

Electronic Supplementary Information
Exploring the feasibility and conduction mechanisms of
P-type nitrogen-doped β -Ga₂O₃ with high hole mobility

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1. Formation energies of phase transition from GaN to β -Ga₂O₃

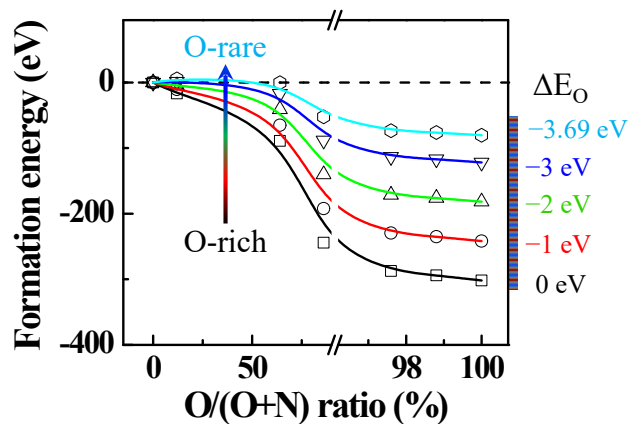


Fig. S1. Formation energy versus O/(O+N) ratio from 0 (GaN) to 1 (β -Ga₂O₃) under different O ambient conditions.

2. Band structure of Ga₄₈O₇₀N₂

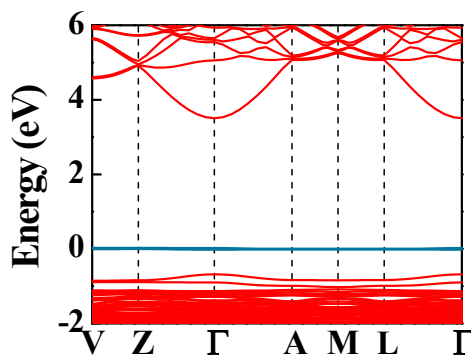


Fig. S2. Band structure of Ga₄₈O₇₀N₂.

The band structure of Ga₄₈O₇₀N₂ shows an impurity level (blue line) at 0.56 eV above VBM.

3. Ultraviolet-photoelectron-spectroscopy (UPS) valence band spectrum of N-doped β - Ga_2O_3

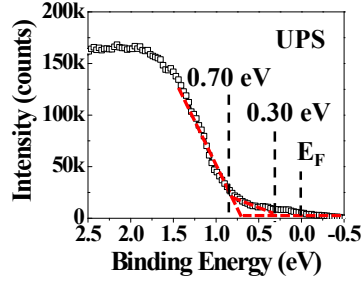


Fig. S3. UPS (21.22 eV He I radiation) valence band photoemission spectrum of the valence band edge.

4. Formation energies of neutral and charged $\text{Ga}_{48}\text{O}_{70}\text{N}_2$

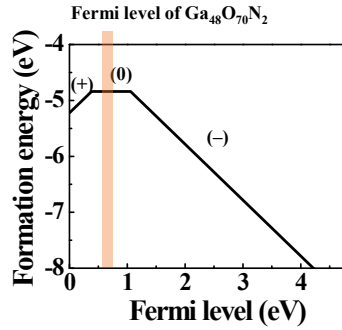


Fig. S4. Formation energies of neutral and charged $\text{Ga}_{48}\text{O}_{70}\text{N}_2$ versus Fermi level.

The relationship between formation energy of neutral $\text{Ga}_{48}\text{O}_{70}\text{N}_2$ and Fermi level is described by $E_f[N_0^0(\text{Ga}_{48}\text{O}_{70}\text{N}_2), E_F = 0] = E_{\text{tot}}[N_0^0(\text{Ga}_{48}\text{O}_{70}\text{N}_2)] - E_{\text{tot}}(\text{Ga}_{48}\text{O}_{69}\text{N}_3) + E_N - E_O$.

The relationship between formation energy of $\text{Ga}_{48}\text{O}_{70}\text{N}_2^-$ and Fermi level is described by $E_f[N_0^-(\text{Ga}_{48}\text{O}_{70}\text{N}_2), E_F = 0] = E_{\text{tot}}[N_0^-(\text{Ga}_{48}\text{O}_{70}\text{N}_2)] - E_{\text{tot}}(\text{Ga}_{48}\text{O}_{69}\text{N}_3) + E_N - E_O - (E_{VBM} + E_F + E_{\text{corr}})$.

