

## SUPPORTING INFORMATION

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# Electronic Supplementary Information (ESI)

## Tellurate Polymorphs with High-performance Nonlinear Optical Switch Property and Wide Mid-IR Transparency

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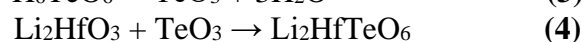
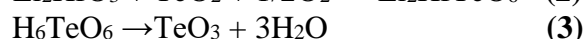
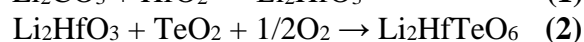
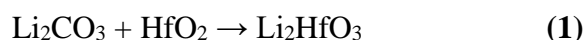
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## Experimental Procedures

**Synthesis.** All the reagents are analytical grade and gained from Aladdin without further processing and purification. Single crystals of  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub> were obtained by using the high-temperature flux method. The mixture of Li<sub>2</sub>CO<sub>3</sub>, HfO<sub>2</sub>, and TeO<sub>2</sub> in a molar ratio of 3:1:5 was fully ground and placed in a platinum crucible. The crucible was heated to 750 °C in 3 h, and held at this temperature for 24 h. Afterward, the solution was cooled slowly to 620 °C with a rate of 2 °C·h<sup>-1</sup>, and then cooled down to room temperature quickly. Colorless crystals of  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub> were obtained from the crushed products. However, under similar conditions, the attempts to obtain single crystals of  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub> were failed due to poor crystallinity.

The polycrystalline samples of  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub> were obtained by the solid-state reaction method in the air in three steps. First, according to equation 1, Li<sub>2</sub>CO<sub>3</sub> and HfO<sub>2</sub> with a stoichiometric ratio were weighted and fully grounded. The mixture was loaded into a corundum crucible and heated to 1000 °C at a rate of 5 °C·min<sup>-1</sup>, held at this temperature for 14 hours to obtain Li<sub>2</sub>HfO<sub>3</sub>. Second, Li<sub>2</sub>HfO<sub>3</sub> and TeO<sub>2</sub> combine with oxygen from the air to create Li<sub>2</sub>HfTeO<sub>6</sub> as suggested by equation 2. The mixture was heated at 700 °C for 12 h and 650 °C for 12 h, respectively, with several intermittent grindings. At this stage, a weak SHG signal can be detected from the as-obtained products, suggesting that both  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub> and  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub> exist. Finally, to prepare single phase of  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub>, the sample was heated again at 550 °C for 2 d and slowly cooled to room temperature in 24 h.



Powder samples of  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub> were prepared in a similar way with some differences in Te source and temperature control. Step 1: Li<sub>2</sub>HfO<sub>3</sub> was pre-synthesized as described above. Step 2: H<sub>6</sub>TeO<sub>6</sub> was calcined at 400 °C for 12h to synthesize TeO<sub>3</sub> (equation 3). Step 3, the as-obtained Li<sub>2</sub>HfO<sub>3</sub> and TeO<sub>3</sub> in the equal molar ratio were calcined at 750 °C for 48 h with several intermittent grindings (equation 4), and the sample was cooled to the room temperature by taking it out of the furnace directly and cooled in ice water to bypass the phase transition. The phase purity was checked by XRD and SHG tests.

**Single-Crystal XRD.** The crystal data were collected on a Bruker D8 VENTURE diffractometer. The diffractometer was equipped with a microfocus sealed X-ray tube and used Mo *K*<sub>α</sub> radiation ( $\lambda = 0.71073$  Å). The data reduction and multi-scan absorption correction were performed using the SAINT and SADABS, respectively. Within Olex2 software,<sup>1</sup> the structure was solved and refined by ShelXT<sup>2</sup> and ShelXL,<sup>3</sup> respectively. The crystal symmetry was checked by PLATON.<sup>4</sup> The information of crystal data and detailed refinement are summarized in Table S1. The atomic coordinates, isotropic displacement parameters, and bond valence sums of each atom are presented in Table S2. The selected bond lengths are given in Table S3.

**Powder XRD.** Powder XRD measurements were performed on a D8 advance X-ray diffractometer with Cu *K*<sub>α</sub> radiation ( $\lambda = 1.5408$  Å) at room temperature. The diffraction patterns were scanned with the  $2\theta$  range from 10° to 70°. The scan step width was 0.02° and the scan time was 2-5 s. Rietveld refinement of the structural model was performed by GSAS-II software.<sup>5</sup>

**Thermal Analysis.** The thermal gravimetric (TG) analysis and differential scanning calorimetry (DSC) of Li<sub>2</sub>HfTeO<sub>6</sub> were measured with a NETZSCH5 instrument under the air. The sample was heated from 30 °C to 1100°C with a heating rate of 10K/min at an alumina crucible.

Supporting Information

**Scanning Electronic Microscope (SEM) Images and Energy Dispersive X-Ray (EDX) Microanalysis.**

SEM images and EDX microanalysis of the crystals were performed on a Hitachi TM4000Plus microscopes with acceleration voltage of 15 kV. In addition, the elemental mapping images were recorded to show the distribution of chemical elements.

