

*Supporting Information for*

**Redispersion Mechanisms of 2D Nanosheets: Combined  
Role of Intersheet Contact and Surface Chemistry**

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## 1. Supplementary Experimental Section

**Evaluation of Hamaker Constants.** To evaluate the vdW forces between nanosheets of pristine and transformed GO and MoS<sub>2</sub>, the Hamaker constants in water were evaluated. Based on the Lifshitz theory for dispersion forces, the Hamaker constant between two semi-infinite media across a medium (water) can be estimated as follows:

$$H_{total} \approx \frac{3hv_e}{8\sqrt{2}} \times \frac{(n_1^2 - n_3^2)(n_2^2 - n_3^2)}{\sqrt{n_1^2 + n_3^2} \times \sqrt{n_2^2 + n_3^2} \times (\sqrt{n_1^2 + n_3^2} + \sqrt{n_2^2 + n_3^2})}$$

where  $n_1 = n_2 = n$  is the refractive index of interacting nanosheets in visible regime,  $n_3$  is the refractive index of water (1.33),  $h$  is the Planck constant ( $6.626 \times 10^{-34}$  Js), and  $v_e$  is the main absorption frequency in the UV region, which could be obtained through the following equation:

$$v_e = v_I \sqrt{3/(n^2 + 2)}$$

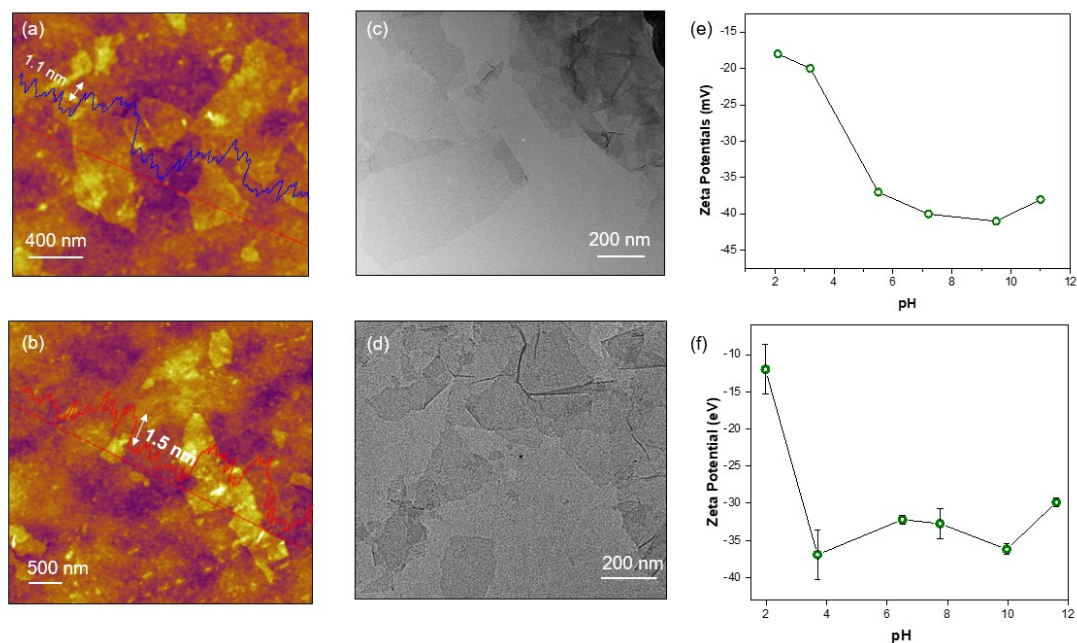
where  $v_I$  is the absorption frequency of a Bohr atom ( $3.3 \times 10^{15}$  s<sup>-1</sup>).

To the best of our knowledge, there's no data available for the refractive index  $n$  of GO, MoS<sub>2</sub> nanosheets and transformed ones. However, it was reported the average  $n$  of GO and rGO in visible regime are around 1.85 and 2.6, respectively;<sup>1,2</sup> and the average  $n$  of monolayer 1T MoS<sub>2</sub> and 2H MoS<sub>2</sub> are around 2.5 and 4.0, respectively.<sup>3,4</sup> Therefore, the  $H$  for GO and rGO in water is thus between 49 and  $260 \times 10^{-21}$  J, and  $H$  of 1T-MoS<sub>2</sub> and 2H-MoS<sub>2</sub> is between 154 and  $320 \times 10^{-21}$  J. A linear increase of  $H$  was assumed during reduction of GO and phase transformation of MoS<sub>2</sub>.

