this, we assume an absolute  
5 accuracy better than a factor 2 for the film thicknesses and the related growth rates. However, the  
6 relative difference between Cu(100) and Cu(111) is not affected by this because the same model  
7 is applied in both cases.

## 8 Position of the Oxide Film Boundaries

9

10 When the copper gets oxidized the atomic concentration decreases from  $n_0 = 84.232 \text{ nm}^{-3}$   
11 down to  $n_+ = 50.740 \text{ nm}^{-3}$  and  $n_{2+} = 52.731 \text{ nm}^{-3}$  for Cu<sub>2</sub>O and CuO respectively (see  
12 **Table S4**). Because the amount of copper atoms is conserved, the crystal has to grow in height  
13 during oxidation (see **Figure 9 (c) and (f)**). If the height of the initial copper surface is defined as  
14  $z = 0$ , then the position of the boundary  $z_{Cu^0 \parallel Cu^+}$  between metallic Cu and the Cu<sub>2</sub>O, the  
15 boundary  $z_{Cu^+ \parallel Cu^{2+}}$  between Cu<sub>2</sub>O and CuO and the top level of the oxide film (boundary to  
16 vacuum)  $z_{top}$

$$17 \quad z_{Cu^0 \parallel Cu^+} = - \left( \frac{n_+}{n_0} \delta + \frac{n_{2+}}{n_0} \Delta \right) \quad (12)$$

$$18 \quad z_{Cu^+ \parallel Cu^{2+}} = z_{Cu^0 \parallel Cu^+} + \delta = \left( \left( 1 - \frac{n_+}{n_0} \right) \delta - \frac{n_{2+}}{n_0} \Delta \right) \quad (13)$$

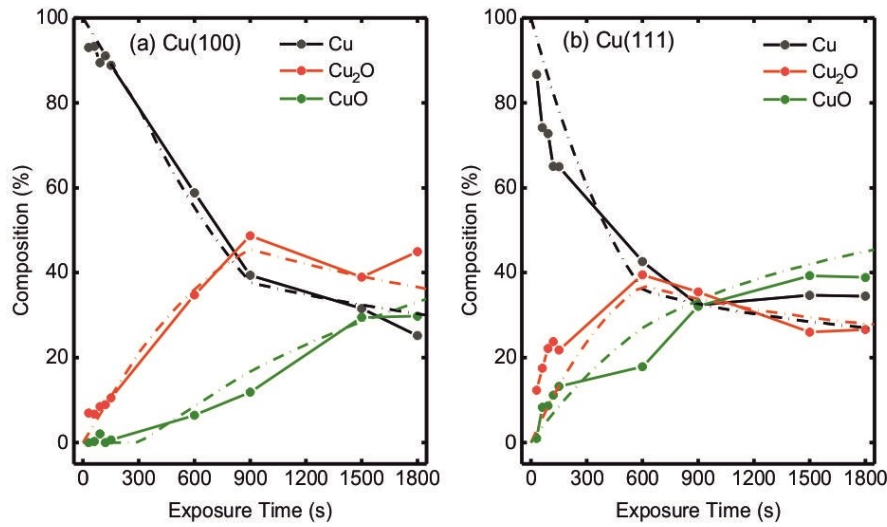
$$19 \quad z_{top} = z_{Cu^0 \parallel Cu^+} + \delta + \Delta = \left( \left( 1 - \frac{n_+}{n_0} \right) \delta + \left( 1 - \frac{n_{2+}}{n_0} \right) \Delta \right) \quad (14)$$

1 Whereas the bottom level  $z_{Cu^0 \parallel Cu^+}$  of the oxide film is always below the level of the initial  
 2 surface (i.e.  $z_{Cu^0 \parallel Cu^+} < 0$ ), the top level  $z_{top}$  is always above the initial surface level (i.e.  
 3  $z_{top} > 0$ ). The boundary between  $Cu_2O$  and  $CuO$  is first above, and at further oxidation below the  
 4 initial surface level.

5

## 6 Composition fit of XPS data

7 Analogous to the fit of the NEXAFS data, the XPS composition data were fitted with the damping  
 8 model. **Fig. S7** shows the experimental data with the model fit.



9

10 **Figure S8:** XPS data fit with the same model applied for the NEXAFS data. The plasma treatment  
 11 was done at  $3 \times 10^{-5}$  mbar  $O_2$ . The fitting parameters are given in **Table S3**.

12

13 **Table S3:** Parameters extracted from the fitting of the composition evolution of Cu,  $Cu_2O$  and  
 14  $CuO$  species as function of exposure time at  $3 \times 10^{-5}$  mbar of oxygen for Cu(100) and Cu(111). Fits  
 15 are displayed in **Fig. S8**.

