

## Supporting Information

### **High-performance short-wave infrared phototransistor based on 2D tellurium/MoS<sub>2</sub> van der Waals heterojunction**

Jinrong Yao<sup>a</sup>, Fangfang Chen<sup>a</sup>, Juanjuan Li<sup>a</sup>, Junli Du<sup>b</sup>, Di Wu<sup>a</sup>, Yongtao Tian<sup>a</sup>, Cheng Zhang<sup>c</sup>, Jinke Yang<sup>d</sup>, Xinjian Li<sup>a,\*</sup>, Pei Lin<sup>a,\*</sup>

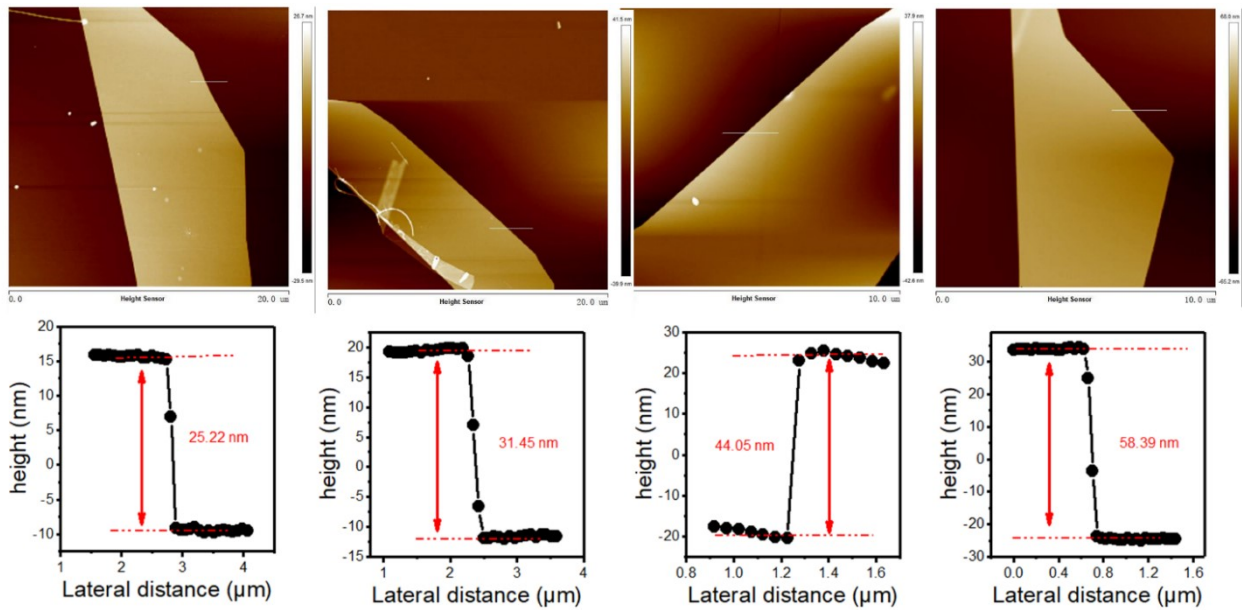
<sup>a</sup> *Key Laboratory of Materials Physics of Ministry of Education, School of Physics and Microelectronics, Zhengzhou University, Zhengzhou 450001, People's Republic of China*

<sup>b</sup> *State Grid Henan Electric Power Research Institute, Zhengzhou 450052, People's Republic of China*

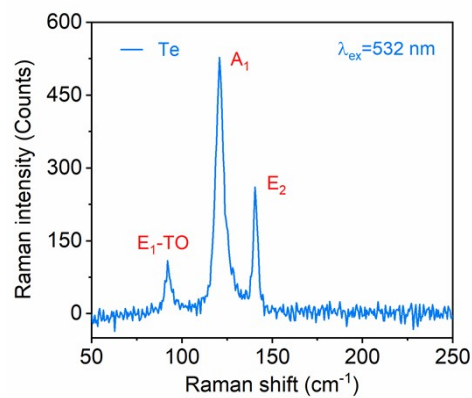
<sup>c</sup> *National Joint Engineering Research Center for Abrasion Control and Molding of Metal Materials, School of Materials Science and Engineering, Henan University of Science and Technology, Luoyang 471003, People's Republic of China*

<sup>d</sup> *School of Physical Science and Technology, Southwest Jiaotong University, Chengdu 610031, People's Republic of China*