**UV-Vis-NIR Diffuse-Reflectance Spectrum.** Polycrystalline samples were used for diffuse-reflectance spectrum measurements. A METASH UV-6100 spectrophotometer was used to measure the diffuse reflectance spectrum in the 200-1100 nm wavelength range at room temperature.

**Infrared Spectroscopy.** The IR spectrum was measured by Shimadzu IR Affinity-1 Fourier transform infrared spectrometer at room temperature in the range of 400-4000  $\text{cm}^{-1}$ . Pellets for measurement were prepared by pressing a thoroughly ground mixture of 5 mg sample and 500 mg dried KBr.

**Dielectric Spectrum.** Polycrystalline samples of  $\alpha\text{-Li}_2\text{HfTeO}_6$  were cold-pressed into a pellet (~12 mm diameter and ~1.3 mm thick), and then sintered at 700 °C for 48 h. The silver paste was applied to both sides of the sintered pellet. The temperature-dependent dielectric spectrum was recorded with a TH2827A Precision LCR Meter. The changes of the dielectric coefficient at variable temperatures (50 to 800 °C, at a heating rate of 2 °C·min<sup>-1</sup>) were recorded in the frequency range from 10 KHz to 1 MHz.

**SHG Measurement.** The SHG effects of the powder samples were measured by using a modified Kurtz and Perry method with a Q-switched 1064 nm Nd: YAG laser.<sup>6</sup> The samples were ground and divided into fractions according to different particle size ranges: 20-38.5, 38.5-55, 55-80, 80-125, 125-160, 160-200, and 200-250  $\mu\text{m}$ . The same particle-sized KDP samples were used as the standard. Temperature-dependent analysis of SHG effects was performed using ground powder from 50 to 700 °C. The SHG switching contrast is defined as the ratio of the SHG intensities at the ON and OFF status near the phase transition temperature.

**Ferroelectric Measurements.** The polarization of  $\alpha\text{-Li}_2\text{HfTeO}_6$  was measured at room temperature by using a ZT-1B ferroelectric test system. The measurements were performed at a fixed frequency of 10 Hz with electric fields ranging from 25 to 100 kV  $\text{cm}^{-1}$ .

**LDT Measurement.** The optical damage induced by laser was tested on crystalline powder pellet samples using a Q-switched 1064 nm Nd: YAG laser (1064 nm, 7 ns, 5 Hz). The pulse energy was changed from 20-100 mJ and a convex lens was used to adjust the beam diameter. The damaged area can be directly observed under a microscope.

**Theoretical Calculation.** Theoretical calculations were performed by using the CASTEP package based on density functional theory (DFT).<sup>7, 8</sup> Perdew-Burke-Ernzerhof generalized gradient approximation (PBE-GGA) functional was adopted to describe the exchange-correlation energy.<sup>9</sup> The energy cut-off for the plane-wave basis was set to be 750 eV. The Norm-conserving pseudopotential (NCP) was adopted with following valence electrons, Li ( $2s^1$ ), Hf ( $5d^26s^2$ ), Te ( $5s^25p^4$ ), O ( $2s^22p^4$ ). The numerical integration of the Brillouin zone was performed using a Monkhorst-Pack scheme with  $k$ -point grid density of  $3 \times 3 \times 2$ . In addition, a scissors operator was adopted to correct the underestimation of band gap induced by PBE-GGA. On the basis of the scissor-corrected electron band structure, the second order NLO coefficients  $d_{ij}$  were calculated using the length-gauge formalism at a zero-frequency limit.<sup>10</sup> The static second-order nonlinear susceptibilities can be ascribed to virtual-electron and virtual-hole processes. In addition, the SHG-weighted electron density method was adopted to visualize the contribution of each atom to the SHG response.<sup>11, 12</sup>

**Results and Discussion****Table S1.** Crystal data and structure refinement for  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub> and  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub>.

Empirical formula	$\alpha$ -Li <sub>2</sub> HfTeO <sub>6</sub>	$\beta$ -Li <sub>2</sub> HfTeO <sub>6</sub>	
Formula weight	415.97	Formula weight	415.97
Crystal system	Trigonal	Crystal system	Trigonal
Space group (number)	<i>R</i> 3 (146)	Space group (number)	$R\bar{3}$ (148)
<i>a</i> [Å]	5.1615	<i>a</i> [Å]	5.153(3)
<i>b</i> [Å]	5.1615	<i>b</i> [Å]	5.153(3)
<i>c</i> [Å]	13.7822	<i>c</i> [Å]	13.753(14)
Volume [Å <sup>3</sup> ]	317.9	Volume [Å <sup>3</sup> ]	316.3(5)
<i>Z</i>	3	<i>Z</i>	3
$\rho_{\text{calc}}$ [gcm <sup>-3</sup> ]	6.516	$\rho_{\text{calc}}$ [gcm <sup>-3</sup> ]	6.552
Radiation	CuK $\alpha$ ( $\lambda = 1.5408$ Å)	$\mu$ [mm <sup>-1</sup> ]	31.454
Instrument type	Bragg-Brentano	<i>F</i> (000)	534
<i>R</i> <sub>wp</sub>	0.0889	Crystal colour	colourless
<i>R</i> <sub>p</sub>	0.0690	Crystal shape	block
Phase fraction	0.980	Radiation	MoK $\alpha$ ( $\lambda = 0.71073$ Å)
		2 $\theta$ range [°]	8.89 to 55.01 (0.77 Å)
		Index ranges	-6 ≤ <i>h</i> ≤ 6, -6 ≤ <i>k</i> ≤ 5, -17 ≤ <i>l</i> ≤ 16
		Reflections collected	1193
		Independent reflections	165, <i>R</i> <sub>int</sub> = 0.0478, <i>R</i> <sub>sigma</sub> = 0.0275
		Completeness	100.0 %
		Data / Restraints / Parameters	165/12/19
		Goodness-of-fit on <i>F</i> <sup>2</sup>	1.311
		Final <i>R</i> indexes [ <i>I</i> ≥ 2σ( <i>I</i> )]	<i>R</i> <sub>1</sub> = 0.0257, w <i>R</i> <sub>2</sub> = 0.0586
		Final <i>R</i> indexes [all data]	<i>R</i> <sub>1</sub> = 0.0279, w <i>R</i> <sub>2</sub> = 0.0589
		Largest peak/hole [eÅ <sup>-3</sup> ]	1.75/-1.28
		Extinction coefficient	0.0051(8)

