

Supporting Information

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Section 1: The characterizations of device

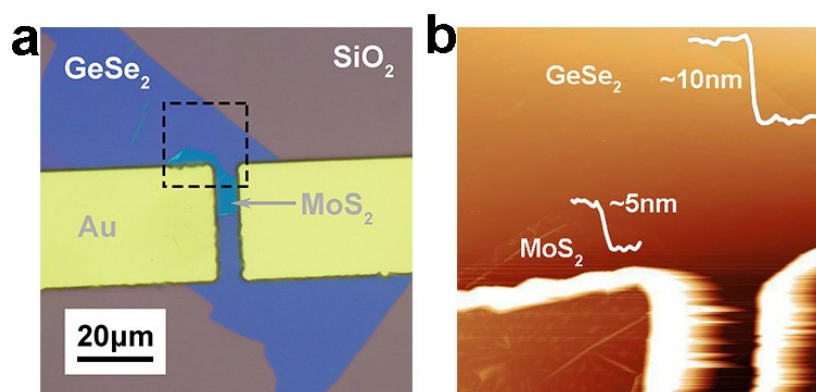


Figure S1. (a) Optical image of the device based on MoS₂/GeSe₂ van der Waals heterojunction. (b) Atomic force microscope image of the area marked with dashed rectangle in (a). From the scan lines (white), the thickness of MoS₂ and GeSe₂ are 5 nm and 10 nm respectively.

Section 2: Calculation of time constants

Here, we used a simple model¹ to estimate the time scale of charge transfer under the perpendicular gate voltage. When the back gate is switched from 0 V to a strong enough positive gate (for example, +80 V) directly, the band of GeSe₂ is tilted and a triangular barrier is formed. At this moment, the electrons in the MoS₂ channel is

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possible to transfer over the barrier under the effect of gate induced electric field. The change of 2D carrier density n is deduced by the equation below

$$dn = T n_{3D} v dt \quad (1)$$

where T is the transmission coefficient of electrons over the barrier, n_{3D} is the 3D density of electrons in the MoS₂ channel and v is the velocity of electrons. The relationship between n_{3D} and n is that $n_{3D} = n/D$, where D is the thickness of MoS₂ channel. Therefore, the time constant of this exponential decay could be obtained to be

$$\tau = \frac{D}{v T} \quad (2)$$

The velocity v could be obtained by $v = \mu E$, where μ is the mobility and E is the gate induced electric field. First, the mobility is calculated from the transfer curve with the formula below

$$\mu = \frac{dI_{ds} L}{dV_g W V_{ds} C} \quad (3)$$

Where L and W are the length and width of the device, C is the total capacitance of GeSe₂ and SiO₂ layer in parallel. The mobility is obtained to be 18.4 cm²V⁻¹s⁻¹ according to the data in the Figure 2a in the main text. The electric field E in the MoS₂ channel is calculated with the relation $E \varepsilon = E_I \varepsilon_I$, where the electric field E_I in the GeSe₂ layer under the gate voltage V_g can be obtained using the Eq. (1) in the main text and the dielectric constants of MoS₂ and GeSe₂ are $\varepsilon = 12.8$ and $\varepsilon_I = 7.0$.

For the transmission coefficient T of electrons over the barrier, calculation based on thermal emission or Fowler-Nordheim tunneling is performed. For thermal emission, there is

$$T_{TE} = e^{-\frac{U}{k_B T}} \quad (4)$$

Where the U is the height of barrier seen by electrons and equals to the conduction band offset here, k_B is the Boltzmann constant and T is the temperature. So $T_{TE} = 9.27 \times 10^{-16}$ at the room temperature $T = 300\text{K}$. For Fowler-Nordheim tunneling, the transmission coefficient could be found²

$$T_{FN} = \exp \left(-\frac{4}{3} \sqrt{\frac{2m^*}{(h/2\pi)^2}} \frac{(U - \varepsilon)^{3/2}}{eE_1} \right) \quad (5)$$

Where U is the barrier height, $m^*=0.6m_0$ ³ is the effective mass of electron, h is the Plank constant, E_1 is the electric field of tunneling layer (GeSe₂) and ε is the energy of electrons which is much less than the barrier height here. In the case of $V_g = +80$ V, $T_{FN} = 6.27 \times 10^{-14}$. Finally, we obtained the estimate of the time constant at $V_g = +80$ V

$$\tau_{TE} = 37s \quad (6)$$

And

$$\tau_{FN} = 0.5s \quad (7)$$

Section 3: The performance of the memory device

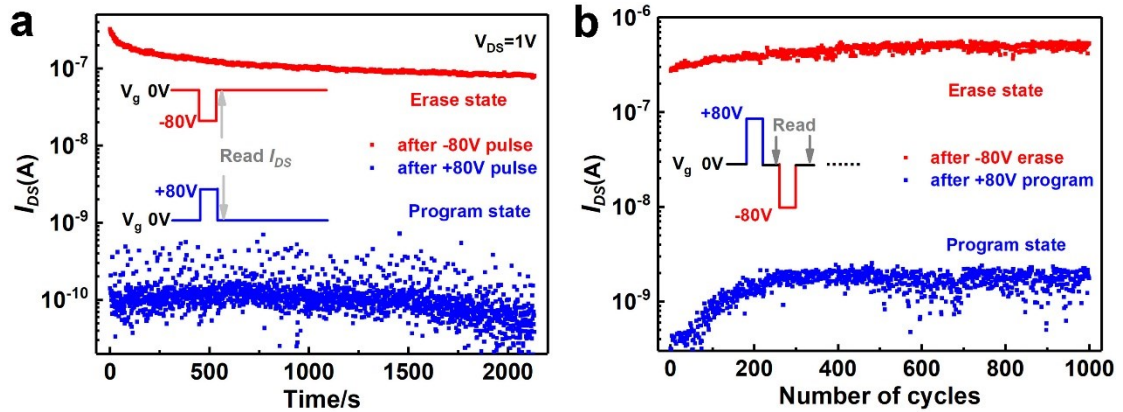


Figure S2. (a) The stability of the two-resistance state. The gate voltage V_g is biased to +80 V (-80 V) first and then reset to 0 V, the time dependence of current in the two states is recorded. (b) Cycle robustness of the device. The V_g is biased to +80 V first so that the device is set to the program state, and after the V_g is reset to 0 V the current is read to confirm the resistance. Then, the V_g is biased to -80 V to set the device to erase state, the current is read again to confirm this resistance after V_g returns to 0 V. The cycle is repeated for thousands of times to test the robustness.

References

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