

## Supporting Information

Ruoyu Wang<sup>1,2</sup>, Zhe Guo<sup>1,2</sup>, Qiang Zhang<sup>1,2</sup>, Jianfeng Cai<sup>1,2</sup>, Guoqiang Liu<sup>1,2,\*</sup>,  
Xiaojian Tan<sup>1,2,\*</sup>, Jun Jiang<sup>1,2,\*</sup>

<sup>1</sup>Ningbo Institute of Materials Technology and Engineering, Chinese Academy of  
Sciences, Ningbo 315201, China

<sup>2</sup> University of Chinese Academy of Sciences, Beijing 100049, China

E-mail: [liugq@nimte.ac.cn](mailto:liugq@nimte.ac.cn); [tanxiaojian@nimte.ac.cn](mailto:tanxiaojian@nimte.ac.cn); [jjun@nimte.ac.cn](mailto:jjun@nimte.ac.cn)

## Figures

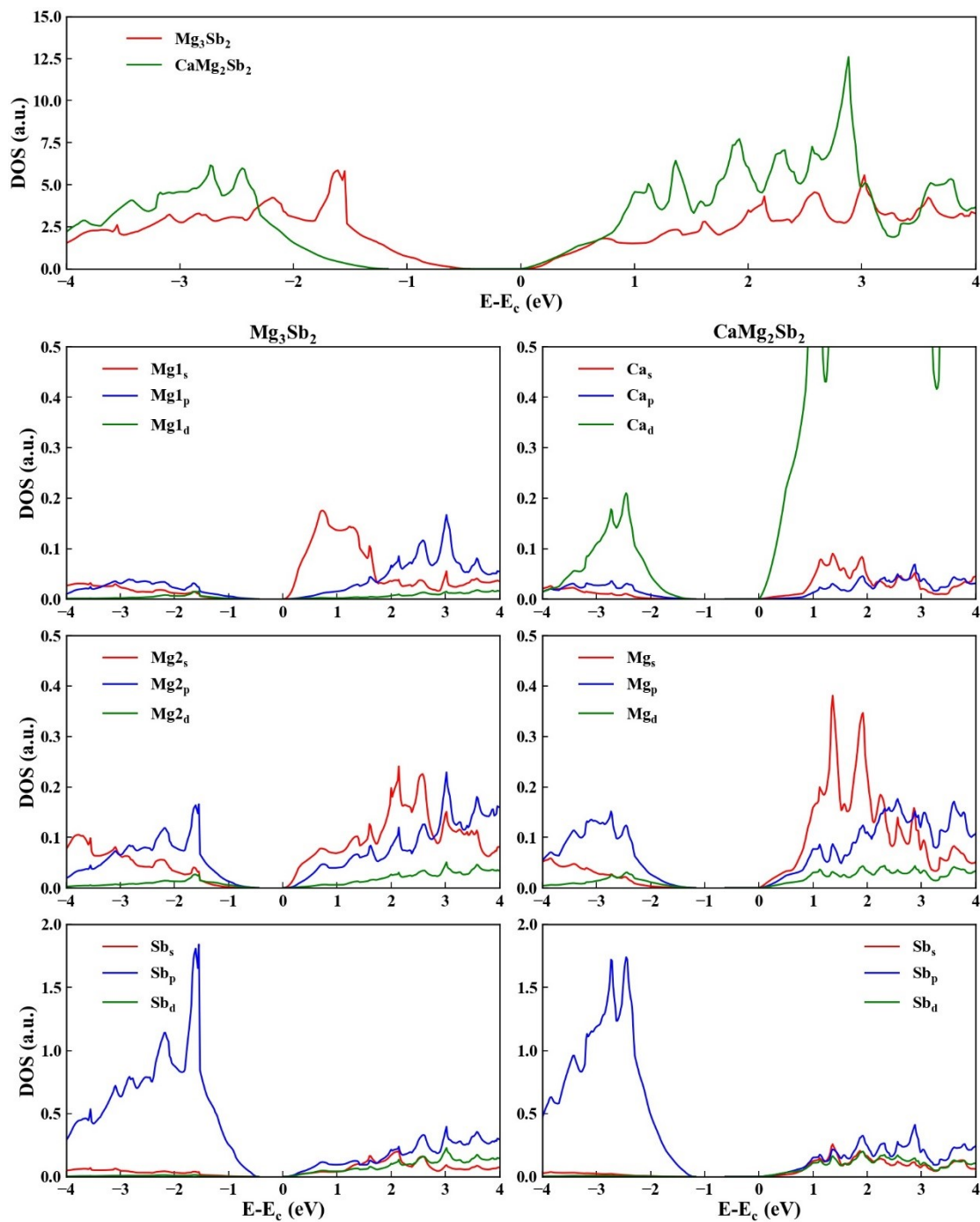


Figure S1. The total electronic density of states of  $\text{CaMg}_2\text{Sb}_2$  and  $\text{Mg}_3\text{Sb}_2$  (Upper). The atomic partial density of states of  $\text{Mg}_3\text{Sb}_2$  (left column), atomic partial density of states of  $\text{CaMg}_2\text{Sb}_2$  (right column).