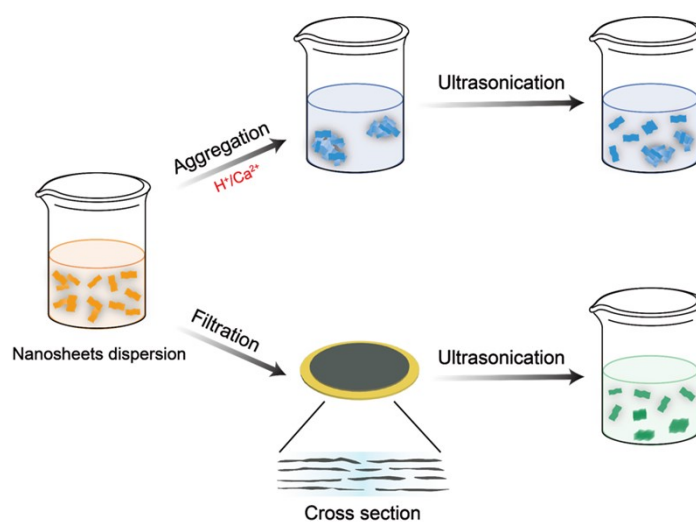
**First-principles calculations.** First-principles calculations were performed in the framework of the density functional theory (DFT) method as implemented in the Vienna Ab Initio Simulation Package (VASP).<sup>5,6</sup> The project-augmented-wave (PAW) method<sup>7,8</sup> with plane wave basis sets were employed to treat the core-electron interaction. The generalized gradient approximation (GGA) in the Perdew-Burke-Ernzerhof (PBE)<sup>9</sup> form was applied to describe the exchange and correlation energy.

The energy cutoff was 400 eV. The Monkhorst-Pack scheme with a  $1 \times 1 \times 1$  k-point grid was adopted to sample the Brillouin zone.<sup>10</sup> All the atomic structures were relaxed until the force was below 0.05 eV/Å. A model consisting of two  $35 \times 39$  Å<sup>2</sup> MoS<sub>2</sub> nanoribbons with a tilt angle was used to model the variation of the system energy with the interlayer distance (see Figure S16d). A  $\sim 15$  Å vacuum layer was used to avoid interaction from neighboring cells in both the  $x$  and  $z$  directions.