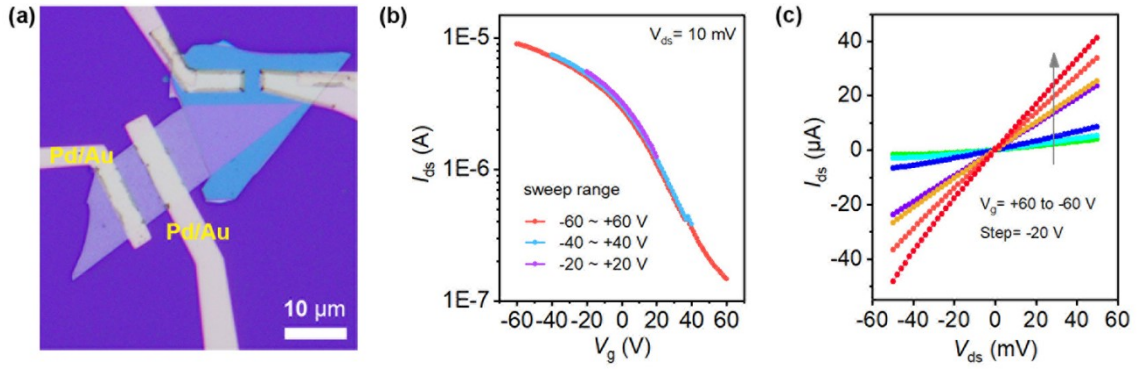
E-mail: linpei@zzu.edu.cn; lixj@zzu.edu.cn



**Fig. S1** Atomic force microscope images of the as-synthesized Te nanoflakes, a typical thickness in the range of  $\sim 20$ -60 nm can be observed.

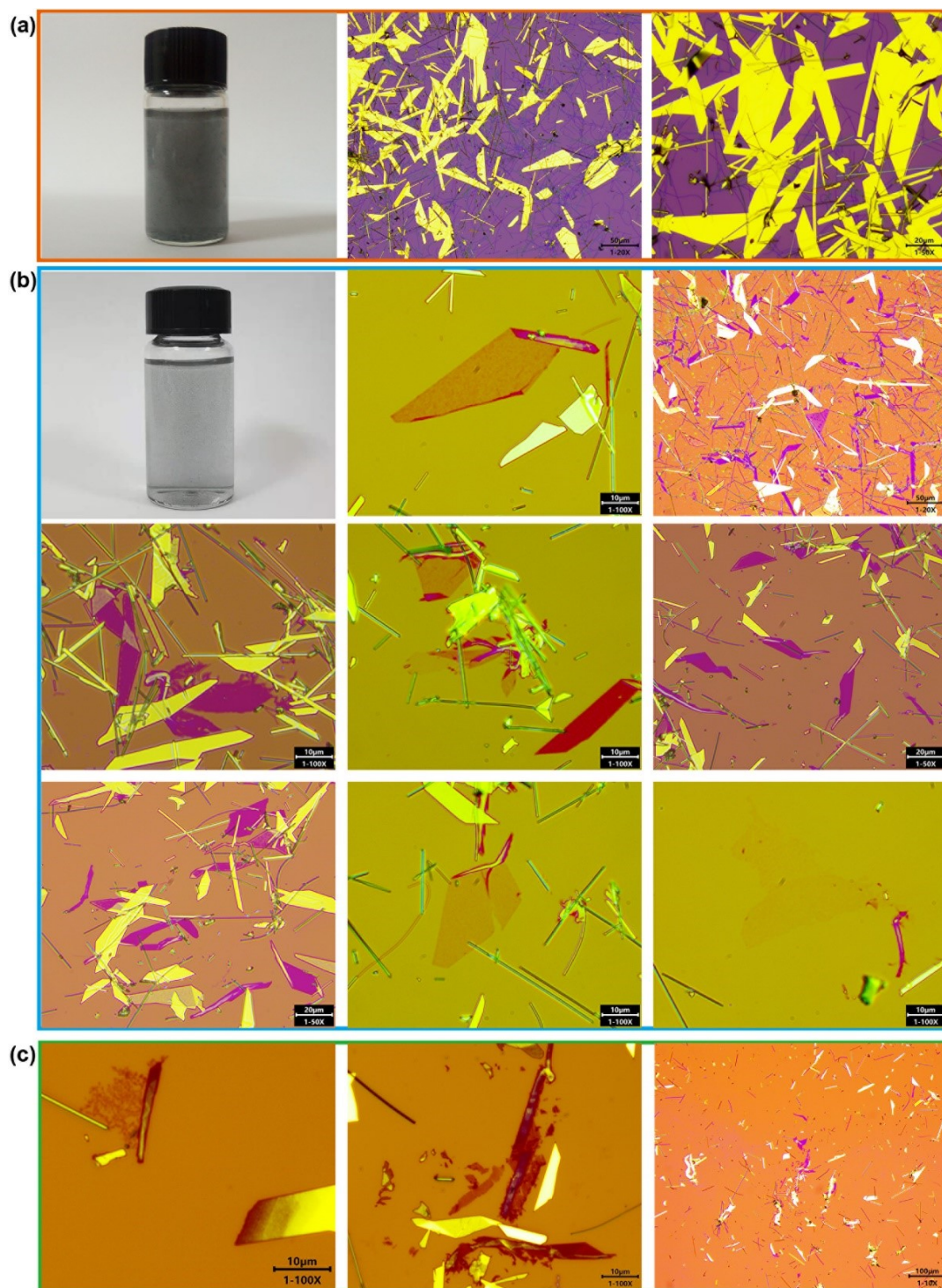


**Fig. S2** Raman spectrum of a typical Te flake, the excitation wavelength is 532 nm. Three characteristic Raman-active modes  $E_1$ -TO,  $A_1$ , and  $E_2$  can be clearly observed.