## Supporting Information

**Table S2.** Atomic coordinates and  $U_{eq}$  [ $\text{\AA}^2$ ] for  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub> and  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub>.

Atom	$x$	$y$	$z$	$U_{eq}^{[a]}$	BVS <sup>[b]</sup>
<b><math>\alpha</math>-Li<sub>2</sub>HfTeO<sub>6</sub></b>					
Li(1)	2/3	1/3	0.4395	0.0487	0.84
Li(2)	0	0	0.5934	0.0064	0.92
Hf(1)	0	0	0.3759	0.0044	4.34
Te(1)	1/3	2/3	0.5456	0.0043	6.11
O(1)	0.3527	0.0581	0.2966	0.0368	2.05
O(2)	0.2844	0.33796	0.4697	0.0020	2.01
<b><math>\beta</math>-Li<sub>2</sub>HfTeO<sub>6</sub></b>					
Li(1)	1	1	0.232	0.0660	0.84
Hf(1)	2/3	1/3	1/3	0.0060	4.23
Te(1)	1/3	2/3	1/6	0.0062	5.92
O(1)	0.6200	0.6321	0.2469	0.0157	1.97

<sup>[a]</sup>  $U_{eq}$  is defined as 1/3 of the trace of the orthogonalized  $U_{ij}$  tensor.

<sup>[b]</sup> Bond valence sums (BVS) are calculated by using bond-valence theory ( $S_i = \exp[(R_o - R_i/B]$ , where  $R_o$  is an empirical constant,  $R_i$  is the length of bond  $I$  (in angstroms), and  $B = 0.37$ ).

**Table S3.** Selected bond lengths [ $\text{\AA}$ ] for  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub> and  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub>.

<b><math>\alpha</math>-Li<sub>2</sub>HfTeO<sub>6</sub></b>		<b><math>\beta</math>-Li<sub>2</sub>HfTeO<sub>6</sub></b>	
Li(1)–O(1) $\times_3$	2.4939	Li(1)–O(1) $\times_3$	1.939
Li(1)–O(2) $\times_3$	2.0281	Li(1)–O(1) $\times_3$	2.600
Li(2)–O(1) $\times_3$	2.0356	Hf(1)–O(1) $\times_6$	2.053
Li(2)–O(2) $\times_3$	2.3553	Te(1)–O(1) $\times_6$	1.922
Hf(1)–O(1) $\times_3$	2.0127		
Hf(1)–O(2) $\times_3$	2.0756		
Te(1)–O(1) $\times_3$	1.9214		
Te(1)–O(2) $\times_3$	1.9001		

## Supporting Information

**Table S4.** Comparison of structure and optical properties of compounds in  $\text{Li}_2\text{MTeO}_6$  (M = tetravalent cation) family.