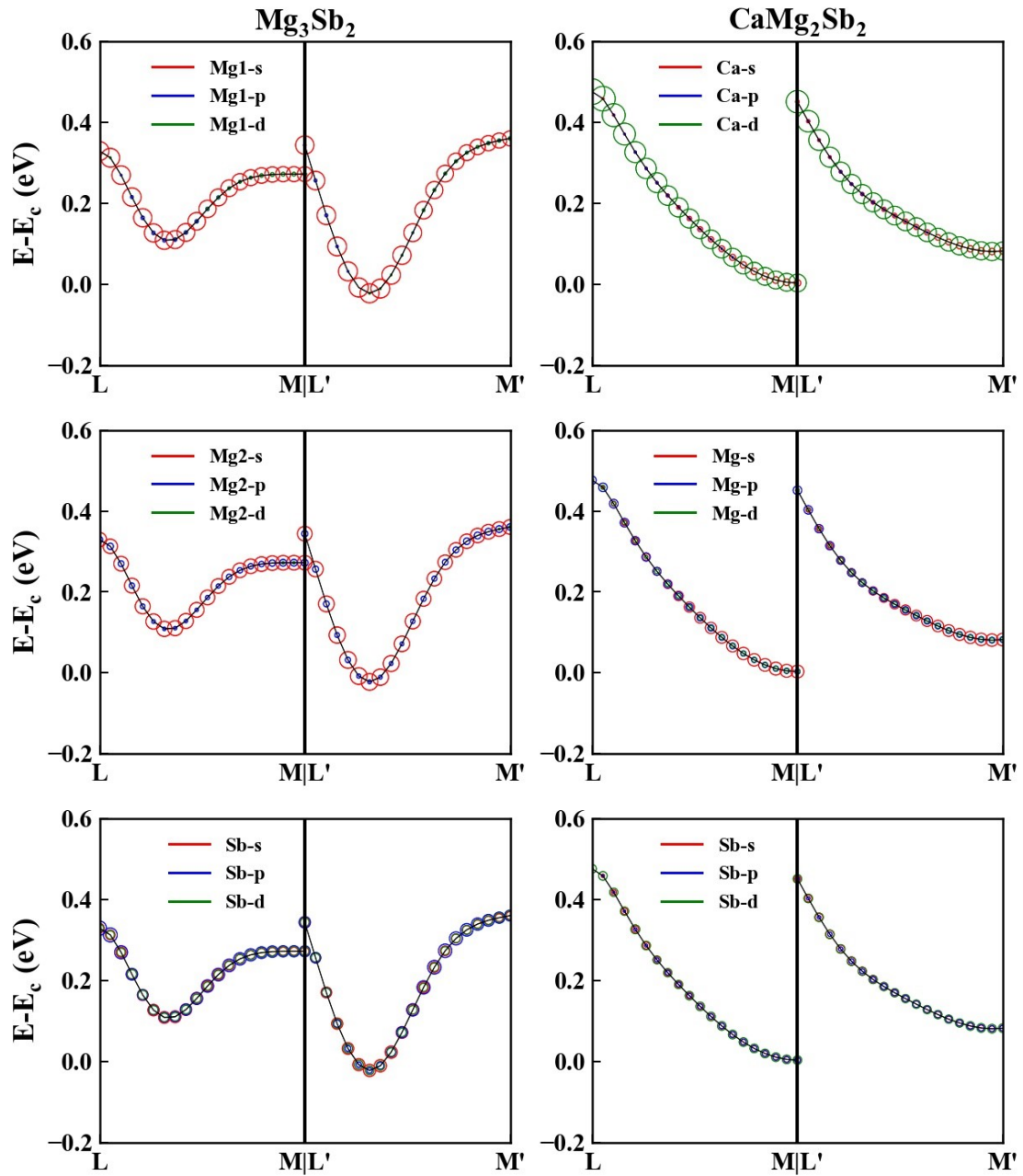


Figure S2. The projected electronic band structure for the conduction band edge of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$ .

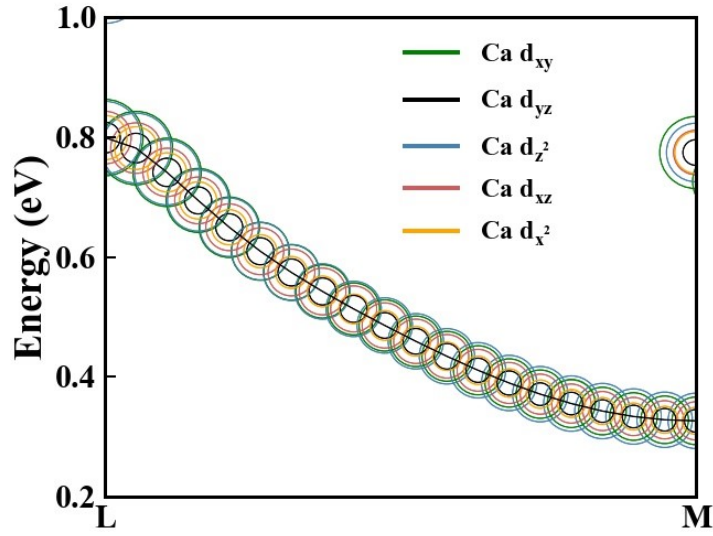


Figure S3. The projected band structure for the five Ca-*d* orbitals at the conduction band minimum of CaMg<sub>2</sub>Sb<sub>2</sub> at the 'M' point. The three most significant components at M point are Ca-*d<sub>z</sub><sup>2</sup>*, Ca-*d<sub>xy</sub>* and Ca-*d<sub>xz</sub>*, respectively.

