

## Supporting Information: Triplet-to-Singlet Exciton Transfer in Hyperfluorescent OLED materials

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# I. CALCULATION DETAILS

## A. Functional Tuning

| Molecule    | $\omega$ (bohr <sup>-1</sup> ) |
|-------------|--------------------------------|
| ACRSA       | 0.1439                         |
| TBPc        | 0.1388                         |
| ACRXTN      | 0.1563                         |
| TTPA        | 0.1262                         |
| PXZ-TRZ     | 0.1438                         |
| TBRb        | 0.1051                         |
| Tri-PXZ-TRZ | 0.1188                         |
| DBP         | 0.1114                         |

TABLE S1. Tuned values of the range separation parameter  $\omega$ .

All electronic structure calculations were performed with the  $\omega$ B97X-D functional and the 6-31G(d,p) basis set. The range separation parameter was tuned for each of the molecules, and the results are shown in Table S1.

## B. Absorption, Fluorescence and Phosphorescence Spectra

All spectra were calculated using the nuclear ensemble approach<sup>1</sup>. A total of  $N = 500$  geometries were sampled for each spectrum simulation from the following distribution

$$\rho(\vec{R}, T) = \prod_{i=1}^{3N-6} \left( \frac{\mu_i \omega_i}{2\pi\hbar \sinh(\hbar\omega_i/k_b T)} \right) \times \exp \left( -\frac{\mu_i \omega_i}{\hbar} R_i^2 \tanh \left( \frac{\hbar\omega_i}{2k_b T} \right) \right), \quad (\text{S1})$$

where  $\vec{R}$  is a displacement,  $T$  is the temperature (set to 300 K), and  $k_b$  is the Boltzmann constant. In addition,  $\mu_i$  and  $\omega_i$  are the reduced mass and frequency of the  $i$ th normal mode, respectively.

For absorption spectra, the normal modes are calculated for the  $S_0$  state, whereas for fluorescence and phosphorescence, they are taken from calculations in the  $S_1$  and  $T_1$  states, respectively.

For each sampled geometry, a single point TD-DFT calculation is performed, and the final spectra are obtained by averaging over results according to the following expressions for the absorption cross section ( $\sigma(E)$ ) and the differential emission rate ( $I(E)$ )<sup>1,2</sup>:

$$\sigma(E) = \frac{\pi e^2 \hbar}{2mc\epsilon_0} \frac{1}{N} \sum_{k=i}^N f_i G(E - \Delta E_i, \sigma), \quad (\text{S2})$$

$$I(E) = \frac{n^2}{3\pi\hbar^3 c^3 \epsilon_0} \frac{1}{N} \sum_{k=1}^N M_i^2 \Delta E_i^3 G(E - \Delta E_i, \sigma). \quad (\text{S3})$$

In these expressions,  $c$  is the speed of light,  $\epsilon_0$  is the vacuum permittivity, and  $e$  and  $m$  correspond to the electron's charge and mass, respectively. In addition,  $f$  is the oscillator strength of the transition,  $\Delta E_i$  is the vertical transition energy and  $M_i$  is the transition dipole moment. The function  $G(E - \Delta E_i, \sigma)$  is a normalized gaussian distribution with mean  $E - \Delta E_i$  and standard deviation  $\sigma = k_b T$ .

For phosphorescence spectra, it is necessary to employ perturbation theory to calculate the appropriate transition dipole moments, which are given by<sup>3</sup>

$$M_\gamma^\beta = \sum_m \frac{\langle T_m^\beta | H_{SO} | S_0 \rangle}{E(S_0) - E(T_m)} \langle T_1^\beta | M_\gamma | T_m^\beta \rangle + \sum_m \frac{\langle S_m | H_{SO} | T_1^\beta \rangle}{E(T_1) - E(S_m)} \langle S_m | M_\gamma | S_0 \rangle, \quad (\text{S4})$$

in which  $\gamma$  refers to the  $x, y$ , and  $z$  coordinates and  $\beta$  indexes the sublevels of the triplet states. The summations include ten singlet and ten triplet excited states and the total transition dipole moment is averaged over the triplet sublevels as

$$M_i^2 = \frac{1}{3} \sum_{\gamma, \beta} |M_\gamma^\beta|^2, \quad (\text{S5})$$

