

**Supplemental Information for "Spontaneous magnetic merons in half-metallic  
 $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer with easy-plane anisotropy"**

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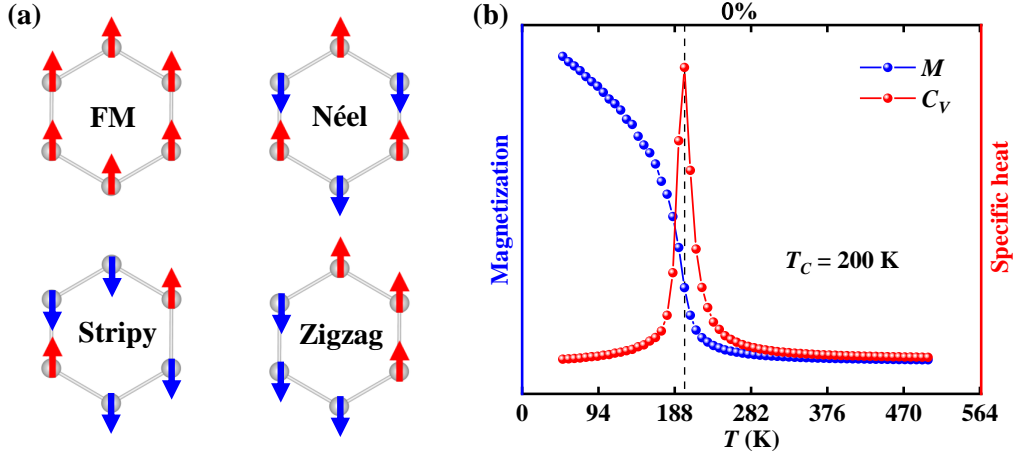


FIG. S1. (a) Schematic diagram of the four spin configurations used to extract the Heisenberg exchange coefficients  $J_1$ ,  $J_2$  and  $J_3$ . Red and blue arrows represent spin up (Z-direction) and down ( $-Z$ -direction), respectively. (b) Magnetic moment and specific heat capacity versus temperature of  $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer.

To obtain the Heisenberg exchange coefficients  $J_1$ ,  $J_2$  and  $J_3$ , we choose four spin configurations named *FM*, *Néel*, *Stripy* and *Zigzag* as shown in Fig. S1(a), with the energies being:

$$E_{FM} = -6J_1 - 12J_2 - 6J_3 + 4K(S^z)^2 + E_0, \quad (\text{S1})$$

$$E_{Néel} = 6J_1 - 12J_2 + 6J_3 + 4K(S^z)^2 + E_0, \quad (\text{S2})$$

$$E_{Stripy} = 2J_1 + 4J_2 - 6J_3 + 4K(S^z)^2 + E_0, \quad (\text{S3})$$

$$E_{Zigzag} = -2J_1 + 4J_2 + 6J_3 + 4K(S^z)^2 + E_0, \quad (\text{S4})$$

where  $J_1$ ,  $J_2$ ,  $J_3$  are the Heisenberg exchange coefficients of first, second and third nearest neighbor, respectively.  $K$  is the single ion anisotropy.  $S^z$  is the off-plane component of the normalized spin and  $E_0$  is the spin-independent energy. When the magnetic parameters are obtained, we perform Monte Carlo simulations with the Metropolis algorithm to evaluate the  $T_C$  of  $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer. As shown in Fig. S1(b), the  $T_C$  of intrinsic  $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer is about 200 K. The  $T_C$  with different strains can be found in Fig. S2

