# Supplementary Materials for Thermal transport across copper-water interface from deep potential molecular dynamics

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## 13 SI. Validation tests

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To test the accuracy of DLP in a large system, we enlarge the original system (139 atoms) along the x- and z-directions four times into a large system (2224 atoms). Both small and large systems were simulated at 300K in the NVT ensemble, for 0.65 ns with a time step of 0.5 fs. We analyzed the average values of the density and radial distribution function (RDF) within the last 0.55 ns. Fig. S1 shows an excellent agreement for RDF between the large system and the small system. Therefore, we conclude that the trained DLP could generalize the accuracy of DFT to a larger system.

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Fig. S1 Comparison of the RDFs of water obtained from DPMD simulations between small
 system and large system. The red line represents the small system, and the black dashed line
 represents the large system.

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### 27 SII. LJ potential model

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The generally accepted Lennard-Jones (LJ) potential model coupling with the coulomb interaction is employed to describe the atomic interactions in the CMD simulations. The specific expression is as follows:

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$$u_{ij} = \frac{q_i q_j}{4\pi\varepsilon_0 r_{ij}} + 4\varepsilon_{ij} \left[ \left( \frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left( \frac{\sigma_{ij}}{r_{ij}} \right)^6 \right]$$
 \\*

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### **MERGEFORMAT**(1)

34 where  $r_{ij}$  is the distance between atom *i* and atom *j*.  $u_{ij}$  represents the interaction energy. 35  $\varepsilon_{ij}$  and  $\sigma_{ij}$  are energy parameters and length parameters, respectively.  $q_i$  is the charge of 36 the *i*-th particle and  $\varepsilon_0$  is the vacuum permittivity.

And the simulation results will be compared with DPMD simulations. Detailed
potential parameters are summarized in Table I, and the parameters between copper and
oxygen are determined by the Lorentz–Berthelot mixing rule<sup>1</sup>. The formula is as
follows:

$$\sigma_{Cu-O} = \frac{\sigma_{Cu} + \sigma_o}{2}$$

42 MERGEFORMAT (2)

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### 44 MERGEFORMAT (3)

45 The particle-particle mesh (PPPM) method<sup>2</sup> with an accuracy of  $1.0 \times 10^{-4}$  was served to 46 solve long-range interactions. The SHAKE algorithm was performed to ensure the rigidity of the 47 water molecules. As for the cutoff distance, it is set to 10 Å.

48 It is noted that CMD simulation adopts the same setup details as DPMD simulation.

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Table I. potential energy and related parameters between particles

Model	Parameters	Values	Units
	$\mathcal{E}_{O-O}$	0.00672	eV
TIP4P	$\sigma_{\scriptscriptstyle O-O}$	3.1536	Å
	$q_{\scriptscriptstyle H}$	0.52	e
	$q_o$	-1.04	e
	${\cal E}_{O=O}$	0.00673	eV
	$\sigma_{\scriptscriptstyle O-O}$	3.166	Å
SPC	$q_{\scriptscriptstyle H}$	0.41	e
	$q_o$	-0.82	e
Cu	$\mathcal{E}_{Cu-Cu}$	0.4096	eV
	$\sigma_{_{Cu-Cu}}$	0.233	nm

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# 52 SIII. Generation of the Training Sets by DP-GEN

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The training sets are generated using the DP-GEN software package, consisting of three steps, exploration, labeling, and training, as shown in Fig. S2. Initial training datasets are extracted from the AIMD simulation using PAW method and PBE functional. The AIMD simulations contains two steps: Firstly, the initial structure was optimized using the Conjugate-Gradient (CG) algorithm. Secondly, the structures were relaxed at 300 K in the canonical ensemble (NVT) with a time step of 1 fs. The cutoff energy for the plane wave expansion is set to be 550 eV, and the Brillouin zone is sampled at the  $\Gamma$  point. The SCF energy convergence threshold was chosen to be 10<sup>-6</sup> Hartree/atom. Simulation runs at 300K in the NVT ensemble, for 4000 steps with a time step of 0.5 fs. Based on the initial training sets, we trained an initial deep potential model for later configuration explorations and labeling. Next, new configurations from DPMD simulations will be added to the training sets for a new round of training. Finally, a relatively accurate and reliable potential model will be constructed after several above cycles.



to the equation (7) in the manuscript, we calculated the ITC at the interface. As shown in the Table
II, the ITC calculated using nosé-hoover thermostat is smaller than the results by langevin
thermostat.

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Fig. S3 The temperature distributions along the y direction calculated from dynamical simulations
with nosé-hoover thermostat and langevin thermostat, respectively. The table in the figure shows
the temperature drop at the high and low temperature surfaces.

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Table II. ITC predicted using different thermostat. ITC is in units of  $W/m^2 \cdot K$ .