### C. Rate Calculations

To calculate Förster transfer rates, we employ a corrected version of the usual point-dipole approximation expression given by<sup>4,5</sup>

$$k_F = \frac{1}{\tau_{emi}} \left( \frac{R_F}{\alpha\mu + r} \right)^6 \quad (\text{S6})$$

where  $R_F$  is the Förster radius of the transfer,  $\tau_{emi}$  is the fluorescence lifetime (phosphorescence lifetime for triplet-to-singlet transfers),  $r$  is the intermolecular distance, and  $\alpha\mu$  is the correction term for small distances in which  $\alpha$  is a constant ( $1.15 e^{-1}$ ,  $e$  being the charge of the electron), and  $\mu$  is the molecule's transition dipole moment.

The Förster radius can be obtained from spectrum simulations by the following expression<sup>6</sup>

$$R_F^6 = \frac{9c^4 \kappa^2 \hbar^3 \tau_{emi}}{8\pi} \int_0^\infty \frac{dE}{E^4} I_D(E) \sigma_A(E) \quad (S7)$$

where  $\kappa$  is the orientation factor (set to 2/3 for isotropic distribution of dipoles),  $I_D$  is the donor molecule’s differential emission rate (for fluorescence or phosphorescence), and  $\sigma_A$  is the acceptor’s absorption cross section.

The radiative emission lifetime ( $\tau_{emi}$ ) for either fluorescence or phosphorescence can be obtained from the corresponding emission spectrum of the molecule and relates to the radiative emission rate  $k_{emi}$  by

$$k_{emi} = \frac{1}{\tau_{emi}} = \frac{1}{\hbar} \int_0^\infty I_D(E) dE \quad (S8)$$

Intersystem crossing (ISC) rates are calculated from the same ensemble of geometries used for spectrum simulations using the following expression<sup>2</sup>

$$k_{ISC} = \frac{2\pi}{\hbar} \frac{1}{N} \sum_{k=1}^N H_{SO_i}^2 G(\lambda + \Delta E_{fi}, \sqrt{2\lambda k_b T + \sigma^2}), \quad (S9)$$

in which  $H_{SO_i}$  is the spin-orbit coupling for the  $i$ -th sampled geometry and  $\lambda$  is the reorganization energy of the  $S_1 \rightarrow T_1$  transition in the case of ISC and the  $T_1 \rightarrow S_1$  transition in the case of reverse ISC (rISC). These reorganization energies are shown in Table S2

| Molecule    | Reorganization energy (eV) |                       |
|-------------|----------------------------|-----------------------|
|             | $S_1 \rightarrow T_1$      | $T_1 \rightarrow S_1$ |
| ACRSA       | 0.189                      | 0.118                 |
| ACRXTN      | 0.218                      | 0.221                 |
| PXZ-TRZ     | 0.192                      | 0.216                 |
| Tri-PXZ-TRZ | 0.192                      | 0.075                 |

TABLE S2. Reorganization energies used to calculate ISC and rISC rates.

## II. KINETIC MONTE CARLO DETAILS

Kinetic Monte Carlo simulations are run on a  $50 \times 50 \times 50$  cubic lattice. Each site in the lattice is considered to be either a TADF molecule or a fluorescent emitter. Host molecules are not taken

| Molecule    | Molar mass (g/mol) |
|-------------|--------------------|
| ACRSA       | 435.51             |
| TBPc        | 476.73             |
| ACRXTN      | 403.47             |
| TTPA        | 568.75             |
| PXZ-TRZ     | 490.55             |
| TBRb        | 757.09             |
| Tri-PXZ-TRZ | 826.90             |
| DBP         | 804.97             |

TABLE S3. Molar masses of all compounds analyzed in this work.

into account in the simulation since they do not participate in the processes of exciton transfer and emission. They are, however, accounted for in the estimation of intermolecular distances.