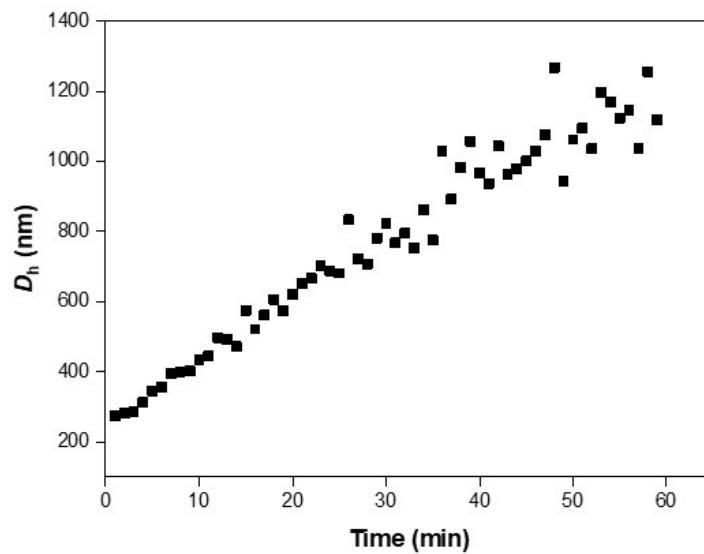
## 2. Supplementary Results



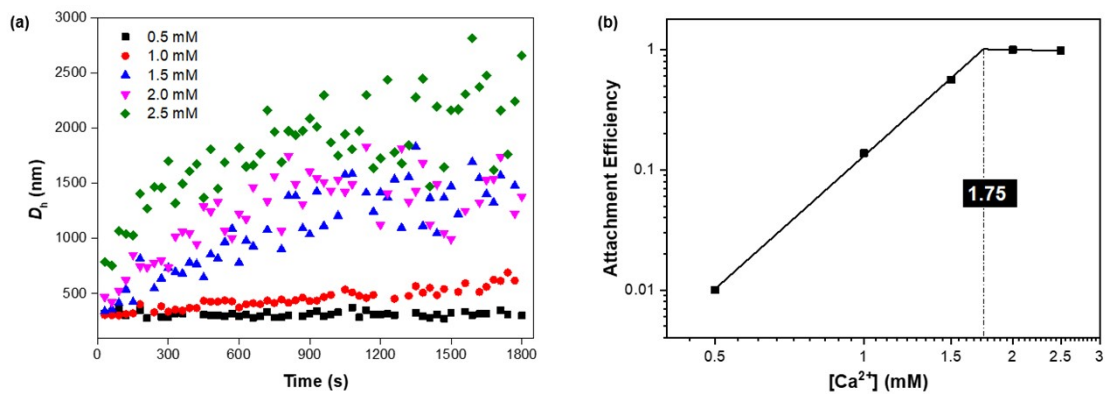
**Figure S1.** Characterization of MoS<sub>2</sub> (a, c, e) and GO (b, d, f) nanosheets: (a-b) representative AFM image, inset: line scan showing the thickness profile along the red line in the image; (c-d) representative TEM image; (e-f) apparent  $\zeta$  potentials as a function of pH.



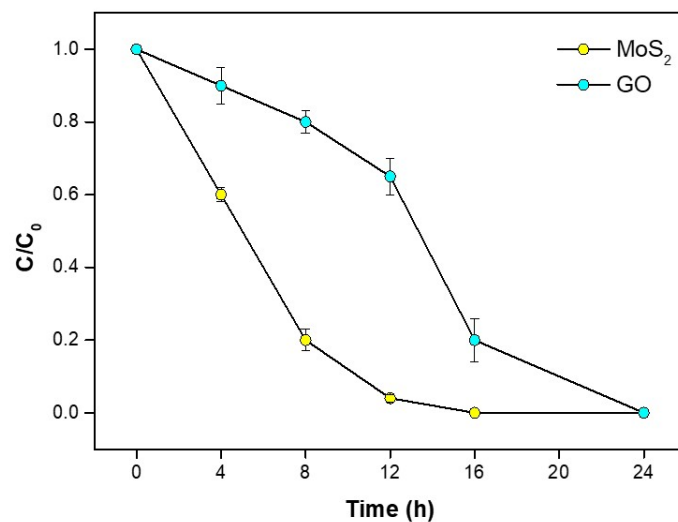
**Figure S2.** Schematic illustration for the creation and the redispersion of random aggregates and aligned stacks.



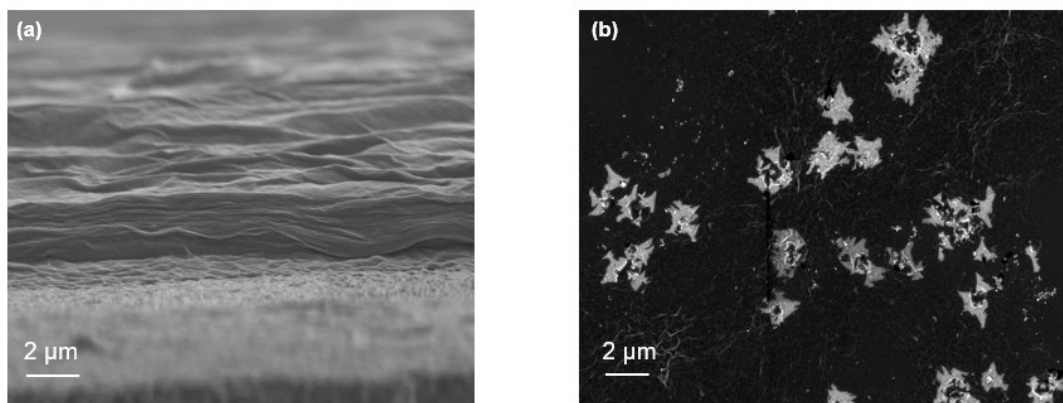
**Figure S3.** Rising hydrodynamic diameter of MoS<sub>2</sub> dispersion at pH 3. The results suggested that a low pH could induce the aggregation of the nanosheets and thus destroy the colloidal stability by providing adequate H<sup>+</sup>.



