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## Electronic Supplementary Information for

## Intrinsic Auxeticity and Negative Piezoelectricity in Two-dimensional Group-IV dipnictide monolayers with In-plane Anisotropy <sup>†</sup>

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Figure S1: Crystal structure for 3D bulk SiP in orthorhombic  $Cmc2_1$  symmetry. Unit cell of bulk SiP is indicated by dashed lines. Si and P atoms are marked in blue and red colors, respectively.

		a (Å)	b (Å)	c (Å)	Symmetry	$\beta$
GeP	Cal	15.16	3.62	9.15	C2/m	$101.1^{\circ}$
	$\operatorname{Exp}$	15.44	3.64	9.19	C2/m	$101.1^{\circ}$
GeAs	Cal	15.53	3.78	9.45	C2/m	$100.7^{\circ}$
	$\operatorname{Exp}$	15.59	3.79	9.49	C2/m	$101.3^{\circ}$
SiAs	Cal	16.04	3.66	9.53	C2/m	$106.1^{\circ}$
	$\operatorname{Exp}$	15.98	3.67	9.53	C2/m	$106.0^{\circ}$
SiP	Cal	13.69	3.51	20.46	$Cmc2_1$	90°
	$\operatorname{Exp}$	13.96	3.54	20.85	$Cmc2_1$	90°
$GeP_2$	Cal	14.33	3.49	9.91	Pbam	
	$\operatorname{Exp}$	NA	NA	NA		
$\operatorname{GeAs}_2$	Cal	14.75	3.72	10.13	Pbam	
	$\operatorname{Exp}$	14.76	3.73	10.16	Pbam	
$\mathrm{SiP}_2$	Cal	14.38	3.42	9.84	Pbam	
	$\operatorname{Exp}$	13.97	3.44	10.08	Pbam	
$SiAs_2$	Cal	14.88	3.65	10.07	Pbam	
	$\operatorname{Exp}$	14.53	3.64	10.37	Pbam	

Table S1: Comparison of our calculated (Cal) lattice parameters with experimental  $(\text{Exp})^{1-6}$  results for all AB and  $AB_2$  (A = Si and Ge, B = P and As) bulk compounds.

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Table S2: Comparison of our optimized the planar lattice parameters with the results reported by other work for 2D AB and  $AB_2$  monolayers with orthorhombic lattice.

		$h(\dot{\lambda})$	c (Å)	Summetry
		0 (A)		Symmetry
$\mathrm{GeP}$	Ours	3.59	21.07	C2/m
	Others	$3.648^{7}$	$21.393^{7}$	C2/m
GeAs	Ours	3.73	21.75	C2/m
	Others	$3.803,^7 \ 3.78^8$	$22.129,^7 22.07^8$	C2/m
$\operatorname{SiP}$	Ours	3.50	20.35	C2/m
	Others	$3.519,^7 \ 3.521^9$	$20.449,^7 \ 20.479^9$	C2/m
SiAs	Ours	3.65	21.08	C2/m
	Others	$3.682,^7 \ 3.66^8$	$21.246,^{7} 21.19^{8}$	C2/m
$\mathrm{GeP}_2$	Ours	3.48	10.15	$Pmc2_1$
	Others	$3.53^{10}$ $3.522^{11}$	$10.51, {}^{10}10.205^{11}$	$Pmc2_1$
$\operatorname{GeAs}_2$	Ours	3.69	10.32	$Pmc2_1$
	Others	$3.761^{11}$	$10.398^{11}$	$Pmc2_1$
$\mathrm{SiP}_2$	Ours	3.42	10.04	$Pmc2_1$
	Others	$3.442^{11}$	$10.045^{11}$	$Pmc2_1$
$SiAs_2$	Ours	3.63	10.29	$Pmc2_1$
	Others	$3.674^{11}$	$10.300^{11}$	$Pmc2_1$

Table S3: The summary of the predicted in-plane NPR for single-layer (SL)  $PdSe_2$ ,<sup>12</sup> SL  $GaPS_4$ ,<sup>13</sup> pentagraphene,<sup>14</sup> SL  $Ag_2S$ ,<sup>15</sup> SL  $Mo_2C$ ,<sup>16</sup>  $Be_5C_2$ ,<sup>17</sup> penta- $B_4N_2$ ,<sup>18</sup> SL  $Se_3P_2$ ,<sup>19</sup> 2D VF<sub>4</sub>,<sup>20</sup>  $\delta$ -phosphorene,<sup>21</sup>  $\delta$  -V<sup>A</sup>N (V<sup>A</sup> = P, As, Sb, Bi),<sup>22</sup> SL SiAs<sub>2</sub>, W<sub>2</sub>C<sup>23</sup> and TcTe<sub>2</sub>.<sup>24</sup> Those 2D materials synthesized either in 3D vdW bulk or 2D layered structural forms in experiment are marked by \*. SiAs<sub>2</sub> monolayer characterizes the in-plane NPR of largest magnitude among all synthetic 2D auxetic materials.