The relationship between formation energy of $\text{Ga}_{48}\text{O}_{70}\text{N}_2^+$ and Fermi level is described by $E_f[N_0^+(\text{Ga}_{48}\text{O}_{70}\text{N}_2), E_F = 0] = E_{\text{tot}}[N_0^+(\text{Ga}_{48}\text{O}_{70}\text{N}_2)] - E_{\text{tot}}(\text{Ga}_{48}\text{O}_{69}\text{N}_3) + E_N - E_O + (E_{VBM} + E_F$

+ E_{corr}).

5. Numerical method and computational details of scattering characteristics

The hole mobility and scattering characteristics are predicted based on the Boltzmann transport theory using the *Ab initio* Scattering and Transport package. The scattering rate from initial state $|nk\rangle$ to final state $|mk+q\rangle$ is defined using Fermi's gold rule as $\tau_{nk \rightarrow mk+q}^{-1} = \frac{2\pi}{\hbar} |g_{mn}(k,q)|^2 \delta(\varepsilon_{nk} - \varepsilon_{mk+q})$, where ε_{nk} denotes the energy of state $|nk\rangle$, ε_{mk+q} the energy of state $|mk+q\rangle$, $g_{mn}(k,q)$ the coupling matrix element for scattering from state $|nk\rangle$ to state $|mk+q\rangle$, k the wave vector, n and m the band index. Acoustic deformation potential scattering (ADP), optical phonon scattering (OPP) and ionized impurity scattering (IOI) are addressed. The impact of various scatterings on hole mobility is expressed via the coupling matrix element.

The matrix element of ADP is given by:

$$g_{nm}^{ADP} = \sqrt{K_B T} \sum_{G \neq -q} \left[\frac{D_{nk} \cdot \tilde{S}_l}{C_l \sqrt{\rho_m}} + \frac{D_{nk} \cdot \tilde{S}_{t_1}}{C_{t_1} \sqrt{\rho_m}} + \frac{D_{nk} \cdot \tilde{S}_{t_2}}{C_{t_2} \sqrt{\rho_m}} \right] \langle mk+q | e^{i(q+G) \cdot r} | nk \rangle$$

, where D_{nk}

denotes the deformation potential tensor, \tilde{S} the unit strain associated with an acoustic mode, ρ_m the density of materials, c the wave velocity, and subscripts l , t_1 and t_2 indicate properties

belonging to longitudinal and transverse modes. The deformation potential tensor D_{nk} is obtained by $D_{nk} + V_{nk} \otimes V_{nk}$, where D_{nk} denotes the second rank deformation potential tensor,

V_{nk} the group velocity. The unit strain associated with an acoustic mode (\tilde{S}) is obtained by $\hat{q} \otimes \hat{u}$, where \hat{u} denotes the unit vector of phonon polarization. The matrix element of OPP is

described by:

$$g_{nm}^{OPP} = \left[\frac{\hbar \omega_{op}}{2} \right]^{\frac{1}{2}} \sum_{G \neq -q} \left(\frac{1}{\hat{n} \cdot \varepsilon_{\infty} \cdot \hat{n}} - \frac{1}{\hat{n} \cdot \varepsilon_s \cdot \hat{n}} \right) \frac{\langle mk+q | e^{i(q+G) \cdot r} | nk \rangle}{|q+G|}$$

, where

ω_{op} denotes the optical phonon frequency, $\hat{n} = (q+G)/|q+G|$ a unit vector in the direction of

scattering, ϵ_s the static dielectric tensor, ϵ_∞ the high-frequency dielectric tensor. The matrix

element of IOI is given by:
$$g_{nm}^{IOI} = \sum_{G \neq -q} \frac{n_{ii}^{1/2} Z e}{2 \hat{n} \cdot \epsilon_s \cdot \hat{n}} \frac{\langle mk + q | e^{i(q+G) \cdot r} | nk \rangle}{|q + G|^2 + \beta^2},$$
 where Z denotes the charge state of the impurity center, $n_{ii} = C_{self} \times (p - n)/Z$ the concentration of ionized impurities, C_{self} the amount of charge compensation, and β the inverse screening length.

The β^2 is obtained by $\frac{e^2}{\epsilon_s K T} \int \frac{D(\epsilon) f(1-f)}{V} d\epsilon$, where V denotes the unit cell volume, $D(\epsilon)$ the density of states, and f Fermi-Dirac distribution.