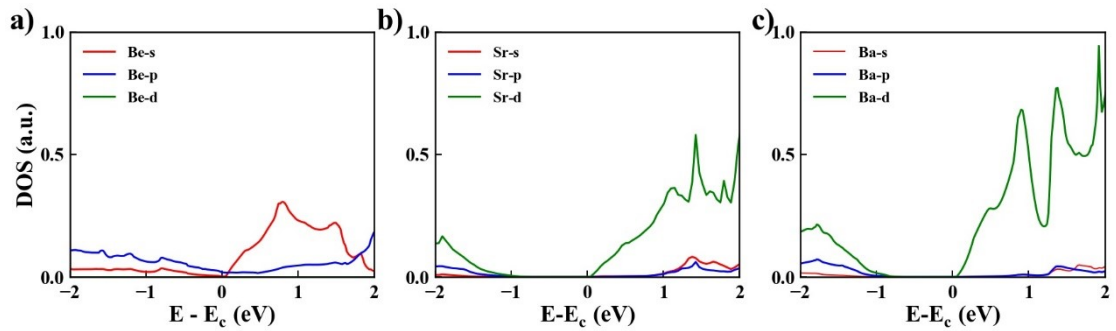


Figure S4. The partial density of states for the 'A'-site atom in BeMg<sub>2</sub>Sb<sub>2</sub>, SrMg<sub>2</sub>Sb<sub>2</sub> and BaMg<sub>2</sub>Sb<sub>2</sub>, respectively.

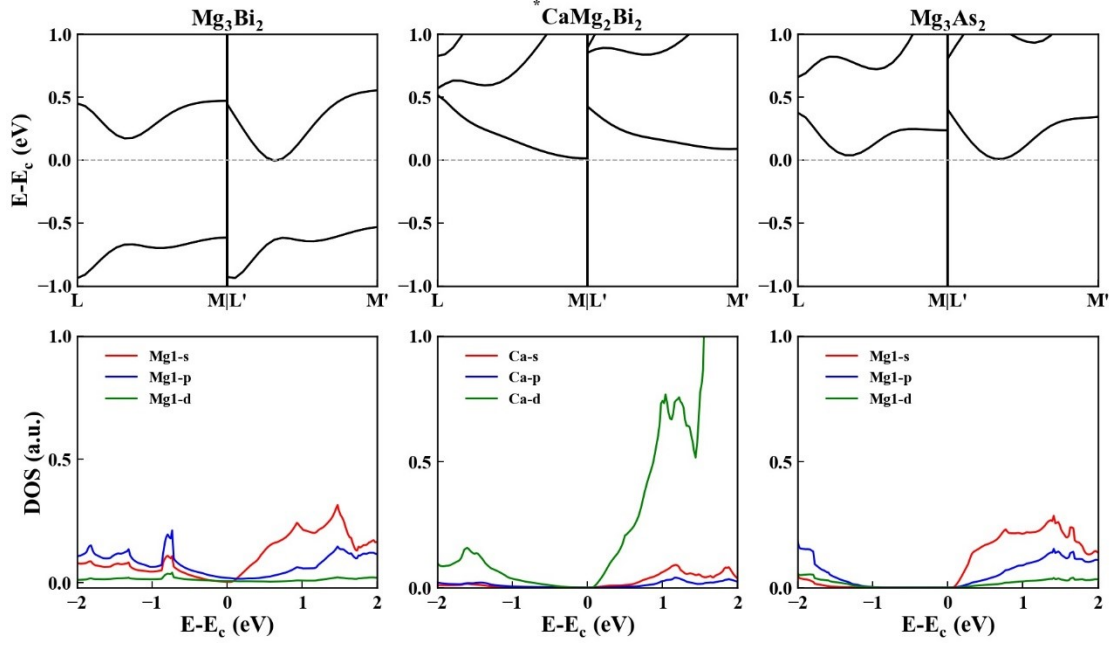


Figure S5. Band dispersion and partial density of states at the conduction band edge of  $\text{Mg}_3\text{Bi}_2$ ,  $\text{CaMg}_2\text{Bi}_2$  and  $\text{Mg}_3\text{As}_2$ . (Note: the CBM of  $\text{CaMg}_2\text{Bi}_2$  is in fact located at the Gamma point. Here the CBM is set to the M point for consistency)

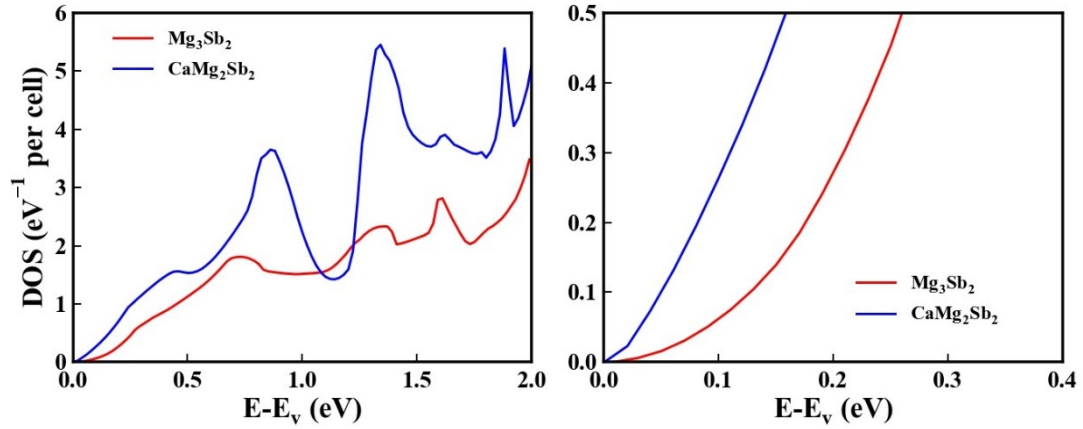


Figure S6. Comparison of total density of states of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$  at conduction band edge.