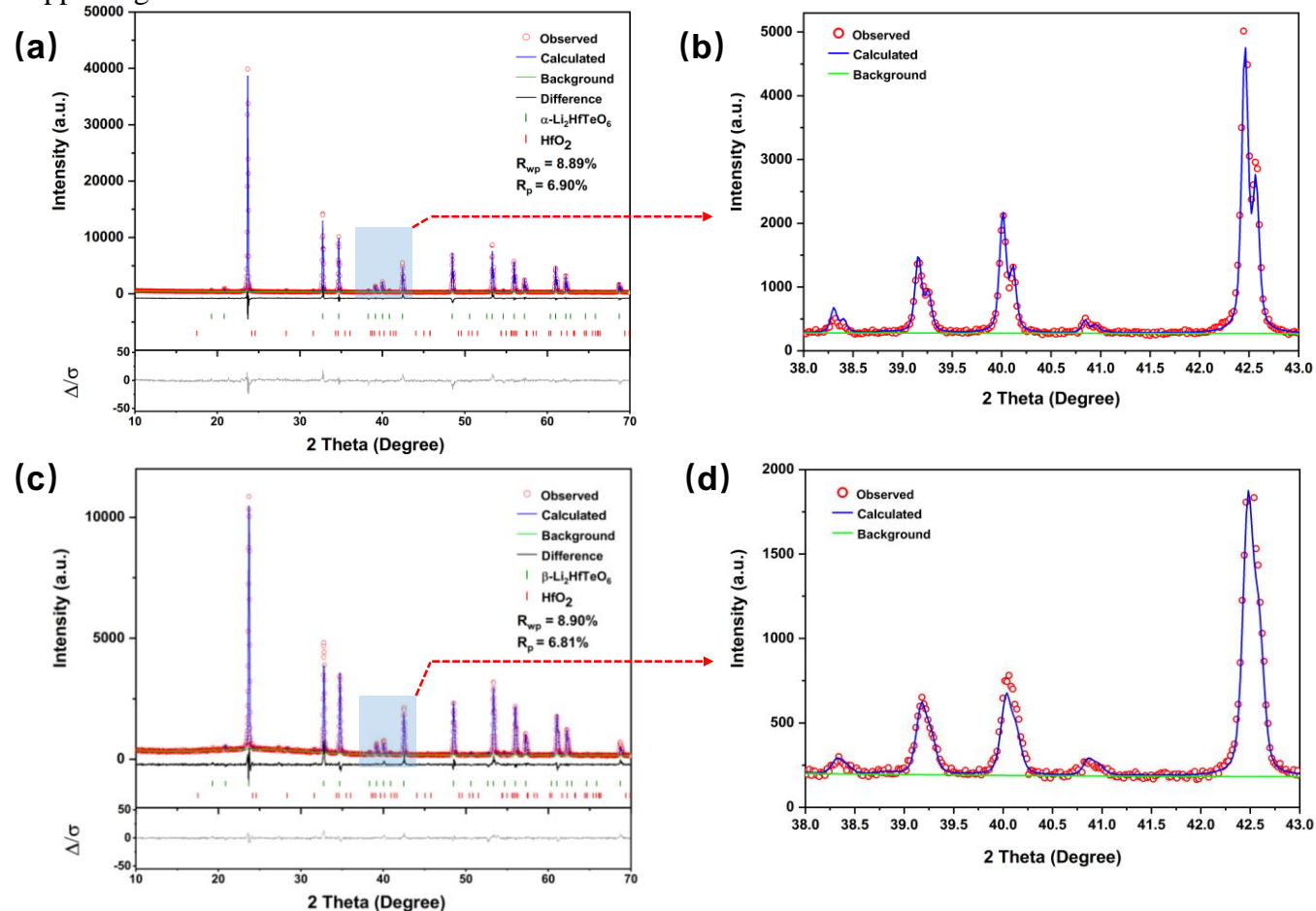
	space group	octahedral distortion $\Delta d$		SHG ( $\times$ KDP)	transparent region ( $\mu\text{m}$ )	LDT ( $\text{MW cm}^{-2}$ )
		$\text{MO}_6$	$\text{TeO}_6$			
$\text{Li}_2\text{TiTeO}_6$	$Pnn2$	0.20	0.06	26	0.38-6.72	>550
$\text{Li}_2\text{ZrTeO}_6$	$R3$	0.12	0.02	2.5	0.29-7.4	>1300
$\alpha\text{-Li}_2\text{HfTeO}_6$	$R3$	0.19	0.06	2.2	0.27-8.6	>1000
$\beta\text{-Li}_2\text{HfTeO}_6$	$R\bar{3}$	0	0	/	0.26-8.6	>1000
$\text{Li}_2\text{SnTeO}_6$	$Pnn2$	0.01	0.09	2.5	0.38-6.86	>672

**Table S5.** Displacive mode amplitudes [ $\text{\AA}$ ] from  $\beta\text{-Li}_2\text{HfTeO}_6$  to  $\alpha\text{-Li}_2\text{HfTeO}_6$ .

mode	atom	$A_s$	$A_p$	$d_{\text{max}}$	Description
$\Gamma_1^+$	Li(1)_A(a)	0.22633	0.22633	0.16004	
	O(1)_A_1(a)	-0.10135	-0.10135	0.04138	
	O(1)_A_2(a)	-0.05727	-0.05727	0.02338	Shift of Li and O atoms, CS
	O(1)_A_3(a)	-0.00292	-0.00292	0.00119	
	all	<b>0.25453</b>	<b>0.25453</b>	/	
$\Gamma_1^-$	Li(1)_A(a)	0.00000	0.00000	0.00000	/
	Hf(1)_Au(a)	0.35849	0.35849	0.35849	Displacement of Hf atom, NCS
	Te(1)_Au(a)	0.40131	0.40131	0.40131	Displacement of Te atom, NCS
	O(1)_A_1(a)	0.05569	0.05569	0.02274	Shift of O atoms, NCS
	O(1)_A_2(a)	-0.06093	-0.06093	0.02487	
	O(1)_A_3(a)	1.12304	1.12304	0.45848	Rotation of $\text{HfO}_6/\text{TeO}_6$ octahedron, NCS
	all	<b>1.24804</b>	<b>1.24804</b>	/	
Overall		<b>1.27373</b>	<b>1.27373</b>	/	

Note: The magnitude of  $A_s$  is the square root of the sum of the squares of the mode-induced changes within the primitive supercell (i.e. the root-summed-squared displacement).  $A_p = A_s \cdot \sqrt{V_p/V_s}$  to be normalized to the parent cell, where  $V_p$  and  $V_s$  are the respective primitive parent and supercell volumes.  $d_{\text{max}}$  is the maximum displacive mode amplitude. CS = centrosymmetric transform; NCS = noncentrosymmetric transform.