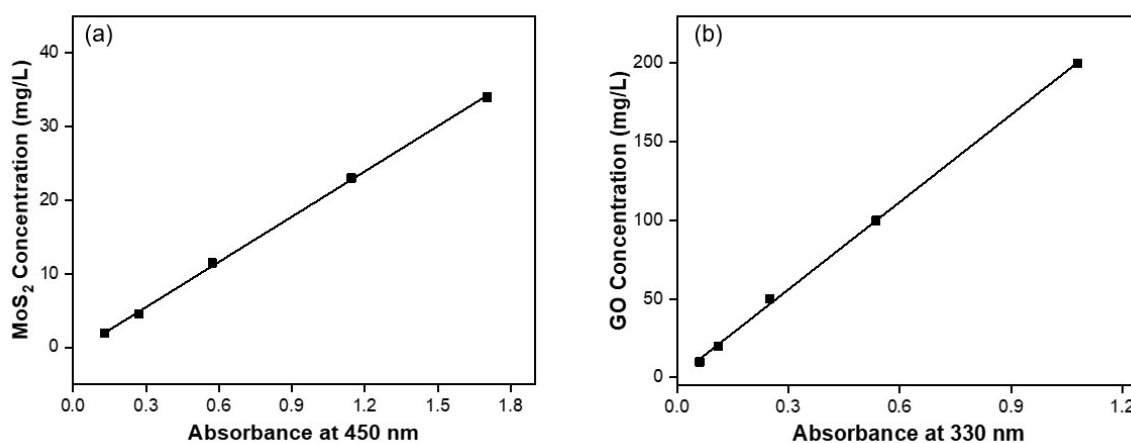
**Figure S4.** Aggregation of GO dispersion: (a) evolution of hydrodynamic diameter of GO in the solutions containing a series of concentrations of Ca<sup>2+</sup>; (b) attachment efficiency as a function of Ca<sup>2+</sup> concentrations, from which the critical coagulation concentration (CCC) of 1.75 mM was obtained. Therefore, a concentration of 2 mM Ca<sup>2+</sup> was employed in the study to induce the nanosheets aggregation, which is higher than CCCs of Ca<sup>2+</sup> for GO (1.75 mM) and MoS<sub>2</sub> (0.9 mM).<sup>11</sup>



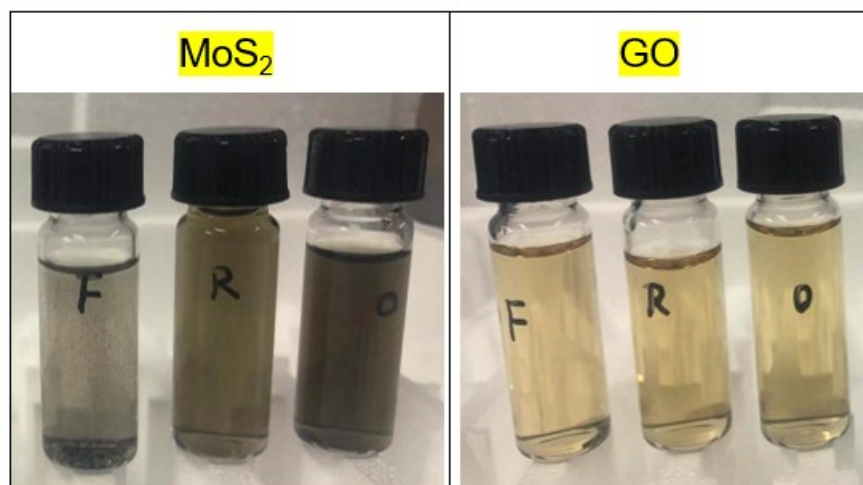
**Figure S5.** Evolution of the normalized suspended concentration of MoS<sub>2</sub> and GO in the solution containing 2 mM Ca<sup>2+</sup>.



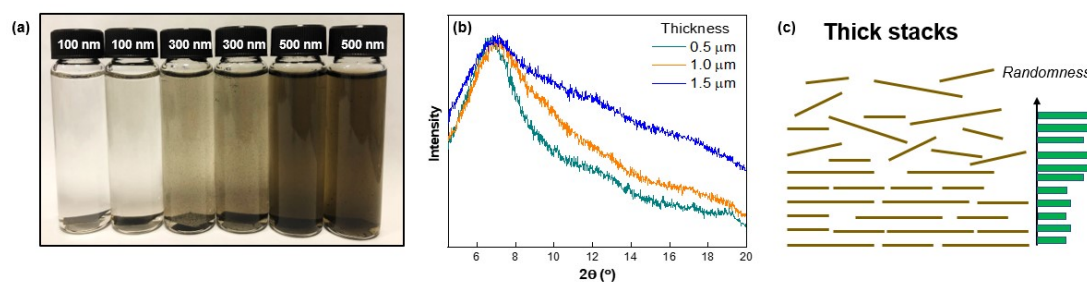
**Figure S6.** SEM characterization of aligned stacks and aggregates of GO: (a) cross-section image of GO stacks obtained by filtration and (b) top view of GO aggregates.