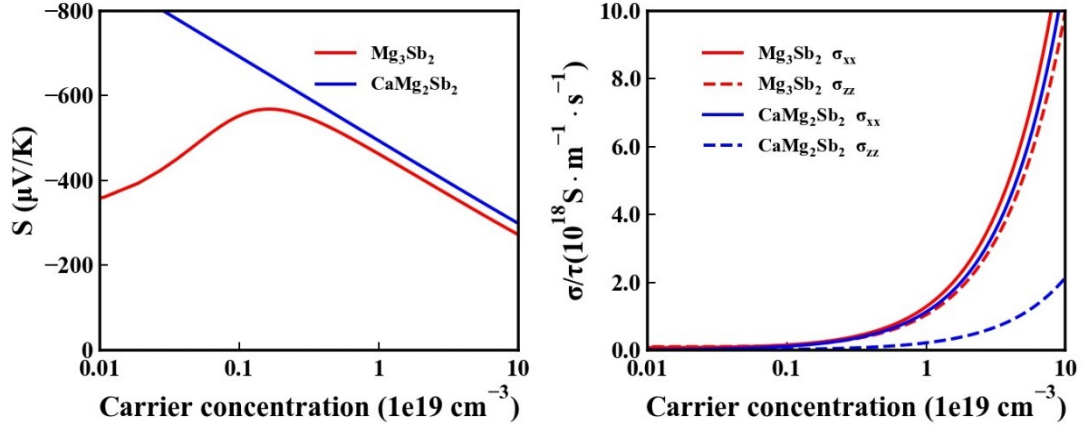


Figure S7. Seebeck coefficient and the electrical conductivity/relaxation time ratio of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$  at 800 K.

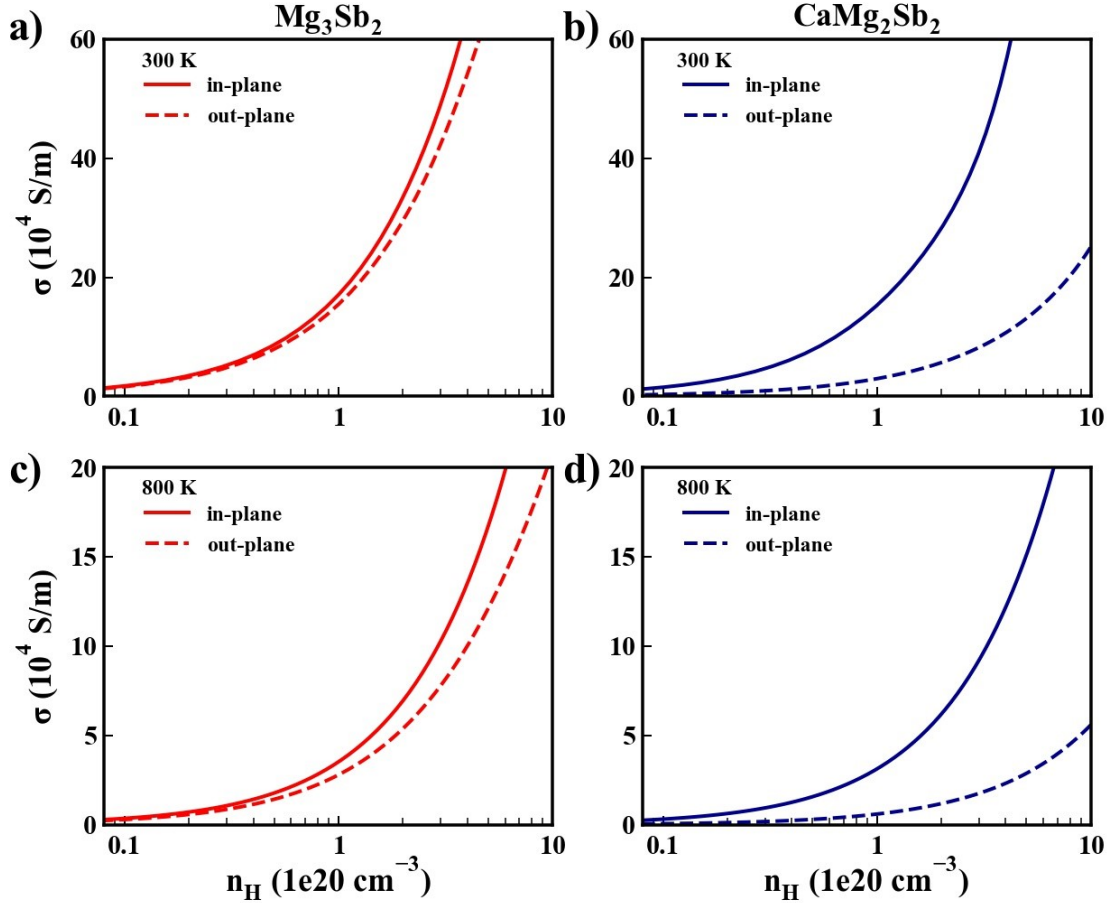


Figure S8. Estimated electrical conductivity of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$ . The in-plane and out-plane  $\sigma$  of  $\text{Mg}_3\text{Sb}_2$  at (a) 300 K and (c) 800 K. The in-plane and out-plane  $\sigma$  of  $\text{CaMg}_2\text{Sb}_2$  at (b) 300 K and (d) 800 K.