## Supporting Information



**Figure S1.** Rietveld refinement of the powder XRD data for (a)-(b)  $\alpha$ - $\text{Li}_2\text{HfTeO}_6$  and enlarged area from 38 to 43°, and (c)-(d)  $\beta$ - $\text{Li}_2\text{HfTeO}_6$  and enlarged area from 38 to 43°. Note that a small quantity of  $\text{HfO}_2$  (1.9% and 1.6%, for  $\alpha$ - $\text{Li}_2\text{HfTeO}_6$  and  $\beta$ - $\text{Li}_2\text{HfTeO}_6$ , respectively) coexists as the impurity, which cannot be avoided after plenty of attempts. On first sight, the two polymorphs show almost the same XRD pattern. However, when zoom in on curves between 38 to 43°, one can see that several peaks of  $\alpha$ - $\text{Li}_2\text{HfTeO}_6$  are split while those of  $\beta$ - $\text{Li}_2\text{HfTeO}_6$  are not divided. This is in line with the systematic absence law of XRD:  $\beta$ - $\text{Li}_2\text{HfTeO}_6$  with higher symmetry shows less diffraction peaks than  $\alpha$ - $\text{Li}_2\text{HfTeO}_6$  due to the extinction rule.



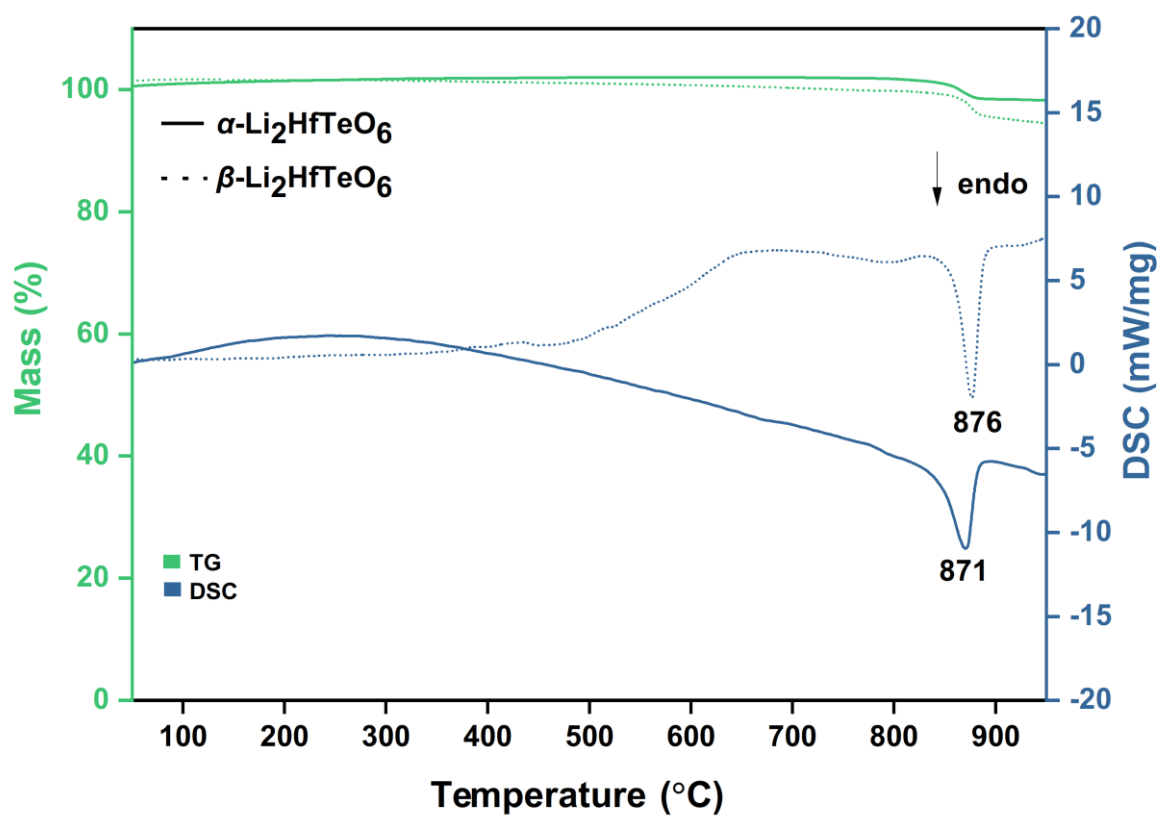
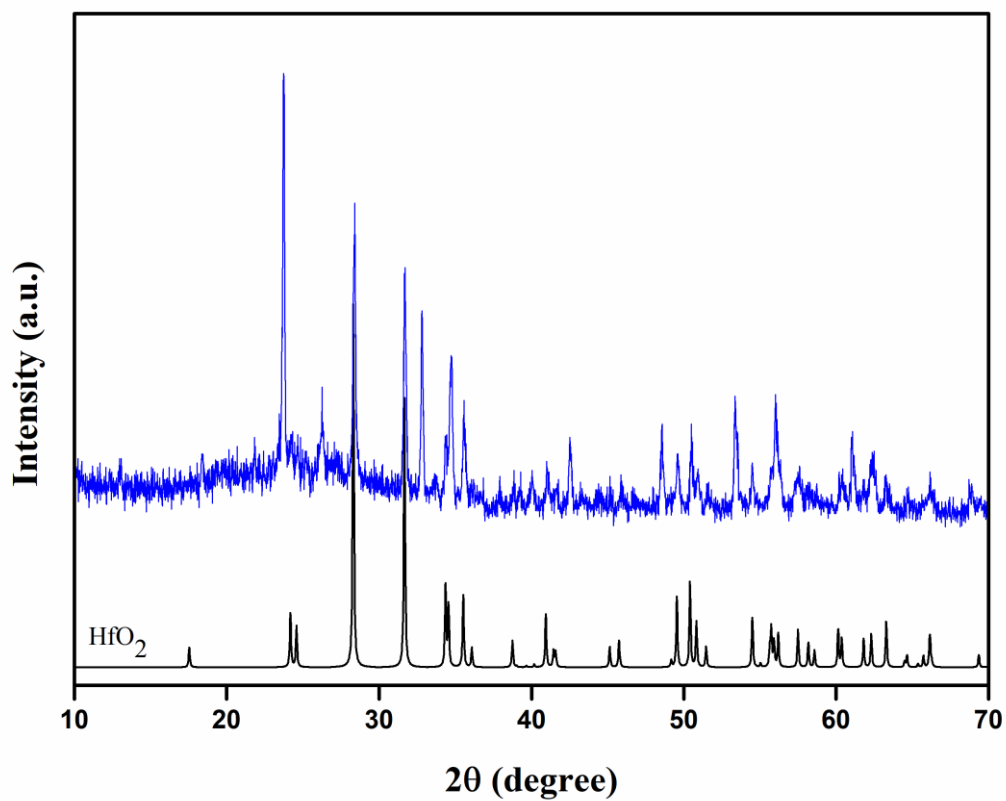
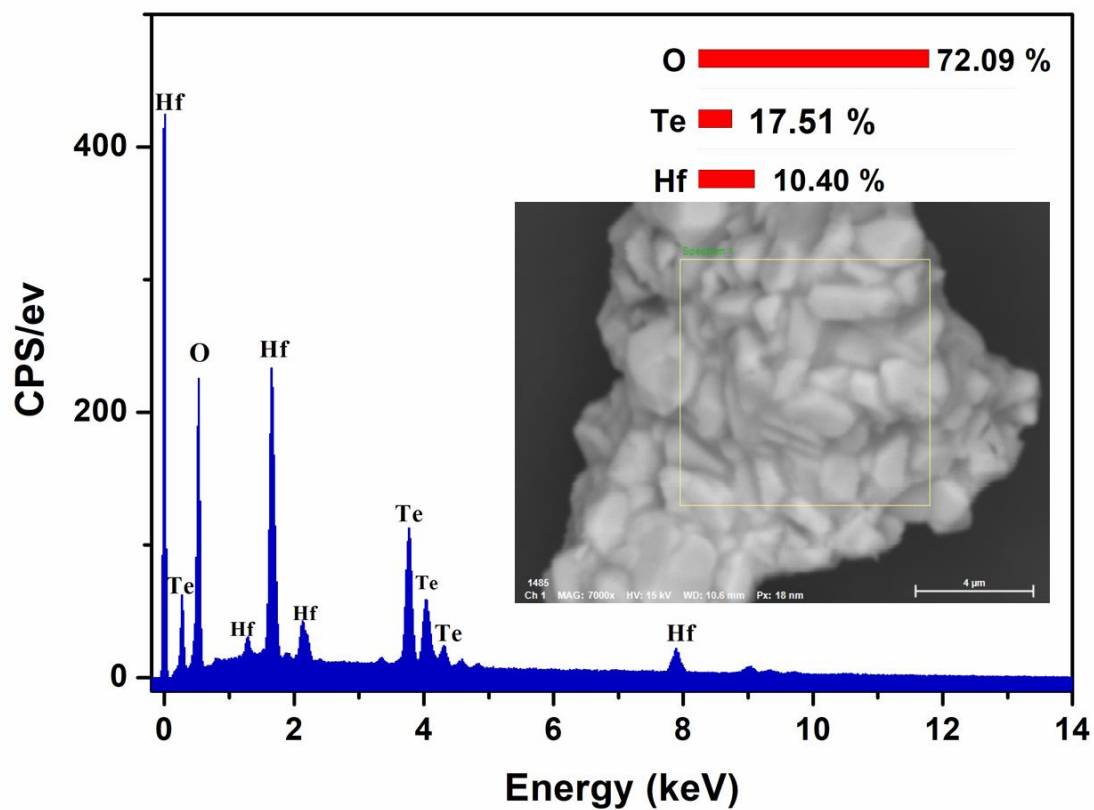