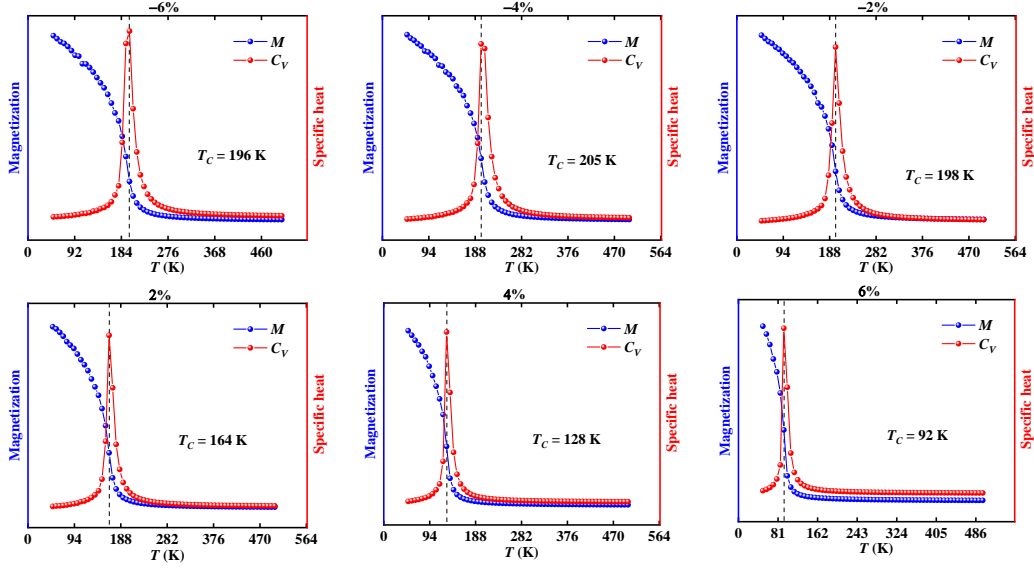


FIG. S2. Magnetic moment and specific heat capacity versus temperature of  $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer with different biaxial strains.

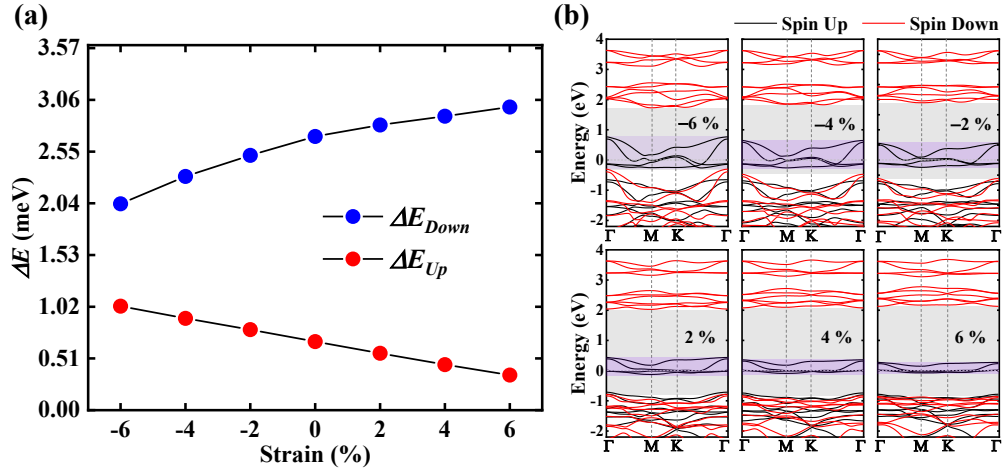


FIG. S3. (a) The band gap of spin down electrons ( $\Delta E_{Down}$ ) and band width of spin up electrons around Fermi level ( $\Delta E_{Up}$ ), and (b) band structures of  $\text{Mn}_2\text{I}_3\text{Br}_3$  with different strains.

As shown in Fig. S3(a), we illustrate the band gap of spin down electrons ( $\Delta E_{Down}$ ) and band width of spin up electrons around Fermi level ( $\Delta E_{Up}$ ) as functions of strains in  $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer. The band structures of  $\text{Mn}_2\text{I}_3\text{Br}_3$  with different strains are illustrated in Fig. S3(b) and  $\Delta E_{Up}$  ( $\Delta E_{Down}$ ) is marked by the purple (gray) shaded areas. It's clearly

seen that the half-metallic feature can be further enhanced by tensile strain due to the more localized spin-up channel around Fermi level and larger band gap of the spin-down channel.

TABLE S1. The optimized lattice constant  $a$ , Heisenberg exchange coefficients ( $J_1, J_2, J_3$ ), in-plane ( $d_{//}$ ) and off-plane ( $d_z$ ) components of DMI, single ion anisotropy ( $K$ ), the magnetic moments of Mn atoms and the  $T_C$  of  $\text{Mn}_2\text{I}_3\text{Br}_3$  monolayer with different strains.

Strain (%)	-6	-4	-2	0	2	4	6
$a$ (Å)	6.24	6.38	6.51	6.64	6.77	6.91	7.04
$J_1$ (meV)	-13.37	-16.16	-15.42	-14.28	-8.40	-4.57	-1.68
$J_2$ (meV)	-1.05	-1.36	-2.10	-2.60	-3.46	-3.38	-3.04
$J_3$ (meV)	-7.43	-5.74	-4.33	-3.85	-2.23	-1.51	-0.98
$d_{//}$ (meV)	1.68	1.91	2.38	3.00	3.71	3.99	3.83
$d_z$ (meV)	1.07	1.20	1.46	1.81	2.19	2.31	2.18
$K$ (meV)	-13.15	-12.47	-10.88	-9.95	-10.02	-9.33	-7.75
$M_{Mn}$ ( $\mu_B$ )	4.00	4.01	4.01	4.02	4.04	4.07	4.12
$T_C$ (K)	196	205	198	200	164	128	93