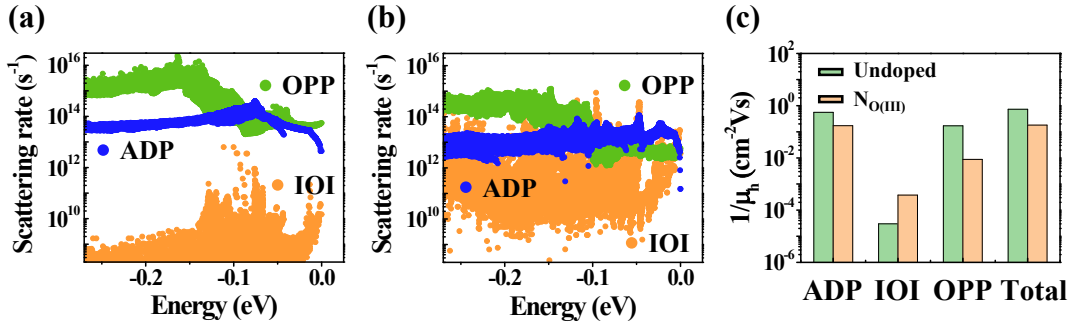


Fig. S5. Scattering rates of ADP, IOI, and OPP for (a) undoped β -Ga₂O₃ and (b) β -Ga₂O₃:N_{O(III)} obtained by the spin-polarized calculations. (c) Calculated $1/\mu_{ADP}$, $1/\mu_{IOI}$, $1/\mu_{OPP}$, and $1/\mu_{Total}$ of undoped β -Ga₂O₃ and β -Ga₂O₃:N_{O(III)}, where μ_{Total} is the total hole mobility, μ_{ADP} , μ_{IOI} , and μ_{OPP} the hole mobilities related to the ADP, IOI and OPP, respectively.

Material	p (cm^{-3})	Temperature (K)	ADP (s^{-1})	IOI (s^{-1})	OPP (s^{-1})
Undoped $\beta\text{-Ga}_2\text{O}_3$	10^{10}	300	4.4×10^{12}	1.4×10^{10}	5.5×10^{13}
$\beta\text{-Ga}_2\text{O}_3$: $\text{N}_{\text{O(III)}}$					
Nonspin	10^{15}	300	4.3×10^{10}	4.8×10^{13}	3.3×10^{13}
Spin polarization	10^{15}	300	1.5×10^{11}	3.6×10^{12}	5.8×10^{12}

Table S1. Summary of room-temperature scattering rates of ADP, IOI, and OPP at VBM for undoped $\beta\text{-Ga}_2\text{O}_3$ and $\beta\text{-Ga}_2\text{O}_3\text{:N}_{\text{O(III)}}$ obtained by the nonspin and spin-polarized calculations.

Material	p (cm^{-3})	Temperature (K)	μ_{Total} ($\text{cm}^2\text{V}^{-1}\text{s}^{-1}$)	μ_{ADP} ($\text{cm}^2\text{V}^{-1}\text{s}^{-1}$)	μ_{IOI} ($\text{cm}^2\text{V}^{-1}\text{s}^{-1}$)	μ_{OPP} ($\text{cm}^2\text{V}^{-1}\text{s}^{-1}$)
Undoped $\beta\text{-Ga}_2\text{O}_3$	10^{10}	300	1.3	1.8	33300	5.8
$\beta\text{-Ga}_2\text{O}_3$: $\text{N}_{\text{O(III)}}$						
Nonspin	10^{15}	300	8.8	300.0	354.3	9.0
Spin polarization	10^{15}	300	5.5	5.8	2643.2	114.0

Table S2. Summary of room-temperature hole mobilities μ_{Total} , μ_{ADP} , μ_{IOI} , and μ_{OPP} for undoped $\beta\text{-Ga}_2\text{O}_3$ and $\beta\text{-Ga}_2\text{O}_3\text{:N}_{\text{O(III)}}$ obtained by the nonspin and spin-polarized calculations.

6. Growth and Materials characterization of N-doped $\beta\text{-Ga}_2\text{O}_3$ films

Growth of N-doped $\beta\text{-Ga}_2\text{O}_3$ films

The β -Ga₂O₃ films were prepared by thermal oxidation of GaN/sapphire substrates at 1000–1100 °C using a modified chemical vapor deposition system. The GaN surface was cleaned by acetone, ethanol, and deionized water for 10 mins in turn. After then the cleaned GaN films were placed in a quartz boat and inserted in a quartz tube furnace, followed by annealing of the GaN films under oxygen ambient. During the thermal oxidation process, the oxygen pressure, high-purity Ar gas purge, and pumping rate are well controlled. The growth rate and crystalline quality were maintained by well control of oxygen pressure and Ar gas purge for efficient N doping within the β -Ga₂O₃ films.

Materials characterization

Ultraviolet photoelectron spectroscopy (UPS, 21.22 eV He I radiation) was used to investigate the VB photoemission. The Fermi level (E_F) was calibrated by referring the Au energy. Temperature-dependent Van der Pauw Hall effect measurements were performed to analyze the p-type properties of the β -Ga₂O₃ films.