To determine intersite distances, we considered a box of volume  $V$  to contain  $n_1$  TADF dopants,  $n_2$  fluorescent emitters, and  $n_3$  host molecules. The average distance between two molecules of either TADF or emitter kind is given by

$$d = \left[ \frac{V}{(n_1 + n_2)} \right]^{1/3} \quad (\text{S10})$$

We consider further that the total volume can be given by  $V = Nd_0^3$ , in which  $N = n_1 + n_2 + n_3$  is the total number of molecules and  $d_0^3$  is an average volume occupied by each molecule. As such, we may write

$$d = d_0 \left[ \frac{N}{(n_1 + n_2)} \right]^{1/3} \quad (\text{S11})$$

The term in square brackets is the inverse of the fraction of TADF plus fluorescent emitters present in the box. This number can be connected to the weight fraction ( $p_i$ ) of each compound  $i$  employed in the experimental setup and their corresponding molar masses ( $m_i$ ), which are shown in Table II. This is given by

$$p_i = \frac{n_i m_i}{M} \Rightarrow n_i = \frac{p_i}{m_i} M \quad (\text{S12})$$

in which  $M$  is the total mass of the box. From this follows that

| System          | $\rho_1$ (%) | $\rho_2$ (%) | $d$ (Å) |
|-----------------|--------------|--------------|---------|
| ACRSA/TBPe      | 94.3         | 5.7          | 15.06   |
| ACRXTN/TTPA     | 98.6         | 1.4          | 10.99   |
| PXZ-TRZ/TBRb    | 97.5         | 2.5          | 13.86   |
| Tri-PXZ-TRZ/DBP | 93.6         | 6.4          | 18.90   |

TABLE S4. Fraction of lattice sites assigned as TADF and emitter molecules for weight fractions reported in Ref. 8 along with the corresponding estimated average intermolecular distances  $d_0$ .

$$\frac{(n_1 + n_2)}{N} = \frac{\frac{p_1}{m_1} + \frac{p_2}{m_2}}{\frac{p_1}{m_1} + \frac{p_2}{m_2} + \frac{p_3}{m_3}} \quad (\text{S13})$$

which allows calculation of  $d$  in Equation S11 once a  $d_0$  value is chosen. To estimate this parameter, we have considered results from molecular dynamics simulations of mCP host molecules<sup>7</sup>. These simulations included 100 mCP molecules with a resulting density of  $1 \text{ g cm}^{-3}$ . This result gives an average intermolecular distance of  $8.79 \text{ \AA}$ , which we use as our  $d_0$  value.

Finally, to determine the fraction of lattice sites that will be considered as TADF molecules ( $\rho_1$ ) or fluorescent emitters ( $\rho_2$ ) we calculate

$$\rho_i = \frac{n_i}{(n_1 + n_2)} = \frac{\frac{p_i}{m_i}}{\frac{p_1}{m_1} + \frac{p_2}{m_2}} \quad (\text{S14})$$

For the weight fractions used in the experimental paper, the proportions of TADF and emitter molecules in the lattices are shown in Table S4.

### III. SUPPORTING RESULTS

In this section, we present a series of supporting results from the various calculations performed. It includes a comparison between calculated and experimental absorption and emission energies, spectral overlap plots indicating the possibility of Förster transfers from emitters to TADF molecules and also between TADF molecules.

We also present tables with estimated fluorescence rates of the fluorescent emitters analyzed in this work, along with the Förster radii for singlet exciton transfers between emitters.

| Molecule    | Absorption (eV) |                   | Fluorescence (eV) |                   |
|-------------|-----------------|-------------------|-------------------|-------------------|
|             | Calc.           | Exp. <sup>8</sup> | Calc.             | Exp. <sup>8</sup> |
| ACRSA       | 3.04            | -                 | 2.91              | 2.55              |
| TBPc        | 3.06            | 2.82              | 2.63              | 2.69              |
| ACRXTN      | 2.91            | -                 | 2.81              | 2.53              |
| TTPA        | 2.75            | 2.61              | 2.27              | 2.34              |
| PXZ-TRZ     | 2.71            | -                 | 2.53              | 2.30              |
| TBRb        | 2.42            | 2.38              | 2.03              | 2.18              |
| Tri-PXZ-TRZ | 2.45            | -                 | 2.40              | 2.27              |
| DBP         | 2.32            | 2.10              | 1.94              | 2.03              |

TABLE S5. Comparison between calculated and experimental values for absorption and fluorescence energy peaks for all TADF/emitter pairs analyzed here.

