Supplementary Information for

Probing Longitudinal Carrier Transport in Perovskite Films via Modified Transient Reflection Spectroscopy

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Preparation of MAPbI₃ perovskite thin films.

The perovskite thin films were prepared from modified nonstoichiometric precursor. For large-grain films, methylammonium iodide (MAI), lead iodide (PbI₂), and methylammonium chloride (MACl) (MAI: PbI₂: MACl =1:1:0.2) were dissolved in N-Methyl-2-pyrrolidone (NMP) and gamma- Butyrolactone (GBL) (NMP: GBL=7: 3, volume ratio) to form 55 wt% and 60 wt% precursor solutions, respectively. In order to prepare MAPbI₃ films with different thicknesses, 55 wt% MAPbI₃ precursor solutions was span at 4000 rpm, 3000 rpm, 2000 rpm for 25 s, and 60 wt% MAPbI₃ precursor solutions was span at 3000 rpm, 2000 rpm for 25 s respectively. After spincoating, these films were promptly transferred into abundant diethyl ether (DEE) solution for 90 s to finish the crystallization process and further annealed at 150 °C for 15 min with a Petri dish covered. The thickness of spun films were 550, 753, 1100, 1200 and 1540 nm, respectively. For small-grain films, 60 wt% precursor (MAI: PbI₂=1: 1) and precursor (MAI: PbI₂: MACl =1:1:0.1) solutions were used. All films were obtained by spin-coating their precursor solutions on FTO glasses followed by the same ether bathing and thermal annealing in air.

Transient absorption (TA) measurement.

The femtosecond TA spectrometer is based on a Yb:KGW laser system (1030 nm, 100 kHz; Light Conversion Ltd.), nonlinear frequency mixing techniques and the Femto-TA100 spectrometer (Time-Tech Spectra LLC). Briefly, the fundamental output at 1030 nm was split in two parts. One part was introduced to a noncollinear optical parametric amplifier to generate a certain wavelength for pump beam. The other was focused onto a YAG crystal to generate continuum white light (550–950 nm) as probe beam. The probe beam was focused with a parabolic reflector onto the sample and the transmitted probe beam was then collimated and focused into a fiber-coupled spectrometer with CMOS sensors. The intensity of the pump pulse used in the experiment was controlled by a variable neutral-density filter wheel. The delay between the pump and probe pulses was controlled by a motorized delay stage. The pump pulses were chopped by a synchronized chopper at 5 KHz and the absorbance change was calculated with two adjacent probe pulses (pump-blocked and pump-unblocked).

Normal and back-excitation transient reflection (TR) measurement.

The TR measurements were performed by using the same pump–probe transient spectrometer with a modified optical path, where the detection signal is the reflected probe beam. Briefly, the fundamental output from a Yb:KGW laser (1030 nm, 100 kHz; Light Conversion Ltd.) was split in two parts. One part was introduced to a noncollinear optical parametric amplifier to generate a certain wavelength for pump beam. The other was focused onto a YAG crystal to generate continuum white light (550–950 nm) as probe beam. The probe beam was focused with a parabolic reflector onto the sample and the reflected probe beam was then collimated and focused into a fiber-coupled spectrometer. The incident angle for the pump beam is around 0° and that for the probe beam is around 45°. For normal TR measurement, the pump and

probe beams impinge on the same side of the sample and were spatially overlapped on the sample surface. For back-excitation TR measurement, the pump and probe beams respectively impinge on both side of the sample. All experiments were performed at room temperature.



Fig. S1 Top-view (upper panels) and cross-sectional (lower panels) SEM images of five large-grain MAPbI₃ films with different thicknesses of (a) 1200 nm, (b) 550 nm, (c) 753nm, (d) 1100 nm and (e) 1540 nm. Scale bar: 2 μ m for top-view images, 1 μ m for cross-sectional images.



Fig. S2 Top-view (upper panels) and cross-sectional (lower panels) SEM images of four small-grain MAPbI₃ films with different thicknesses of (a) 980 nm, (b) 1060 nm, (c) 1150nm, (d) 850 nm. Scale bar: 2 μ m for top-view images, 1 μ m for cross-sectional images.



Fig. S3 Lateral grain size statistical distribution of all as-prepared MAPbI₃ films determined from their top-view SEM images as shown in Fig. S1 and S2.



Fig. S4 (a) The UV–Vis absorption and PL spectra of a typical MAPbI₃ film (film 1) as shown in Fig. S1a. (b) The measured linear absorption spectra (red squares) of the film and its corresponding simulations (black line) based on Elliott's model. The blue and green dashed lines represent the excitonic and continuum contributions to the total absorption. For photon energies greater than ~1.8 eV, the absorption no longer follows Elliott's model since the contribution from higher energy bands is ignored.



Fig. S5 (a) TA spectra measured in MAPbI₃ film at a delay time of 5 ps. (b) TA kinetics probed at 760 nm in the MAPbI₃ film. The red line is a bi-exponential fit of TA kinetics, yielding the average bulk carrier lifetime of $\tau = 6.1$ ns.



Fig. S6 The TR spectra recorded at 2100 ps delay under back-excitation at 635 nm in the same perovskite film (film 1) as used in Fig. 2. The Hilbert transform of TA spectrum (blue trace) and the interference signal (green trace) deconvolved from total TR spectrum are also shown. The spectral features in the grey shaded region result from the interference of the reflected probe beam from both the front and back surface of thin film.