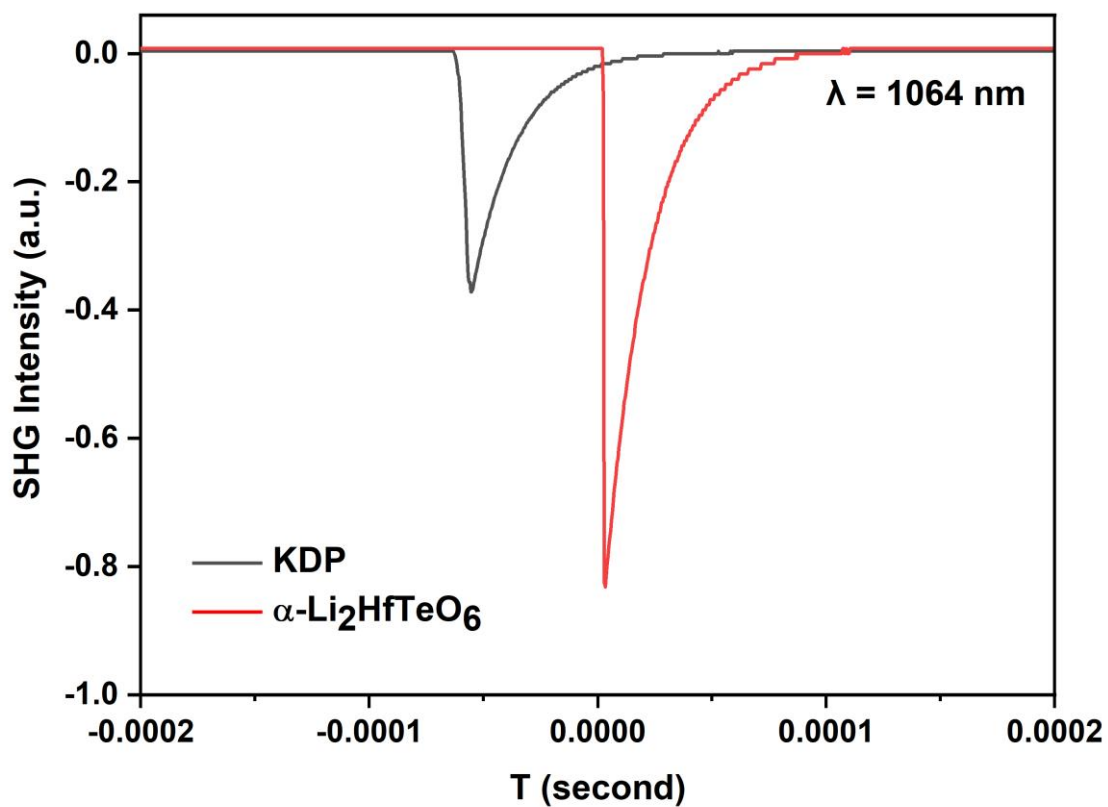
Figure S2. TG and DSC curve of  $\text{Li}_2\text{HfTeO}_6$  polymorphs.



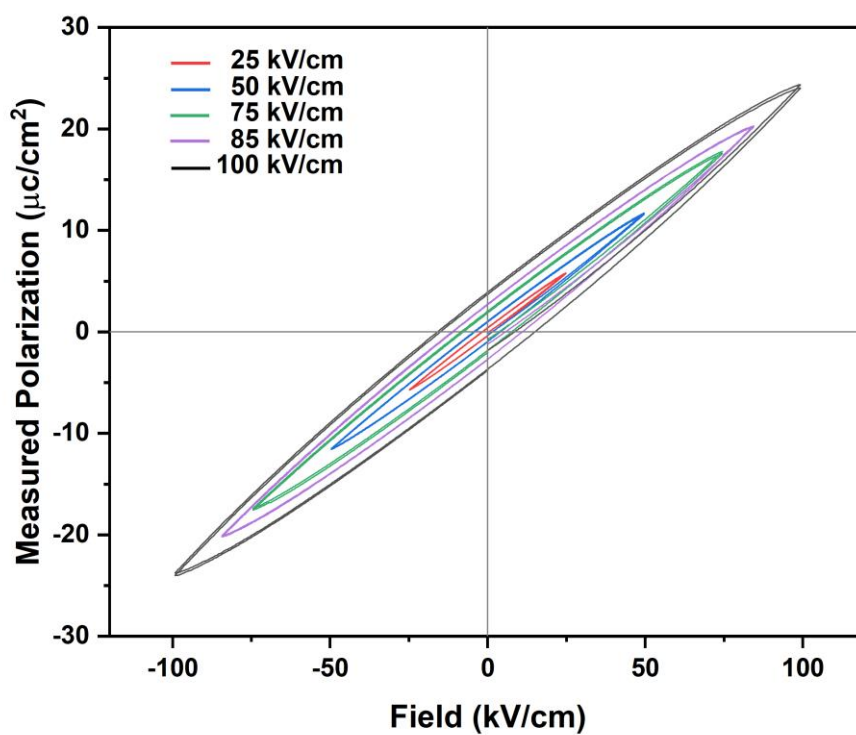
**Figure S3.** XRD pattern of  $\text{Li}_2\text{HfTeO}_6$  after heating to 850 °C. The main decomposition product is  $\text{HfO}_2$ .