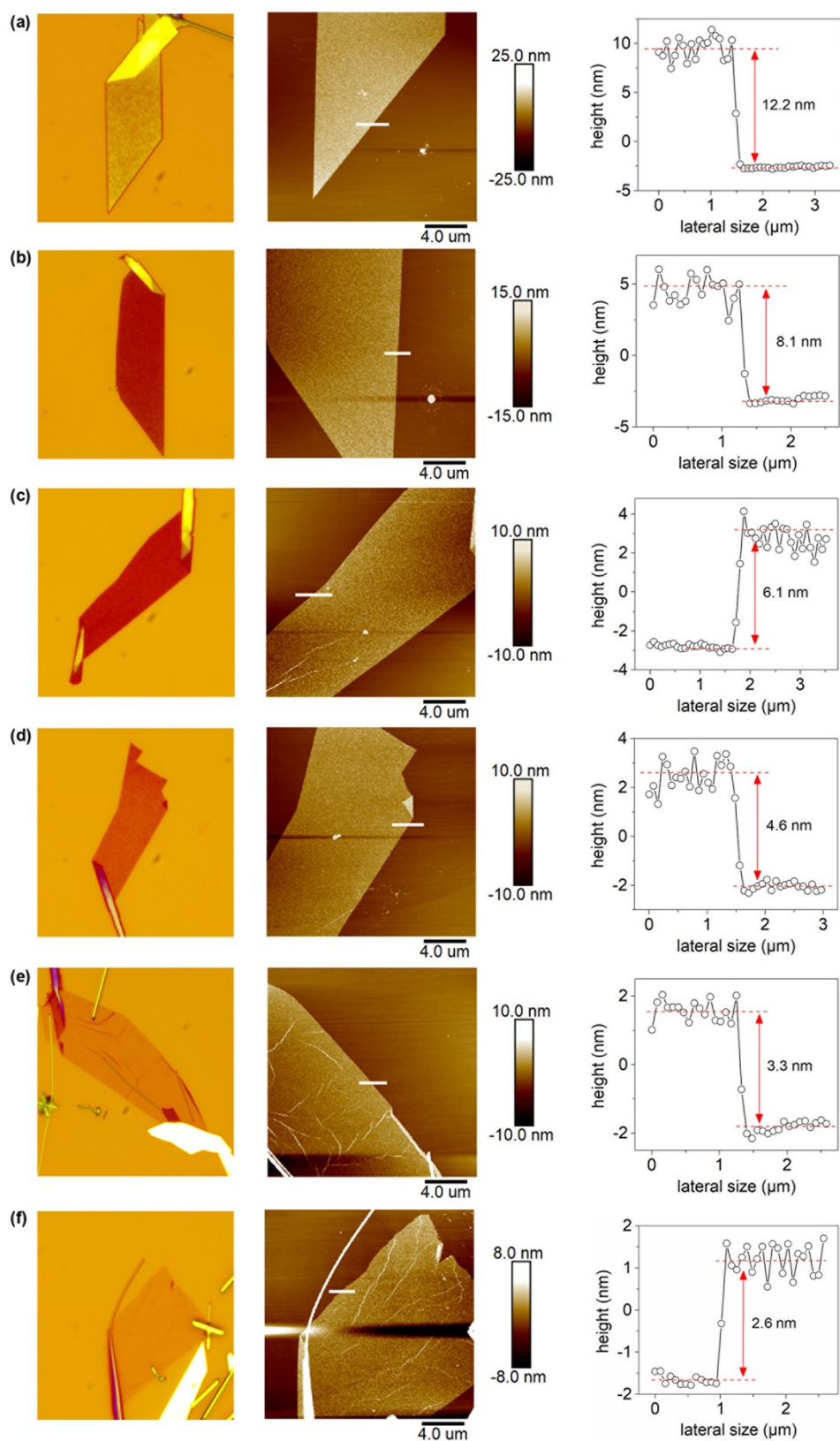
**Fig. S3** Electrical characterization of the Te field effect transistor. (a) Optical image of the measured Te transistor, the channel length and width are 4.3 μm and 12 μm, respectively. Thickness of the SiO<sub>2</sub> dielectric layer is 300 nm. The evaporated Pd/Au was used as contact electrode. (b) Transfer curve ( $I_{ds}$ - $V_g$ ) of the transistor at  $V_{ds}=10$  mV, showing a typical p-type characteristic. (c) Output curves ( $I_{ds}$ - $V_{ds}$ ) of the transistor under different gate voltages. Field-effect mobility of  $\sim 340$  cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup> was calculated from the transfer curves by using equation  $\mu=[dI_{ds}/dV_g] \times [L/(WC_{oxide}V_{ds})]$ , where  $L/W$  is the ratio between channel length and width, and  $C_{oxide}$  is the capacitance between the channel and the back gate per unit area.