**Figure S7.** Linear correlation of the absorbance to the concentration of nanosheets dispersion: (a) SL-MoS<sub>2</sub> at 450 nm and (b) GO at 330 nm.

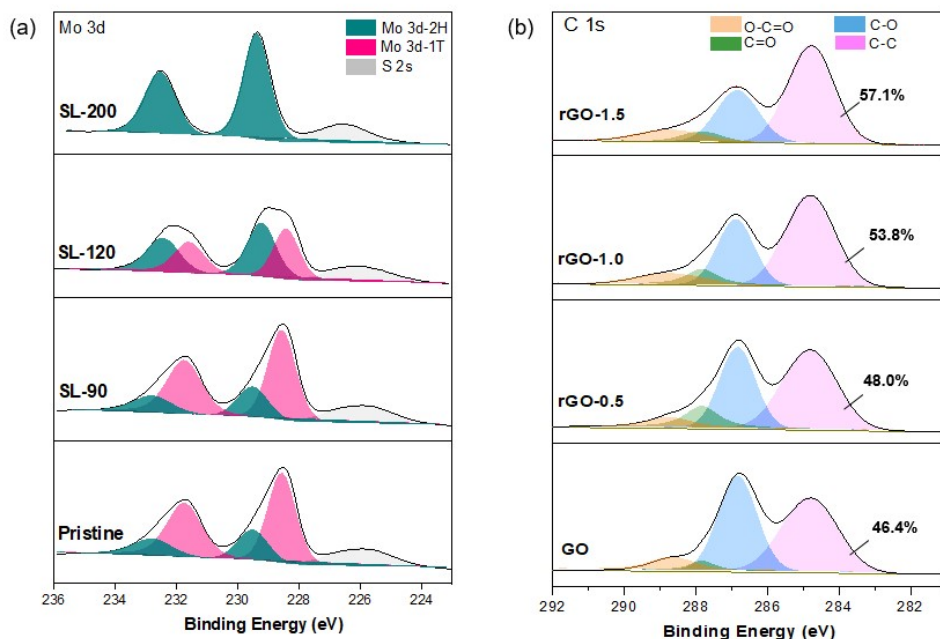


**Figure S8.** Photographs of the original dispersion and redispersion obtained from the random aggregates and aligned stacks. Left: MoS<sub>2</sub>; right: GO (F: redispersion of stacks created by filtration; R: redispersion from random aggregates; O: original dispersion).