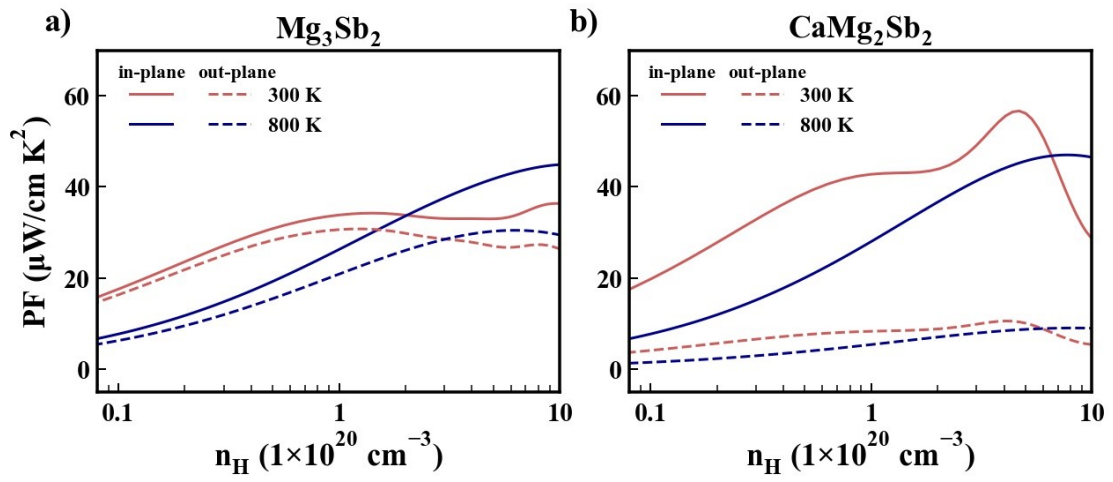


Figure S9. The estimated PF of (a)  $\text{Mg}_3\text{Sb}_2$  and (b)  $\text{CaMg}_2\text{Sb}_2$  at 300 and 800 K.

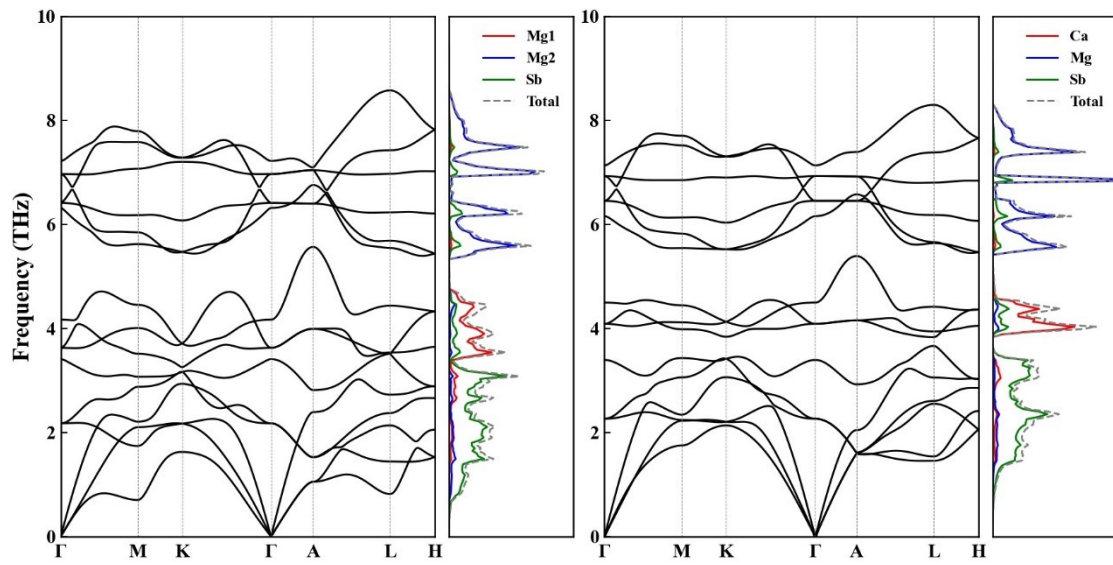


Figure S10. The phonon dispersion and projected density of states of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$ , respectively.