Despite high carrier mobility, the current on/off ratio is typically low ( $< 10^3$ ) due to the large thickness of as-grown Te flakes. This will lead to a large dark current, which is unfavorable for IR photodetection. Therefore, to suppress the dark current and noise, construction of p-n junction with high gate tunability is still needed for high-performance Te photodetector.

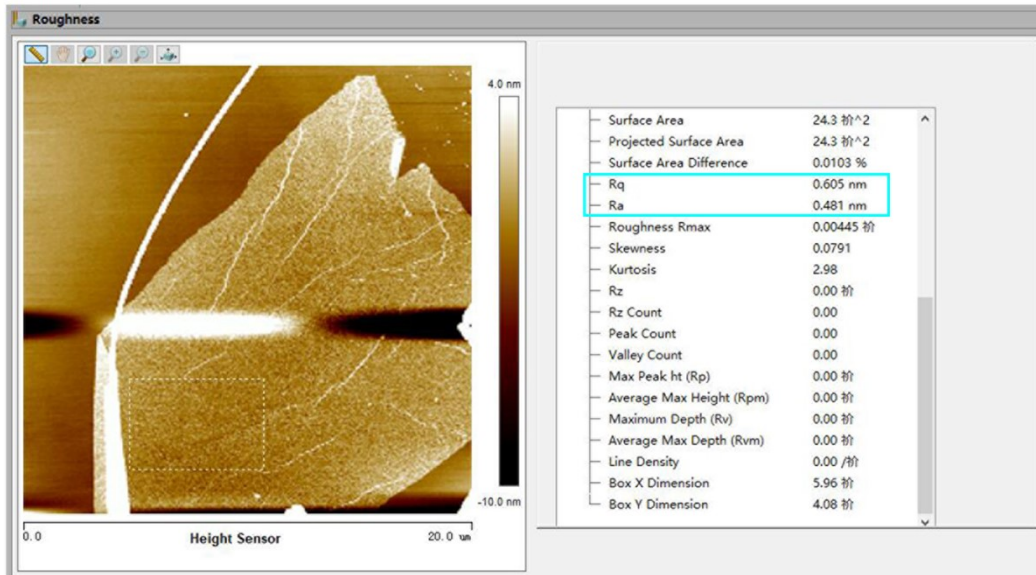


**Fig. S4** (a) Optical images of the as-grown Te flakes, which typically show yellow color because of the large material thickness. (b) After storage in DI water (4 °C) for about two weeks, large-scale thin flakes start to emerge and show different colors depending on the thickness. (c) With time further prolonged (about several days), the Te are completely oxidized and dissolved in DI water. Almost no intact ultrathin flakes can be observed.

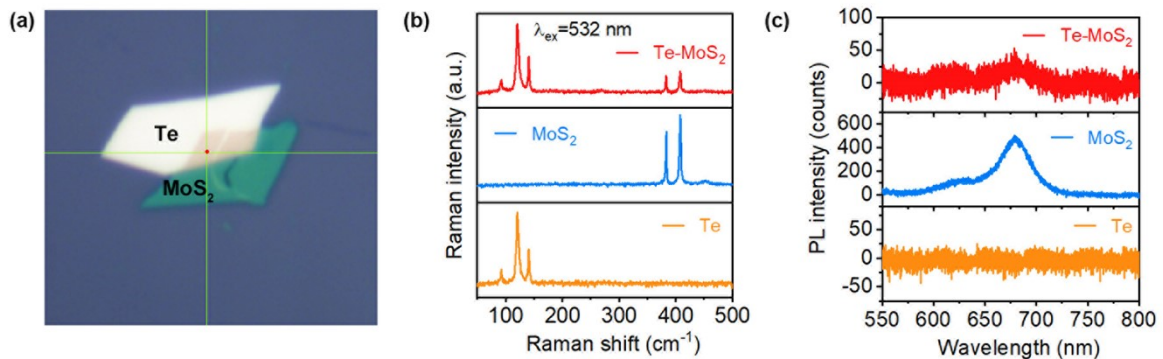




**Fig. S5** Optical and atomic force microscope images of thin tellurium. It shows that the Te nanoflakes present different colors depending their thickness, which can be used as a quick guide to estimate the thickness of 2D Te.

