# Extreme-Modulation of Liquid Crystal Viscoelasticity

# via Altering Ester Bond Direction

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#### 1. Dielectric spectroscopy

The dielectric spectroscopy was made by using a multichannel electrochemical workstation (VMP3 Potentiostat, BioLogic). As shown in Fig. S9, a large dielectric relaxation was observed for COO<sub>4</sub> and OCO<sub>4</sub> molecules (also for all the other homologues) with the relaxation strength,  $\delta\varepsilon$ , of over 100. Similar relaxation process, but with small  $\delta\varepsilon$ , was observed previously, and assigned to a goldstone mode that corresponds to the long-scale fluctuations of the azimuthal angle of the local nematic director along the coarse-grained the coarse grained director.<sup>S1-S3</sup> The correlation length of the precession fluctuation can be estimated by using a non-stationary solution of the Smoluchowski equation.<sup>S1</sup> The correlation length is written as  $\delta z^2 \approx D/q^2 \omega$ . *D*, *q* and  $\omega$  are the orientational di usion coe cient, wavenumber of the heliconical structure and the frequency. We used  $D=10^9$  s<sup>-1</sup>, *q* from Ref. S4 and  $\omega$  from the measured dielectric spectrum (e.g. Fig. S9) for calculating the correlation length  $\delta z$ . Fig. S10 demonstrates COO<sub>4</sub> exhibits considerably longer correlations than others. With increasing the number of the carbon spacer in COO<sub>n</sub>, the correlation length decreases to nearly similar to that of OCO<sub>n</sub>.

## 2. Switching time tuning by mixing

To seek the possibility for tuning the viscoelastic properties in a wide range of values, we studied the LCD electro-optical characteristics by mixing  $COO_4$  and  $OCO_4$  molecules. We confirmed the solubility of  $COO_4$  in  $OCO_4$  very nice independent on the weight ratio. The phase diagram is shown in Fig. 3(a). Strikingly, the increasing of the weight of  $COO_4$  in  $OCO_4$  allows us for extremely modulate the switching time, continuously going from c.a. 0.1 s (pure  $OCO_4$ ) up to several hundreds of seconds (pure  $COO_4$ ). Fig. S11 represents the temperature dependence of offswitching time in various mixtures with different mixing ratio of  $COO_4$  by weight.

#### 3. Switching time for $COO_n$ and $OCO_n$ homologs (n = 8 and 10)

As discussed in the manuscript, decreasing the carbon number in the spacer causes the slow-down of their dynamics, suggesting the stronger molecular packing at longer flexible spacers. Here we show the temperature dependence of off-switching time in  $COO_n$  and  $OCO_n$  (n = 8 and 10) in Fig. S12. As shown in the figure, a decrease in the off-switching time was observed in COO homologs with increasing carbon number of spaces. Their off-switching time is in the order of hundred seconds, which is still one order longer than that in  $OCO_n$ . Based on the Arrenius plots, the rotation activation energy for  $COO_{10}$  (0.52ev) is about two times of that in  $OCO_{10}$ . (0.34ev) while the rotation activation energy for  $COO_6$  (1.61ev) is around four times magnitude for  $OCO_6$  (0.42ev). It is revealed that the activation energy of  $COO_n$  (n = 8 and 10) is comparable with thoes in  $OCO_n$  (n = 4 and 6). Thereby, reducing the carbon number in the spacer results in the decrease of the activation energy in the  $COO_n$ . This result is well consistent with the dielectric measurements.



### 4. Supporting figures

**Fig. S1.** Summary of DFT calculation at B3LYP-G3/6-311+G(d,p) basis for COO<sub>4</sub> and OCO<sub>4</sub> molecules.



**Fig. S2.** The rotational potentials about Ph–C or Ph–O bonds (represented by dotted circles in the upper space filling representations for  $COO_4$  and  $OCO_4$ , respectively) as a function of the rotation angle.



**Fig. S3.** The rotational potentials about C–O or O–C bonds (represented by dotted circles in the upper space filling representations for  $COO_4$  and  $OCO_4$ , respectively) as a function of the rotation angle.



**Fig. S4.** The rotational potentials about O–C or C–C bonds (represented by dotted circles in the upper space filling representations for  $COO_4$  and  $OCO_4$ , respectively) as a function of the rotation angle.



**Fig. S5.** The rotational potentials about -SC bonds (represented by dotted circles in the upper space filling representations for  $COO_4$  and  $OCO_4$ ) as a function of the rotation angle.



Fig. S7. (a-f) Stripe textures of  $COO_4$  and  $OCO_4$  molecules are presented at the different thicknesses of the samples. (g) The temperature dependence of the stripe periodicity is shown as a function of temperature.



Fig. S8. Temperature dependencies of the splay elastic constant of the N phase in the four molecules,  $COO_4$ ,  $OCO_4$ ,  $COO_6$  and  $OCO_6$ .



**Fig. S9.** Dielectric spectroscopy of  $COO_4$  and  $OCO_4$ : (a) The real and imaginary dielectric permittivity. The full squares show the real dielectric permittivity data. The solid lines show the imaginary dielectric permittivity data.; (b) The temperature dependence of the relaxation frequency of the slowest relaxation mode, m4.



Fig. S10. The correlation lengths of the fluctuation of the azimuthal angle of the local director for  $COO_4$ ,  $OCO_4$ ,  $COO_6$  and  $OCO_6$  as a function of a scaled temperature.



 $T/T_{NI}$ **Fig. S11.** The temperature dependencies of the off-switching time for COO<sub>4</sub>-OCO<sub>4</sub> mixtures at different weight ratios of COO<sub>4</sub>.



Fig. S12. The temperature dependencies of the off-switching time for  $COO_n$  and  $OCO_n$  (n = 8 and 10). The inset demonstrates the corresponding Arrhenius plot.  $T_{NI}$  is the transition temperature between the N and the isotropic liquid state.

## References

- S1. K. Merkel, A. Kocot, C. Welch and G. H. Mehl, Soft Modes of the Dielectric Response in the Twist-Bend Nematic Phase and Identification of the Transition to a Nematic Splay Bend Phase in the CBC7CB Dimer, *Phys. Chem. Chem. Phys.*, 2019, 21(41), 22839–22848.
- S2. K. Merkel, A. Kocot, J. K., Vij, G. Shanker, Distortions in Structures of the Twist Bend Nematic Phase of a Bent-Core Liquid Crystal by the Electric Field, *Phys. Rev. E*, 2018, 98(2), 1–8.
- S3. K. Merkel, C. Welch, Z. Ahmed, W. Piecek and G. H. Mehl, Dielectric Response of Electric-Field Distortions of the Twist-Bend Nematic Phase for LC Dimers, J. Chem. Phys., 2019, 151(11), 114908.
- S4. Y. Arakawa, K. Komatsu, J. Feng, C. Zhu and H. Tsuji, Distinct Twist-Bend Nematic Phase Behaviors Associated with the Ester-Linkage Direction of Thioether-Linked Liquid Crystal Dimers, *Mater. Adv.*, 2021, 2(1), 261–272.