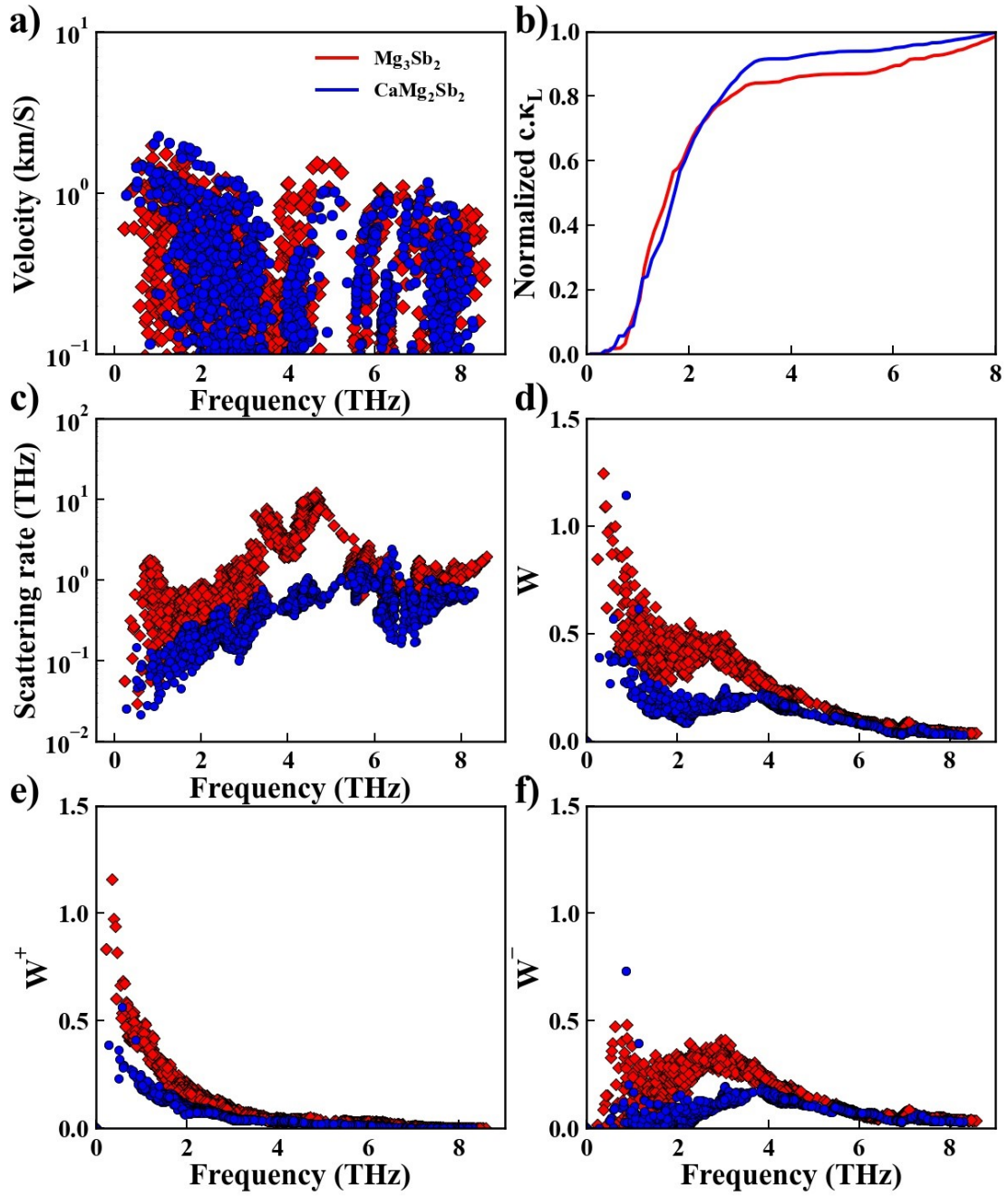


Figure S11. a) The sound velocity of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$ . b) The cumulative  $\kappa_L$  of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$  normalized to their respective  $\kappa_L$ . c) Anharmonic scattering rate of  $\text{Mg}_3\text{Sb}_2$  and  $\text{CaMg}_2\text{Sb}_2$ . d) The weighted scattering phase space of three-phonon processes. e) The weighted scattering phase space of three-phonon annihilation process. And f) the weighted scattering phase space of three-phonon absorption process.



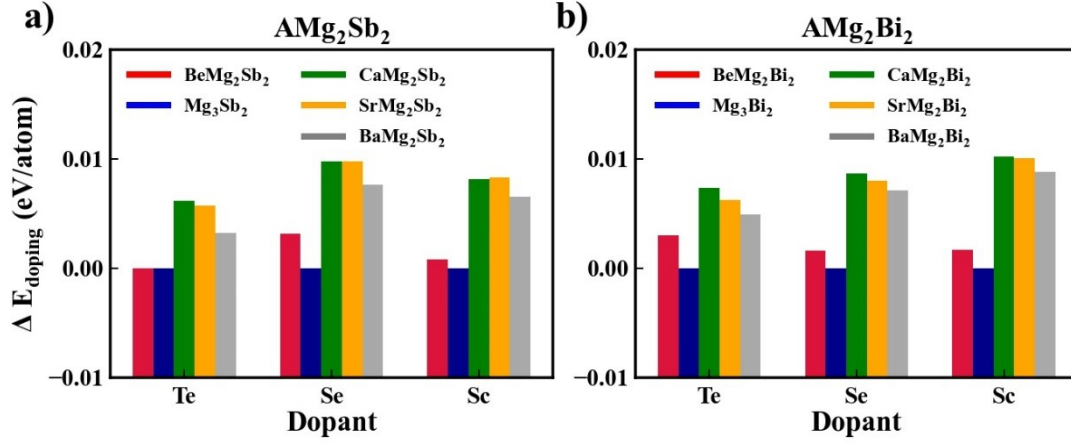


Figure. S12 Doping energy  $\Delta E_{doping}$  of  $\text{BeMg}_2\text{X}_2$  ( $\text{X}=\text{Sb}, \text{Bi}$ ) compared with their  $\text{AMg}_2\text{X}_2$  ( $\text{A}=\text{Mg}, \text{Ca}, \text{Sr}$  and  $\text{Ba}$ ) counterparts. The  $\Delta E_{doping}$  of  $\text{BeMg}_2\text{X}_2$  significantly lower than  $\text{CaMg}_2\text{X}_2$ ,  $\text{SrMg}_2\text{X}_2$  and  $\text{BaMg}_2\text{X}_2$ , due to the absent  $p$ - $d$  bonding.