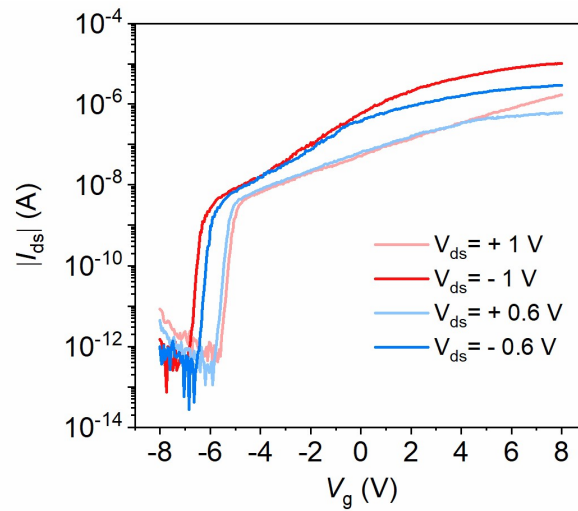


**Fig. S6** Evaluation of the surface roughness from AFM image of the Te flake ( $\sim 2.6$  nm thick). The root-of-mean-square (RMS) roughness ( $R_q$ ) was measured to be  $\sim 0.6$  nm, and the arithmetic mean roughness ( $R_a$ ) is only  $\sim 0.48$  nm.

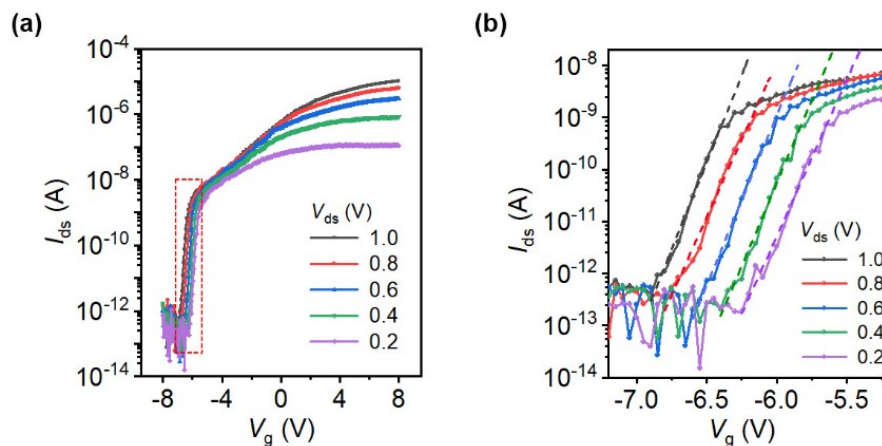


**Fig. S7** Raman (b) and photoluminescence (c) characterization of the Te-MoS<sub>2</sub> heterojunction with few-layer MoS<sub>2</sub> on top of Te, the excitation wavelength is 532 nm. Both the characteristic Raman peaks of Te and MoS<sub>2</sub> can be observed from the overlapped area. Compared with individual MoS<sub>2</sub>, the PL intensity of MoS<sub>2</sub> measured from the junction area is significantly reduced. As shown in Fig. 4d in the manuscript, a type I heterostructure is formed between Te and MoS<sub>2</sub>, and the valence band of MoS<sub>2</sub> is lower than Te. As a result, the photogenerated holes in MoS<sub>2</sub> would transfer to Te instead of

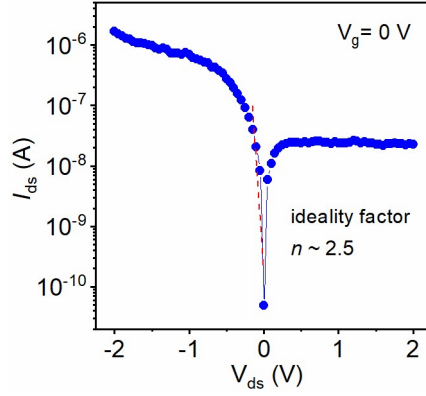
radiatively combining with electrons, resulting in the quenching of MoS<sub>2</sub> PL. This result confirms the effective charge transfer between Te and MoS<sub>2</sub>, indicating the formation of high-quality vdW heterojunction.