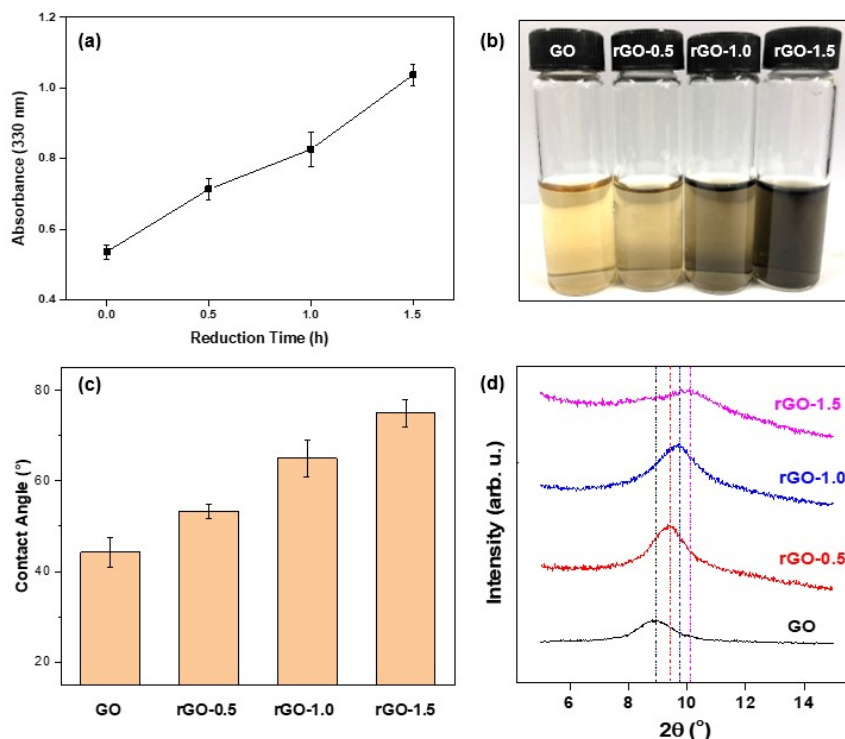


**Figure S9.** Redispersion from MoS<sub>2</sub> stacks as a function of thickness: (a) images of redispersion, with redispersion efficiency of 1%, 29% and 54% from the stacks with thickness of 100, 300 and 500 nm, respectively; (b) XRD patterns, showing a broadening peak as the increase of thickness; (c) schematic illustration of randomness as influenced by the thickness of MoS<sub>2</sub> stacks.

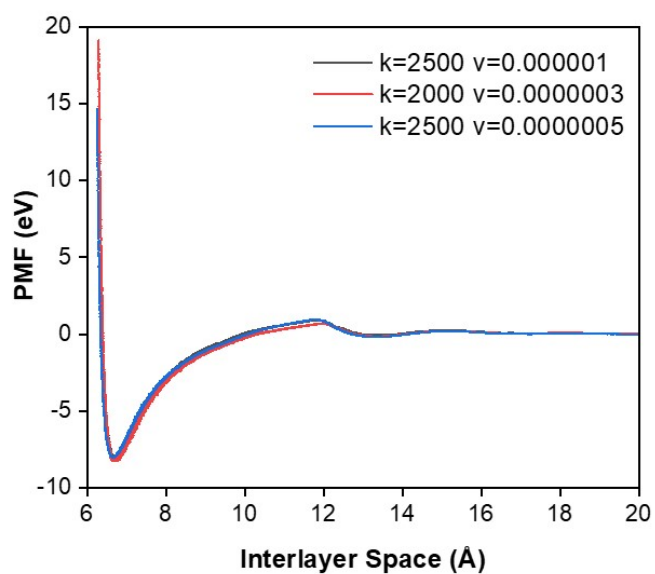




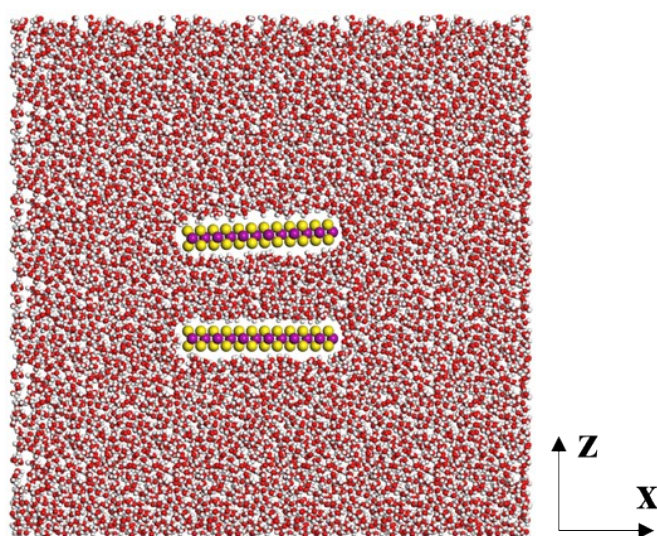
**Figure S10.** XPS characterizations on the composition deconvolution of MoS<sub>2</sub> and GO during transformation: (a) Mo 3d of MoS<sub>2</sub> nanosheets, showing the increasing fraction of 2H phase during hydrothermal treatment, and (b) C 1s of GO during the reduction reaction, suggesting the restoration of C–C/C=C during reaction.