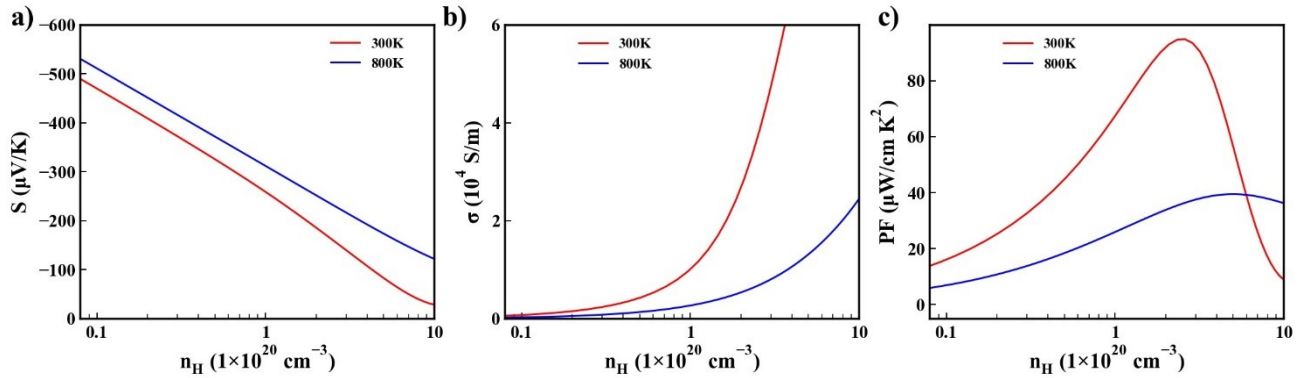


Figure S13. The estimated *in-plane* thermoelectric transport of  $\text{BaMg}_2\text{Sb}_2$ . a) Seebeck coefficient b) electrical conductivity and c) power factor.

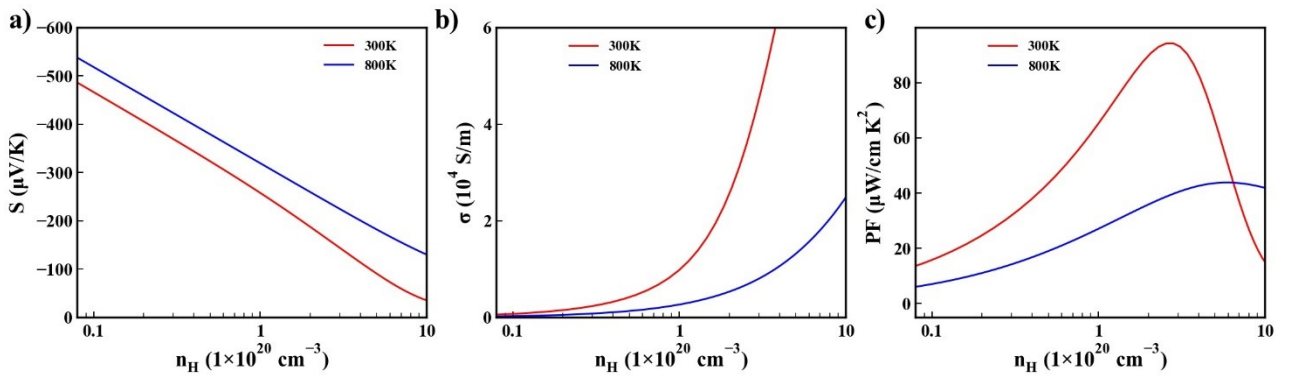


Figure S14. The estimated *in-plane* thermoelectric transport of  $\text{BaMg}_2\text{SbBi}$  alloy. a) Seebeck coefficient b) electrical conductivity and c) power factor.