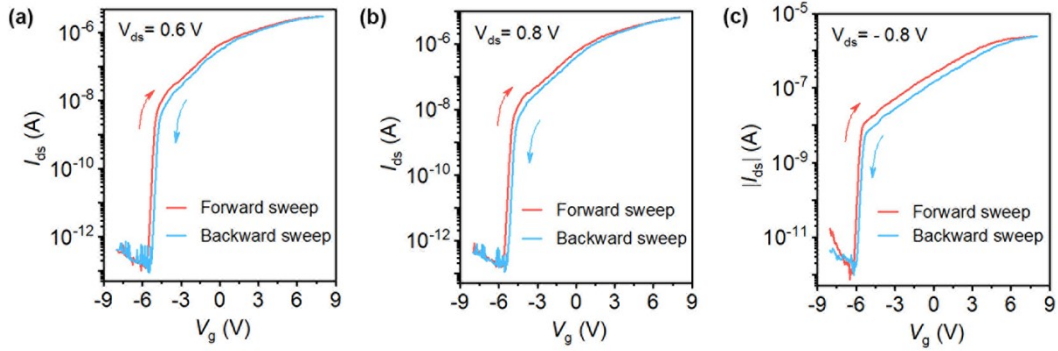
**Fig. S8** Transfer curves ( $I_{ds}$ - $V_g$ ) of the Te/MoS<sub>2</sub> heterojunction FET at positive and negative drain bias, exhibiting typical n-type transport characteristic. As presented in Fig. S3, the Te shows poor gate tunability because of its large thickness. Therefore, the modulation of heterostructure with  $V_g$  should be dominated by the bottom, gate-facing MoS<sub>2</sub>.



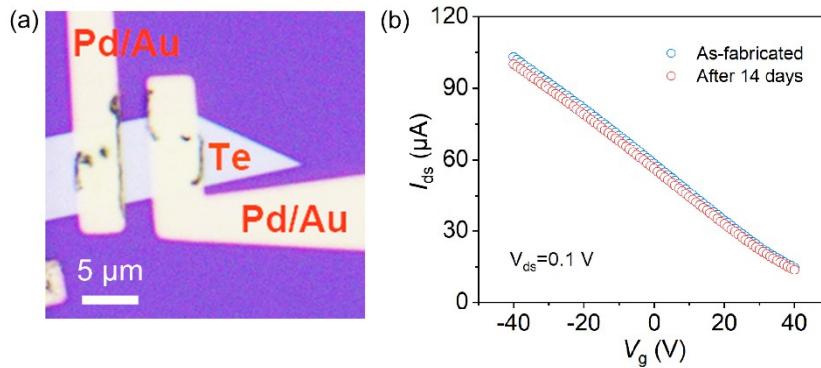
**Fig. S9** (a) Transfer curves ( $I_{ds}$ - $V_g$ ) of the Te/MoS<sub>2</sub> heterojunction at drain voltages of 0.2-1.0 V. (b) The enlarged view of red dashed box shown in (a), which shows a subthreshold swing (SS) of  $\sim 150$  mVdec<sup>-1</sup>.



**Fig. S10**  $I_{ds}$ - $V_{ds}$  curve of the Te/MoS<sub>2</sub> heterojunction at  $V_g=0$  V, from which an ideality factor of  $\sim 2.5$  could be obtained.

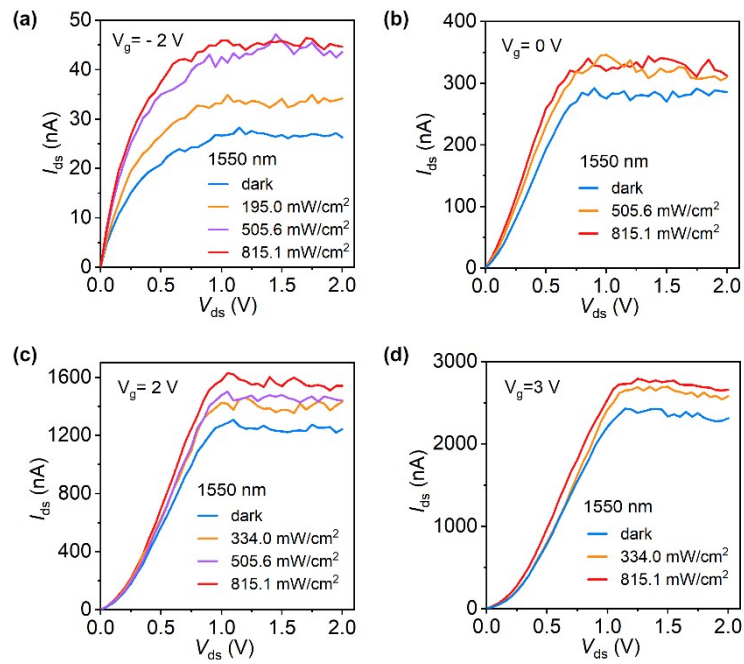


**Fig. S11** Hysteretic behavior of the device at fixed drain voltage of 0.6, 0.8, and -0.8 V, demonstrating negligible hysteresis. This can be attributed to the dangling-bond-free surface of h-BN, which reduces the density of trap states.