	<b>Damping model:</b> $\Delta(t) = \Lambda \ln \left( 1 + \frac{R}{\Lambda} t \right)$
<b>Cu(100)</b>	
Mean free path length is set to	$\Gamma = 1.55 \text{ nm}$
Thickness of buried $Cu_2O$ layer	$\delta = 1.7 \text{ nm}$
Initial oxidation rate of $CuO$	$R = 0.0006 \text{ nm/s}$
Effective oxidation length (damping)	$\Lambda = 1.36 \text{ nm}$
<b>Cu(111)</b>	

Mean free path length <sup>9</sup> is set to	$\Gamma = 1.55$ nm
Thickness of buried Cu <sub>2</sub> O layer	$\delta = 1.55$ nm
Initial oxidation rate of CuO	$R = 0.0016$ nm
Effective oxidation length (damping)	$\Lambda = 0.6$ nm

1 Considering that the results presented in Table S2 and Table S3 have been determined from data  
2 measured in two different systems, one can calculate a rough calibration of the two plasma sources  
3 involved in the study. The initial oxidation rates of CuO, applying the damping model equal  
4  $R = 0.009$  nm/s for Cu(100), and  $R = 0.03$  nm/s for Cu(111) for the LEEM/XPEEM and  
5  $R = 0.0006$  nm/s for Cu(100) and  $R = 0.0016$  nm/s for Cu(111) for the STM/XPS. This yields a  
6 higher oxidation rate in the LEEM/XPEEM system (about 20 times higher) as compared to the  
7 STM/XPS system. This factor can be very well explained by the pressure ratio of a factor 13,  
8 together with the different sample to source distances and slight differences in the individual  
9 sources. Additionally, the different pumping speeds of the systems might play a role, resulting in  
10 a different pressure in the plasma chamber, though the nominal measured pressure in the vacuum  
11 chamber is nearly the same.

## 12 Structure

13

14 **Table S4:** Comparison of the crystal parameters for the metallic Cu crystal, cuprite Cu<sub>2</sub>O and  
15 tenorite CuO.

	Space group	Lattice parameter (Å)	Volume unit cell (Å <sup>3</sup> )	$n_{\text{Cu}}$ Cu density (1/nm <sup>3</sup> )	$a_{\text{NN}}$ of Cu (Å)	d row distance on (111) (Å)	Ratio to $a_{\text{NN}}$ (Cu)	reference
<b>Cu</b>	Fm-3m [225] Cubic	$a=b=c=3.621$ Å $\alpha=\beta=\gamma=90^\circ$ fcc, 4 Cu atoms per unit cell	47.488	84.232	2.561	2.218	1	<a href="https://materialsproject.org/materials/mp-30/#">https://materialsproject.org/materials/mp-30/#</a> DOI: 10.17188/1204433
<b>Cu<sub>2</sub>O</b>	Pn-3m [224] Cubic	$a=b=c=4.288$ Å $\alpha=\beta=\gamma=90^\circ$ 4 Cu atoms on fcc position 2 O atoms on bcc position	78.833	50.740	3.032	2.626	1.184	<a href="https://materialsproject.org/materials/mp-361/">https://materialsproject.org/materials/mp-361/</a> DOI: 10.17188/1207131
<b>CuO</b>	Fm-3m [225] cubic	$a=b=c=2.993$ Å $\alpha=\beta=\gamma=60^\circ$	18.964	52.731	2.993	2.592	1.169	<a href="https://materialsproject.org/materials/mp-14549/">https://materialsproject.org/materials/mp-14549/</a> DOI: 10.17188/1190720

16

17

18

1  
2  
3  
4  
5  
6  
7  
8  
9

## 10 **References**

11

- 12 1. C. Schmidt, A. Witt and G. Witte, *J Phys Chem*, 2011, **115**, 7234-7241.
- 13 2. B. Cook, A. Russakoff and K. Varga, *Appl Phys Lett*, 2015, **106**, 211601.
- 14 3. J. I. Flege, A. Meyer, J. Falta and E. E. Krasovskii, *Physical Review B*, 2011, **84**, 115441.
- 15 4. J. I. Flege and E. E. Krasovskii, *physica status solidi (RRL) - Rapid Research Letters*, 2014, **8**, 463-  
16 477.
- 17 5. J. Deuermeier, H. J. Liu, L. Rapenne, T. Calmeiro, G. Renou, R. Martins, D. Munoz-Rojas and E.  
18 Fortunato, *Apl Materials*, 2018, **6**, 096103.
- 19 6. M. C. Biesinger, 2017, **49**, 1325-1334.
- 20 7. H. T. Nguyen-Truong, *J Phys Condens Matter*, 2017, **29**, 215501.
- 21 8. O. Y. Ridzel, V. Astašauskas and W. S. M. Werner, *Journal of Electron Spectroscopy and Related  
22 Phenomena*, 2020, **241**, 146824.
- 23 9. M. P. Seah and W. A. Dench, *Surface and Interface Analysis*, 1979, **1**, 2-11.

24