2D materials	In-plane NPR
$SL PdSe_2^*$	-0.022
$SL GaPS_4^*$	-0.033
Penta-graphene	-0.068
$SL Ag_2S$	-0.12
$SL Mo_2C^*$	-0.15
$\mathrm{Be}_5\mathrm{C}_2$	-0.16
$SL Se_3P_2$	-0.199
$2D VF_4^*$	-0.26
$\delta$ -Phosphorene	-0.267
$\delta$ -PN	-0.268
$\delta$ -AsN	-0.177
$\delta ext{-SbN}$	-0.296
$\delta$ -BiN	-0.260
$SL SiAs_2^*$	-0.32
$TcTe_2$	-0.37
$W_2C$	-0.43

Table S4: Comparison of our calculated energy band gaps ( $E_g$  in eV) and anisotropic carrier effective masses  $m_e^*/m_0$  and  $m_h^*/m_0$  using HSE06 functional with the results reported by other work for all AB and  $AB_2$  monolayers. The simulated direct and indirect  $E_g$  are marked by subscript indices  $_D$  and  $_I$ , respectively.

		$m_e^*$	$/m_{0}$	$\mathrm{m}_h^*$	$/m_{0}$	$E_g$
		Г-Х	Г-Ү	Г-Х	Г-Ү	
GeP	Ours	0.36	0.72	0.93	0.56	$2.48_{I}$
GeAs	Others Ours Others	$0.053, 0.40^{-5}$ 0.13 0.118, 7, 0.1228, 8, 0.1226	$0.004, 0.72^{-5}$ 0.9 0.004, 7, 0.0178, 8, 0.2526	$0.175, 0.98^{-5}$ 0.78 0.152, 7, 0.6221, 8, 0.8426	$0.004, 0.57^{20}$ 0.22 $0.002, 7, 0.0068, 8, 0.24^{26}$	$2.309_{I}$ , $2.22_{D}$ , 2.077, $7.2.06$ , $8$
SiP	Ours Others	0.113, 0.1328, 0.12 0.14 $0.111, 7, 0.24^9$	0.004, 0.0178, 0.25 0.57 $0.062, 7, 0.70^9$	0.152, 0.0221, 0.04 1.06 $0.351, 7, 0.42^9$	0.002, 0.0008, 0.24 0.73 0.050 7 $0.44$ 9	$2.64_I(2.65_D)$ $2.64_I = 7.250 p^9$
SiAs	Ours Others	$\begin{array}{c} 0.111, & 0.24 \\ 0.11 \\ 0.050 & 7 & 0.1025 & 8 & 0.1426 \end{array}$	0.002, 0.10 0.37 $0.003$ <sup>7</sup> $0.0195$ <sup>8</sup> $0.88^{26}$	0.331, 0.42 0.84 $0.137, 0.6975, 8, 0.94^{26}$	$\begin{array}{c} 0.030, & 0.44\\ 0.22\\ 0.002 & 7 & 0.0062 & 0.22^{26} \end{array}$	2.041D, 2.05D 2.43D $2.353D, 7, 2.50D^8$
$\mathrm{GeP}_2$	Ours Others	$\begin{array}{c} 0.1020, & 0.1020, & 0.14 \\ 0.13 \\ 0.14^{10} \end{array}$	1.23 $1.20^{10}$	1.39 $1.54^{10}$	$\begin{array}{c} 0.002, & 0.0002, & 0.22\\ 0.32\\ 0.38^{10} \end{array}$	$2.08_I(2.10_D)$ $2.01_I 101_{10} 06_D 11$
$\operatorname{GeAs}_2$	Ours Others	0.13 NA	0.56 NA	0.43 NA	0.63 NA	$1.78_I$ $1.64_r$ <sup>11</sup>
$\mathrm{SiP}_2$	Ours Others	0.12 NA	2.10 NA	3.99 N 4	1.12 N 4	$2.25_I(2.28_D)$ $2.23_D^{-11}$
$\operatorname{SiAs}_2$	Ours Others	0.13 NA	0.41 NA	0.85 NA	2.20 NA	$2.04_{I}$ $1.92_{I}$ <sup>11</sup>



Figure S2: The evolution of cation polar displacement relative to anion tetrahedral center of  $AB_2$  monolayers, as a function of uniaxial strain applied along the polar c axis. Four tetrahedrons (marked as  $T_1$ ,  $T_2$ ,  $T_3$ , and  $T_4$ ,) within  $AB_2$  monolayers can yield two distinct polar displacement. Reducing (increasing) the magnitude of polar displacement along the polar c axis by applying tensile (compressive) uniaxial strain will lead to the overall negative piezoelectric coefficients  $e_{33}^{2D}$  ( $d_{33}^{2D}$ ) for  $AB_2$  monolayers.



Figure S3: The HSE06 calculations predicted photon energy dependent optical absorption coefficients  $\alpha_x$  and  $\alpha_y$  for all AB and  $AB_2$  monolayers under incident light polarized along crystallographic b and c axes, respectively.

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