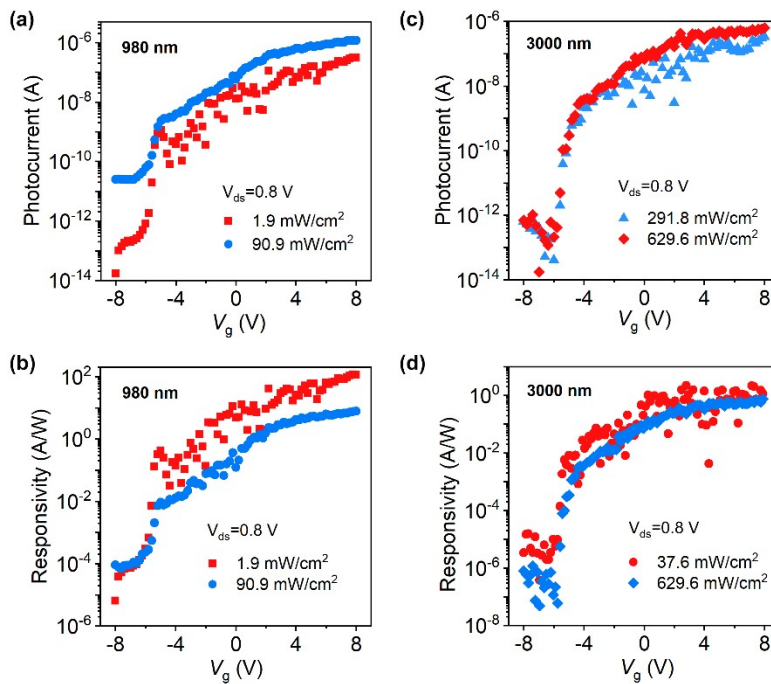


**Fig. S12** (a) Optical image of the Te field-effect transistor constructed on SiO<sub>2</sub>/Si substrate. (b) Transfer curves of the device, showing negligible change after being exposed in air for 14 days.



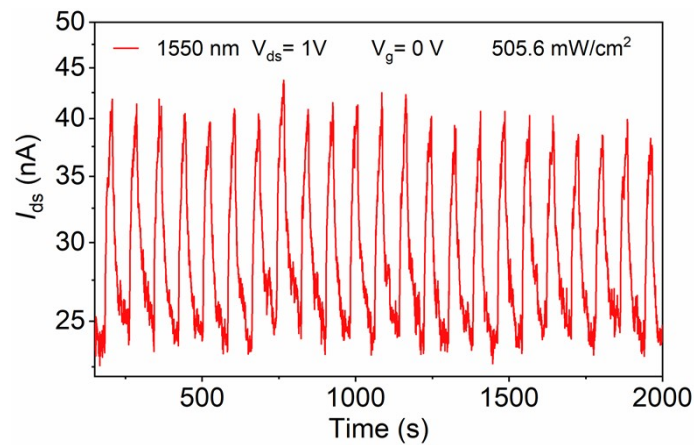


**Fig. S13** Photoresponse of output ( $I_{ds}$ - $V_{ds}$ ) curves of the Te/MoS<sub>2</sub> heterojunction device under different 1550 nm illumination power densities and gate voltages.