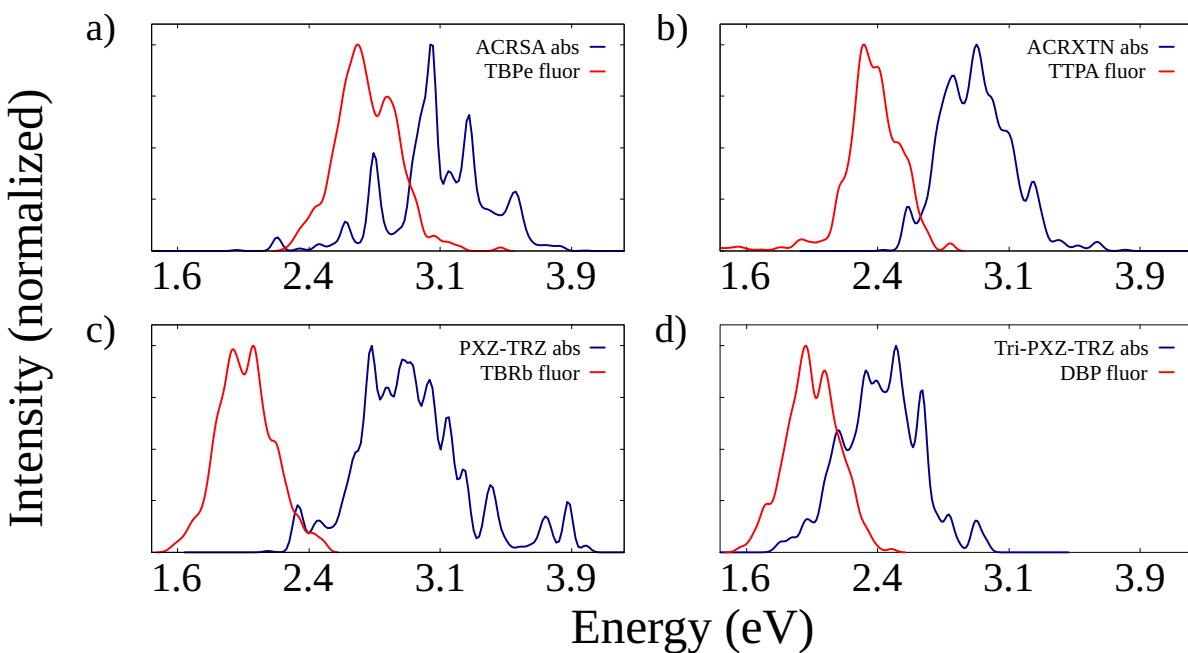


FIG. S1. Spectral overlaps between the fluorescence spectra of the fluorescent emitters and the absorption spectra of the TADF molecules they are paired with.