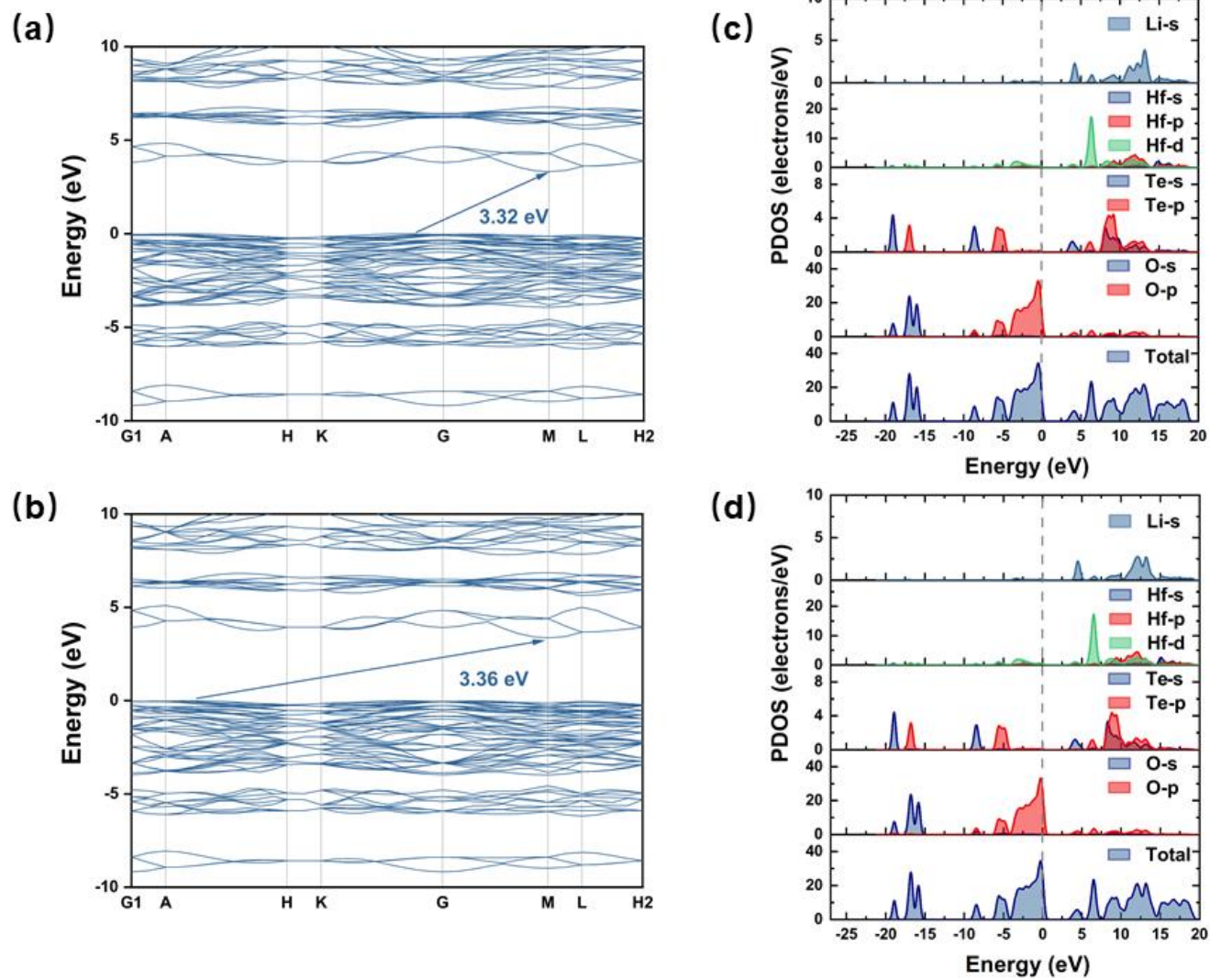
**Figure S4.** SEM images and EDX microanalysis for  $\beta$ -Li<sub>2</sub>HfTeO<sub>6</sub>. Noted that the EDX only provides semi-quantitative results, and the Li elements are ignored.



**Figure S5.** Oscilloscope traces showing the SHG intensities for the powder KDP and  $\alpha$ -Li<sub>2</sub>HfTeO<sub>6</sub>.



**Figure S6.** Polarization vs. voltage plot for  $\alpha\text{-Li}_2\text{HfTeO}_6$ . Although the curves reveal hysteresis loops, the observed loops do not represent ferroelectric behavior, but dielectric loss.<sup>13, 14</sup>



**Figure S7.** Bandgap of (a)  $\alpha\text{-Li}_2\text{HfTeO}_6$  and (b)  $\beta\text{-Li}_2\text{HfTeO}_6$ . The PDOS of (c)  $\alpha\text{-Li}_2\text{HfTeO}_6$  and (d)  $\beta\text{-Li}_2\text{HfTeO}_6$ .

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