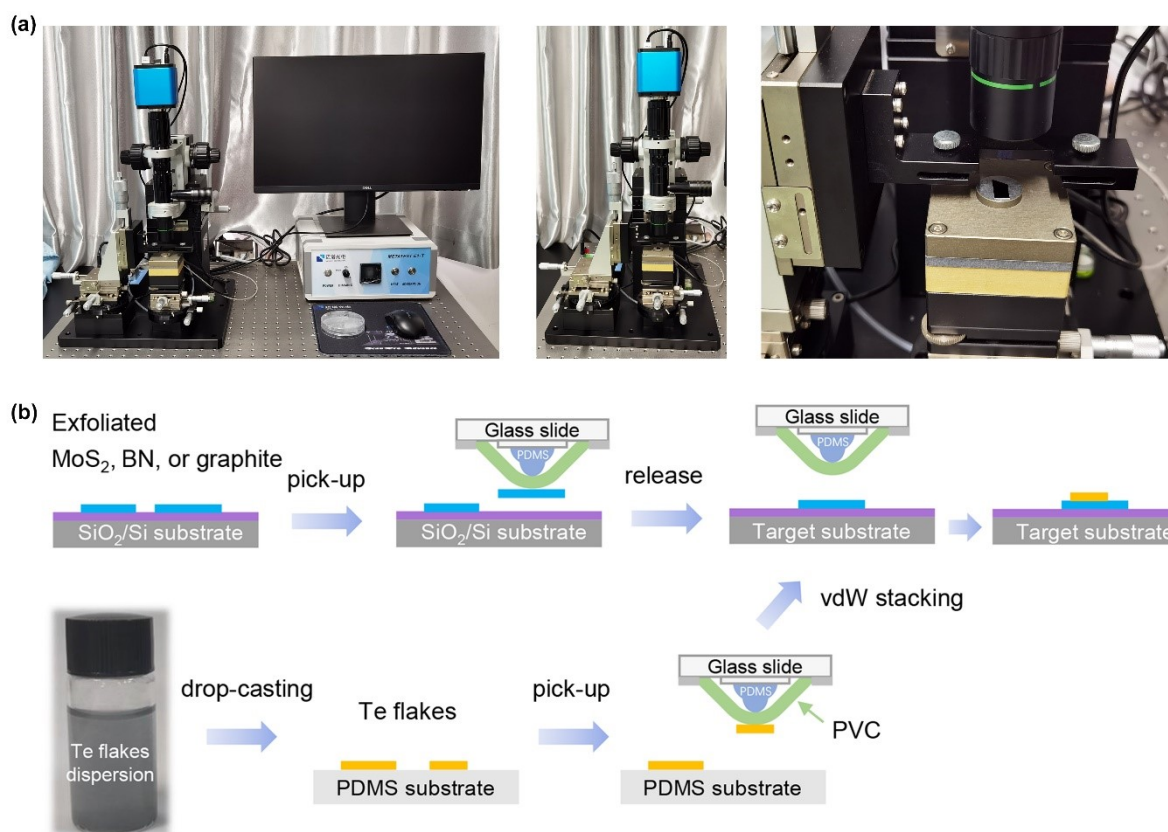


**Fig. S14** Calculated net photocurrent (a) and responsivity (b) of the Te/MoS<sub>2</sub> heterojunction at different gate voltages and 980 nm incident laser power densities. The photocurrent (c) and responsivity (d) of the device under different gate voltages and 3000 nm illumination densities.

As mentioned in the manuscript, the net photocurrent was calculated by  $I_{\text{ph}}=I_{\text{ds}}-I_{\text{dark}}$ , where the  $I_{\text{ds}}$  and  $I_{\text{dark}}$  indicate the drain current with and without illumination. The photoresponsivity  $R$  is defined as  $R=I_{\text{ph}}/(P_{\text{in}}\times A)$ ,  $P_{\text{in}}$  is the power density,  $A$  is the effective device area  $\sim 138.6 \mu\text{m}^2$  (measured from the device's optical image). In the case of an illumination power density ( $P_{\text{in}}$ ) of  $12.4 \text{ mW/cm}^2$  (980 nm), the  $I_{\text{ds}}$  and  $I_{\text{dark}}$  is  $5.31\times 10^{-6} \text{ A}$  and  $4.82\times 10^{-6} \text{ A}$ , respectively, yielding a photocurrent  $I_{\text{ph}}$  of  $4.9\times 10^{-7} \text{ A}$ . As a result, the photoresponsivity  $R$  is calculated to be  $\sim 28.4 \text{ A/W}$ . The detectivity is defined as  $D=R\cdot A^{1/2}/(2e\cdot I_{\text{dark}})^{1/2}$ , where  $R$  is the responsivity ( $28.4 \text{ A/W}$ ),  $A$  is the effective area ( $\sim 138.6 \mu\text{m}^2$ ),  $e$  is the electronic charge ( $1.602\times 10^{-19} \text{ C}$ ), and  $I_{\text{dark}}$  is the dark current ( $4.82\times 10^{-6} \text{ A}$ ). Accordingly, a detectivity of  $2.7\times 10^{10}$  Jones could be achieved.



**Fig. S15** Time-dependent photoresponse of the Te/MoS<sub>2</sub> vdW heterojunction under the switched on/off 1550 nm illumination. The applied drain voltage is 1.0 V and the gate voltage is 0 V.



**Fig. S16** (a) Optical images of the accurate transfer platform, with which the alignment of materials could be precisely controlled. (b) Schematic representation of the Te/MoS<sub>2</sub>/BN/graphite heterostructure stacking procedure following the previous report.<sup>1</sup> The MoS<sub>2</sub> (BN, or graphite) nanosheets were mechanically exfoliated on SiO<sub>2</sub>/Si substrate. Then, using a micro-domed PDMS covering poly(vinyl chloride) film (PVC, Reynolds), the material with proper thickness was selectively picked up. Here, the micro-domed PDMS structure was created on the glass slide by sequential dropping and curing of the increasingly small PDMS droplets. Because the adhesion force between PVC and MoS<sub>2</sub> (BN, or graphite) reaches a maximum at  $\sim 70$  °C, the picking-up procedure was carried out at that temperature. The adhesion force becomes negligible at  $\sim 130$  °C, allowing materials to be easily released onto the target substrate. However, due to the stronger adhesion force between Te and SiO<sub>2</sub>, PVC cannot pick up the Te flakes drop-casted on SiO<sub>2</sub>/Si substrate. As a result, the PDMS substrate was used for the Te flakes drop-casting.

**Supporting Reference:**

1. Y. Wakafuji, R. Moriya, S. Masubuchi, K. Watanabe, T. Taniguchi and T. Machida, *Nano Lett.*, 2020, **20**, 2486-2492.