Thermostat	langevin thermostat	nosé-hoover thermostat		
High-temperature	2 17×10 <sup>8</sup>	1 15×10 <sup>8</sup>		
interface	2.17*10	1.15~10		
Low-temperature	$2.84 \times 10^{8}$	$1.67 \times 10^{8}$		
interface	2.01.10	1.07 10		

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### 90 Temperature Effect

91 To investigate the effect of temperature on ITC, we calculated the ITC at 290 K, 300K and 310 K,
92 respectively. The bottom plate was heated up to 320 K, 310K and 300 K, respectively to create a
93 temperature gradient. The temperature of the top plate remained at 300 K, 290 K and 280 K,
94 respectively. Note that the ITC at 300 K is calculated by the temperature difference of 20 K, which
95 is different from the result in the manuscript (the temperature difference of 40 K). Fig. S4 shows
96 the temperature profiles and temperature drop. As shown in the Table III, the ITC increases with
97 rising temperature. Specially, the ITC of high-temperature interface at 300 K is less than that of 290
98 K and 310 K, and the ITC of low-temperature interface at 300 K lies between 290K and 310K.
99 The reason for this is the difference in temperature difference and the statistical averaging process.
100 Combined with the results in the manuscript, in our view, the ITC should follow the trend of



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103 Fig. S4 The temperature distributions along the y direction calculated from dynamical simulations104 at 290 K, 300 K and 310 K, respectively. The table in the figure shows the temperature drop at the

- 105 high and low temperature surfaces.
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Table III. ITC predicted at different temperature. ITC is in units of  $W/m^2 \cdot K$ .

Temperature	290 K	300 K	310 K	This work (300 K)	
High-temperature	$2.02 \times 10^{8}$	1 95×10 <sup>8</sup>	2 29×10 <sup>8</sup>	$2.17 \times 10^{8}$	
interface	2.02~10	1.95~10	2.27~10	2.17~10	
Low-temperature	2 25×10 <sup>8</sup>	$2.48 \times 10^{8}$	3 52×10 <sup>8</sup>	2 84×10 <sup>8</sup>	
interface	2.23 10	2.10 10	5.52 10	2.01 10	
Average value	(2.135 ±	(2.215 ±	$(2.905 \pm$	$(2.505 \pm$	
Average value	0.115)×10 <sup>8</sup>	0.265)×10 <sup>8</sup>	0.615)×10 <sup>8</sup>	0.335)×10 <sup>8</sup>	

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### 109 System Size

110 Size effect has been a challenge for NEMD simulations. Here, we investigate the effect of different

111 cross-sectional area and system length on interfacial thermal conductance. The main findings are 112 presented in Table IV and Fig. S5. As shown in the Table IV, the average ITC and the ITC at low-113 temperature interface tend to increase with rising cross-sectional area and to decrease with rising 114 channel height. As for the ITC at high-temperature interface, the trend is not particularly obvious, 115 and the additional calculations are needed to perform.





Fig. S5 (a) The temperature distributions along the y direction calculated from dynamical simulations with different cross-sectional area. (b) the temperature distributions along the y direction calculated from dynamical simulations with different channel height. The table in the figure shows the temperature drop at the high and low temperature surfaces.

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Table IV. ITC predicted with different size. ITC is in units of W/m<sup>2</sup>·K

Size	Cross-sectional area (Å <sup>2</sup> )			Channel Height (Å)		
ІТС	39×10	39 ×25	39 ×34	31	39.5	46.5
High-temperature interface	2.02×10 <sup>8</sup>	1.95×10 <sup>8</sup>	2.73×10 <sup>8</sup>	2.41×10 <sup>8</sup>	2.04×10 <sup>8</sup>	2.17×10 <sup>8</sup>
Low-temperature	2.25×10 <sup>8</sup>	2.48×10 <sup>8</sup>	4.70×10 <sup>8</sup>	4.09×10 <sup>8</sup>	3.67×10 <sup>8</sup>	2.84×10 <sup>8</sup>

interface

Average value	$(2.135 \pm$	(2.215 ±	(3.715 ±	$(3.25 \pm$	$(2.855 \pm$	$(2.505 \pm$
	0.115)×10 <sup>8</sup>	0.265)×10 <sup>8</sup>	0.985)×10 <sup>8</sup>	0.840)×10 <sup>8</sup>	0.815)×10 <sup>8</sup>	0.335)×10 <sup>8</sup>

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120 In summary, it can be concluded that size, temperature and thermostat all expert an effect on the

121 interfacial thermal conductance, and more exploration is needed. In addition, previous

122 computational and experimental studies demonstrated a proportional relation between work of solid-

123 liquid adhesion or solid-liquid interaction strength and interfacial thermal conductance (ITC). This

- 124 presents the idea to design different coating surfaces for modulating the interfacial thermal
- 125 conductance. Then, a reliable interfacial thermal conductance can be calculated based on the
- 126 accurate machine learning potential.
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