Fig. S7 The comparison of the normalized TR kinetics probed at 720 nm in typical perovskite film (film 1) as shown in Fig. S1a under normal- (red trace) and back-excitation (blue trace) TR modes respectively, as well as its TA kinetics (black trace) probed at 760 nm of the perovskite film.



Fig. S8 The normal TR kinetics probed at 720 nm in typical perovskite film (film 1) as shown in Fig. S1a with different excitation fluence under indicated pump wavelength at (a) 400nm, (b) 540nm, (c) 635nm, respectively. These data indicate that the higher-order recombination can be avoided and thus these TR kinetics are independent of the excitation intensities when the excitation fluence $< 7 \mu J/cm^2$.



Fig. S9 The back-exc. TR kinetics probed at 720 nm in perovskite film with 1200nm thickness under 400 nm excitation with different fluences. These data indicate that the TR kinetics is independent of the excitation intensities when the excitation fluence $< 7 \mu J/cm^2$ and thus the effect of high-order recombination in this case is negligible.



Fig. S10 The simulated initial carrier distributions along the film thickness direction at different excitation wavelengths.



Fig. S11 The measured absorption coefficient of $MAPbI_3$ perovskite film as a function of wavelength.



Fig. S12 Normal TR kinetics of the typical perovskite film as shown in Fig. S1a. The kinetics are extracted from the normal TR spectra collected from the front surface (blue trace) and the back surface (red trace). The back surface is the interface between the perovskite film and FTO glass. The consistency of the two kinetic traces suggests that the surface recombination velocity (S) for the front and back surfaces are similar.



Fig. S13 The normal TR kinetics probed at 720 nm in the same film with larger grains as in Fig. 3a (film 1) under excitation at indicated wavelengths. The solid lines are their global fits according to the diffusion model discussed in the main text.



Fig. S14 Simulated longitudinal carrier distribution at indicated delay times in 1200-nm-thickness perovskite film with D=1.5 cm²s⁻¹, S=2000 cm/s and $k_1=10^8$ s⁻¹. The excitation wavelengths are indicated in the figures. The top row (a, d) shows the

carrier distribution in the whole film. The middle row (b, e) shows the carrier distribution in the effective detection depth (~30 nm) of the front surface. The bottom row (c, f) shows the carrier distribution in the effective detection depth of the back surface. It is evident that the carrier density at the front surface exhibits an uneven distribution during initial delays after excitation, while that at the back surface always shows a uniform distribution during any delays.



Fig. S15 The simulated TR kinetics in 1200-nm-thickness perovskite film (with $D=1.5 \text{ cm}^2\text{s}^{-1}$, S=2000 cm/s and $k_1=10^8 \text{ s}^{-1}$) under 400 nm excitation by using different detection depth in (a) normal- and (b) back-excitation TR mode, respectively. The remarkable difference of simulated TR kinetics among different detection depths in normal TR mode suggests that the detection depth necessarily influences their fitting results, while the consistency of TR kinetics among different detection depths in back-excitation TR mode suggests that the detection depth should not affect their fitting results.



Fig. S16 XRD patterns of five large-grain MAPbI₃ thin films (film 1-5) and a smallgrain film (film 8), as shown in Fig. S1 and S2c. For large-grain films, their XRD patterns only exhibit two very strong diffraction peaks, suggesting that each grain within the film is highly oriented and high crystallinity. It is quite different from small-grain film where other diffraction peaks from different crystallographic planes also appear because the small grains prefer random orientation.



Fig. S17 The back-excitation TR kinetics probed at 720 nm in various MAPbI₃ perovskite films (film 2-9) with different thicknesses and grains under excitation at the indicated pump wavelength. The solid lines are their global fits according to the diffusion model discussed in the main text and the detailed fitting results are summarized in Table S1. Note that the bulk carrier lifetimes used during the fitting for the films with large grains (i.e., film 1-5) are extracted from their TA kinetics due to their very short lifetimes, while those for the films with small grains (i.e., film 6-9) are extracted from their relatively long lifetimes.



Fig. S18 Comparison of PL kinetics collected at 770 nm in various perovskite films with (a) large grains (film 1–5) and (b) small grains (film 6-9).

Table	S1 .	Fitting	parameters	for	back-excita	ation TI	R kine	etics	of	different	perov	skite
films s	show	n in Fig	g. 3a and S1	7. (I	D and S are	set as fr	ee fitt	ing	para	ameters.)		

film No.	Thickness	Grain size	R	$D(cm^2s^{-1})$	S (cm s ⁻¹)
	(nm)	(um)			
film 1	1200	1.81	0.66	1.52	5807.5
film 2	550	1.25	0.44	1.75	3035.5
film 3	753	1.67	0.45	1.72	3669.0
film 4	1100	2.13	0.51	1.45	4913.0
film 5	1540	2.58	0.59	1.70	1611.2
film 6	980	0.96	1.02	1.19	0*
film 7	1060	0.76	1.39	0.77	1213*
film 8	1150	0.46	2.5	0.49	0*
film 9	850	0.36	2.36	0.48	0*

* Note: The S values obtained for four films (film 6-9) with small grains are of insignificance due to the absence of the evident decay process in their TR kinetics (see Fig. S17).