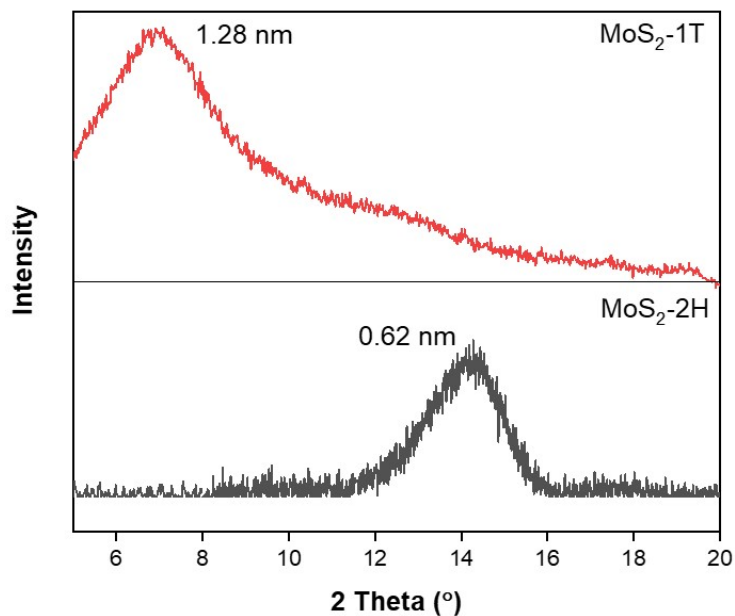
**Figure S11.** Characterization of GO and rGO obtained through chemical reduction: (a) absorbance at 330 nm in the concentration of 100 mg/L; (b) photograph of the dispersion (100 mg/L); (c) water contact angles and (d) XRD patterns of restacked GO and rGO.



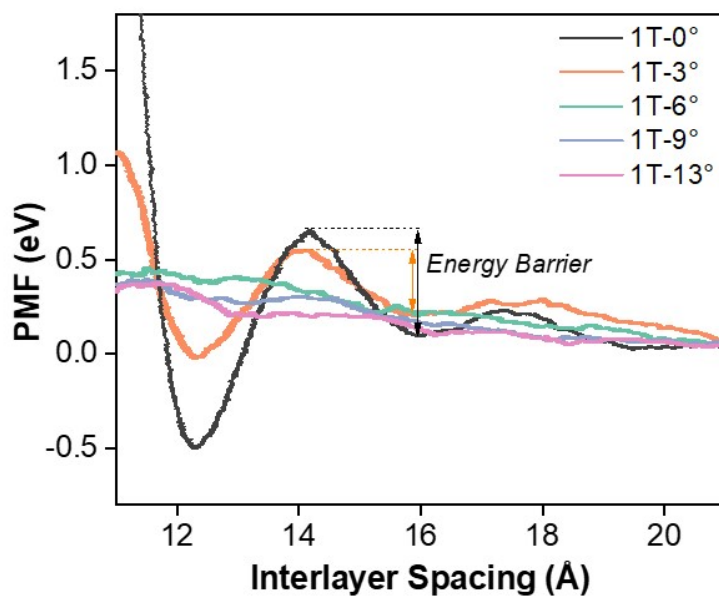
**Figure S12.** SMD simulation results of system at  $k_{\text{spring}} = 2500 \text{ kcal/mol/\AA}^2$  with various pulling speed from 0.0000005 to 0.000001  $\text{\AA}/\text{fs}$ .