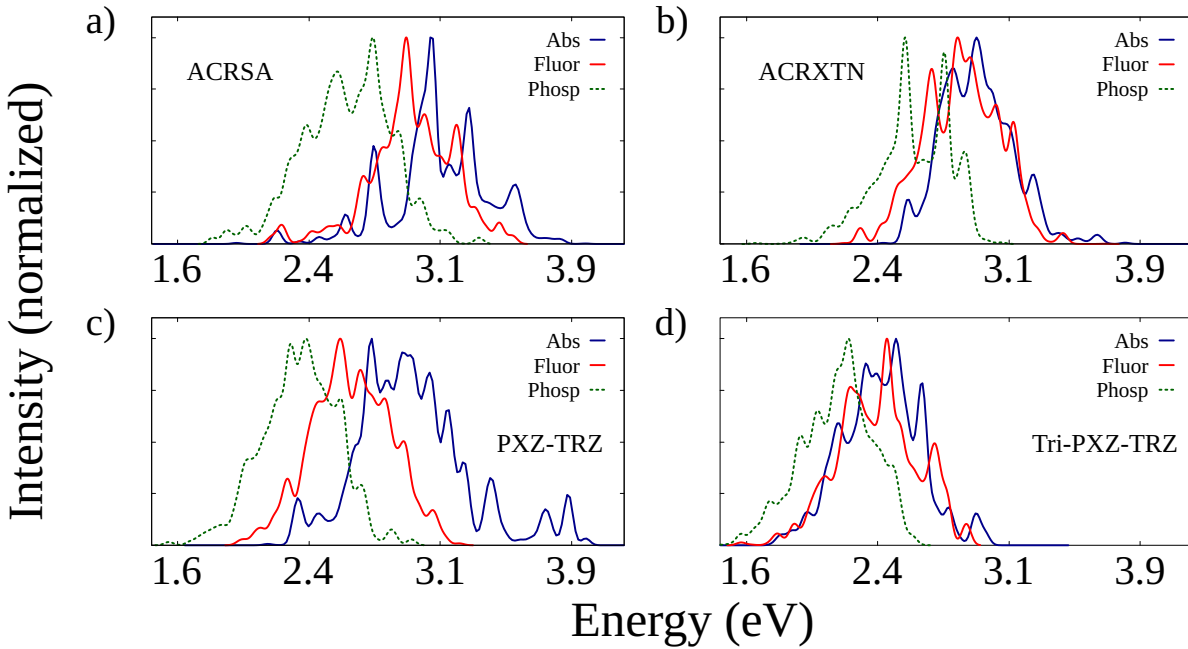


FIG. S2. Spectral overlaps between the fluorescence and phosphorescence spectra of the TADF molecules and their own absorption spectra.

| Molecule | Förster radius (Å)    |
|----------|-----------------------|
|          | $S_1 \rightarrow S_1$ |
| TBPe     | 50.1                  |
| TTPA     | 37.5                  |
| TBRb     | 51.8                  |
| DBP      | 73.9                  |

TABLE S6. Förster radii for singlet exciton transfers between different fluorescent emitters of the same kind.

| Emitter | Fluor. ( $s^{-1}$ ) |
|---------|---------------------|
| TBPe    | $8.13 \times 10^8$  |
| TTPA    | $6.66 \times 10^7$  |
| TBRb    | $2.46 \times 10^8$  |
| DBP     | $1.07 \times 10^9$  |

TABLE S7. Estimated fluorescence rates for the fluorescent emitters analyzed in this work.



| System          | $r_1 = \frac{\phi_{TTS}}{\phi_{rISC}\phi_{STS}}$ |                                   | $r_2 = \frac{\phi_{TTS}}{\phi_{rISC}}$ |                                   |
|-----------------|--|-----------------------------------|--|-----------------------------------|
|                 | $T_1^{TADF} \rightarrow S_1^{TADF}$              | $T_1^{TADF} \rightarrow S_1^{EM}$ | $T_1^{TADF} \rightarrow S_1^{TADF}$    | $T_1^{TADF} \rightarrow S_1^{EM}$ |
| ACRSA/TBPe      | 40.1   | 125.9                             | 0.7                                    | 99.0                              |
| ACRXTN/TTPA     | 41.6   | 68.8                              | 40.9                                   | 67.3                              |
| PXZ-TRZ/TBRb    | 1.3  | 56.3                              | 1.2                                    | 56.2                              |
| Tri-PXZ-TRZ/DBP | 0.0  | 1.2                               | 0.0                                    | 1.2                               |

TABLE S8. Ratios between the quantum yields of two processes that allow a triplet exciton to be transferred to a nearby molecule as a singlet ( $r_1$ ): TTS and rISC followed by singlet-to-singlet (STS) FRET. On the right, ratios between the quantum yields of two triplet conversion mechanisms ( $r_2$ ): TTS and rISC. The table shows the ratios for transfers between TADF molecules (superscript TADF) and from TADF to emitter molecules (superscript EM).

## REFERENCES

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