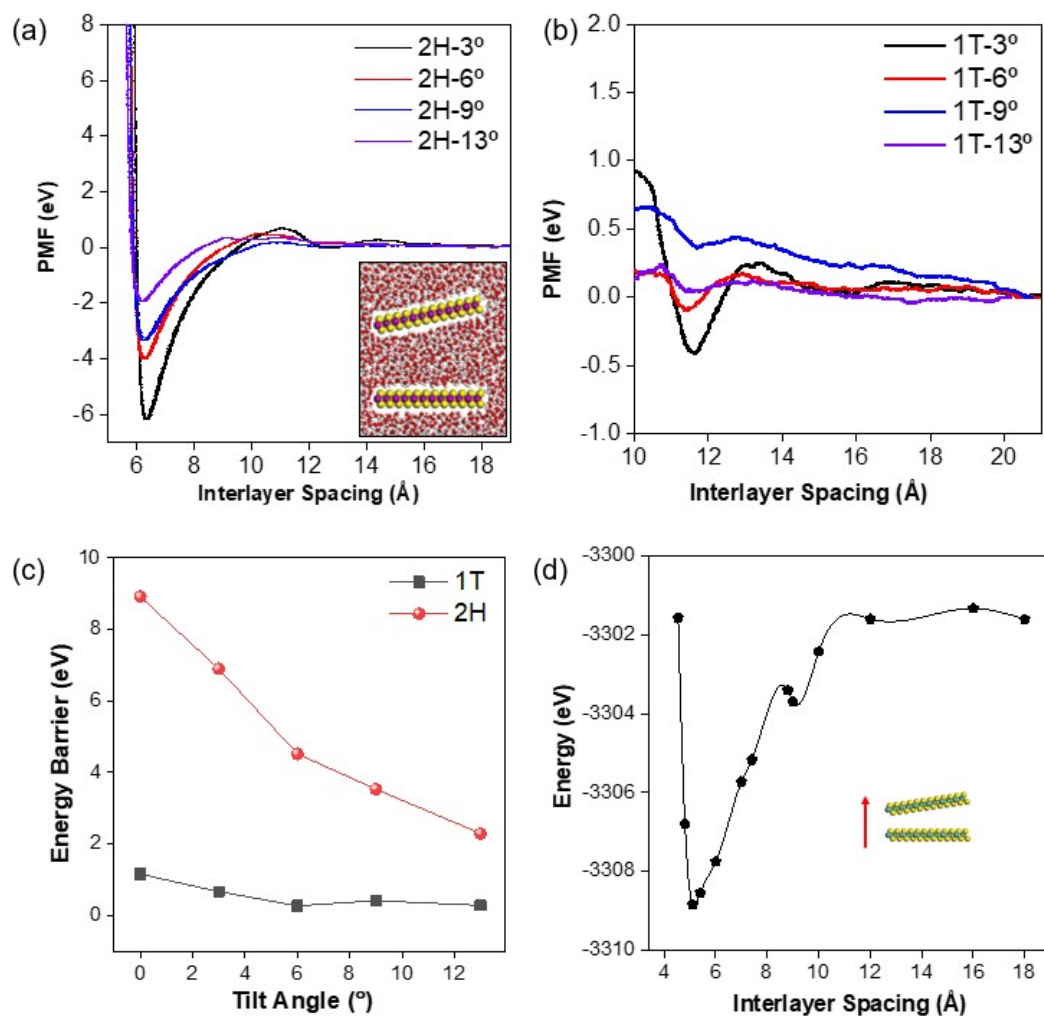
**Figure S13.** Snapshot of an MD simulation used to compute the PMFs showing the full extent of the simulation box in the x- and z-directions.



**Figure S14.** XRD patterns of as-created (wet) MoS<sub>2</sub> stacks from SL-pristine dominant in 1T phase and SL-200 in pure 2H phase, with the interlayer spacing of 1.28 and 0.62 nm identified, respectively.



**Figure S15.** PMF as a function of interlayer spacing for 1T phase of MoS<sub>2</sub> with various tilt angles. The simulated models are built with S atoms being exposed outside.



**Figure S16.** PMF calculations for 1T and 2H MoS<sub>2</sub> tilted at the Mo end: (a-b) the PMF profiles of 2H (a) and 1T (b) MoS<sub>2</sub> at various tilt angles (inset: the schematic illustration of the tilt model of two MoS<sub>2</sub> layers in water solution); (c) the redispersion barrier of 1T and 2H MoS<sub>2</sub> changes with the tilt angle; and (d) the total energy of the tilted 1T MoS<sub>2</sub> in vacuum (tilt angle=9°) changes with interlayer spacing based on first-principles calculations.

**Table S1** Water contact angle and Hamaker constant of nanosheets

<b>Nanosheets</b>	<b>Hamaker constant (<math>\times 10^{-21}</math> J)</b>	<b>Water contact angle (<math>^{\circ}</math>)</b>
<b>GO</b>	49	45
<b>rGO-0.5</b>	62	53.3
<b>rGO-1.0</b>	108	69
<b>rGO-1.5</b>	135	88
<b>MoS<sub>2</sub></b>	154	58
<b>MoS<sub>2</sub>-90</b>	159	60
<b>MoS<sub>2</sub>-120</b>	218	72
<b>MoS<sub>2</sub>-200</b>	320	98
<b>WS<sub>2</sub></b>	320 <sup>12</sup>	78 <sup>13</sup>
<b>h-BN</b>	764 <sup>12</sup>	81 <sup>14</sup>
<b>MXene</b>	48.7 <sup>15</sup>	24.8 <sup>16</sup>
<b>AlMg-LDH</b>	14 <sup>17</sup>	21.4 <sup>18</sup>

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