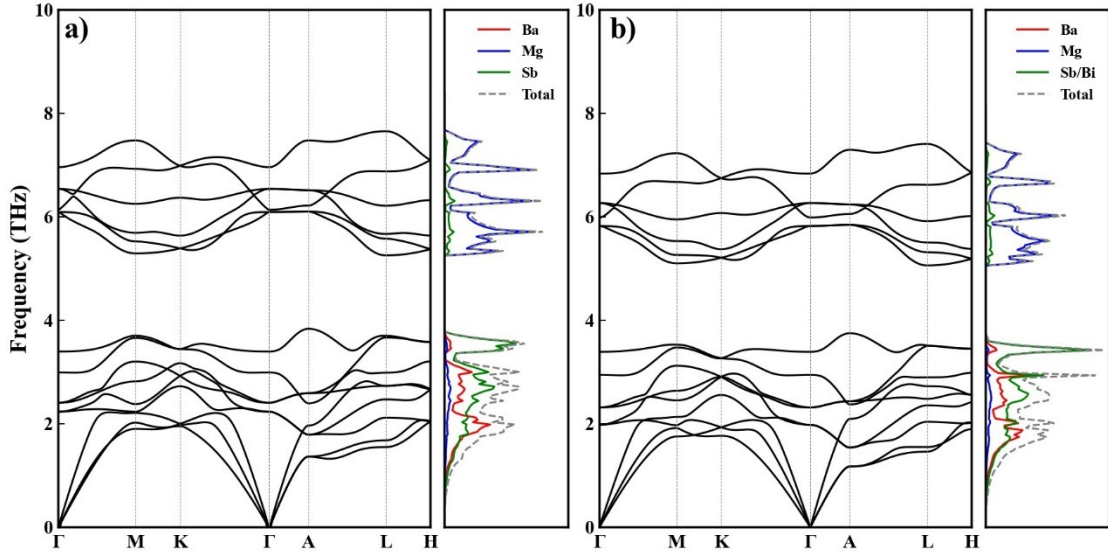


Figure S15. The phonon dispersion and projected density of states of  $\text{BaMg}_2\text{Sb}_2$  and  $\text{BaMg}_2\text{SbBi}$ , respectively. The phonon dispersion of  $\text{BaMg}_2\text{SbBi}$  is calculated using the “virtual crystal” technique.

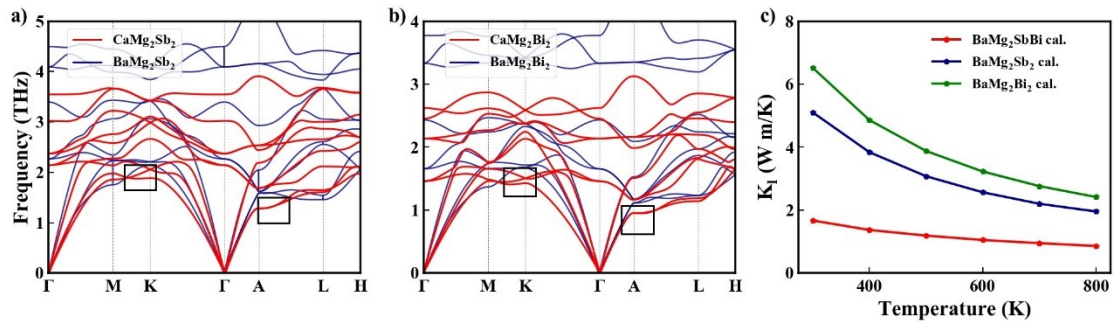


Figure S16. The phonon dispersion and calculated  $\kappa_L$  of  $\text{BaMg}_2\text{X}_2$  ( $X=\text{Sb}, \text{Bi}$ ). (a) The phonon dispersion spectrum of  $\text{CaMg}_2\text{Sb}_2$  and  $\text{BaMg}_2\text{Sb}_2$ . Note that moderate band softening can be observed for  $\text{BaMg}_2\text{X}_2$ . (b) The phonon dispersion spectrum of  $\text{CaMg}_2\text{Bi}_2$  and  $\text{BaMg}_2\text{Bi}_2$ . (c) The calculated of  $\text{BaMg}_2\text{Sb}_2$ ,  $\text{BaMg}_2\text{Bi}_2$  and  $\text{BaMg}_2\text{SbBi}$  solid solution.

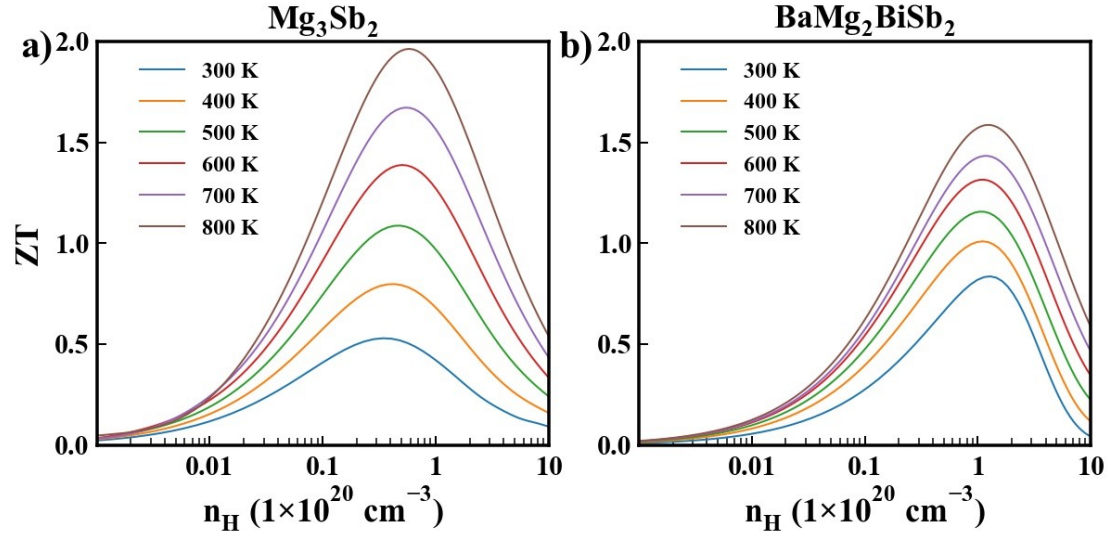


Figure S17. The estimated *in-plane* ZT values of  $\text{Mg}_3\text{Sb}_2$  and  $\text{BaMg}_2\text{Sb}_2$ .

### Tables:

**Table S1.** The 2<sup>nd</sup> order IFC between Sb and up to the 3<sup>rd</sup> nearest neighbor in various  $\text{AMg}_2\text{Sb}_2$  compounds. The unit of 2<sup>nd</sup> order IFC is  $\text{eV}\text{\AA}^{-2}$

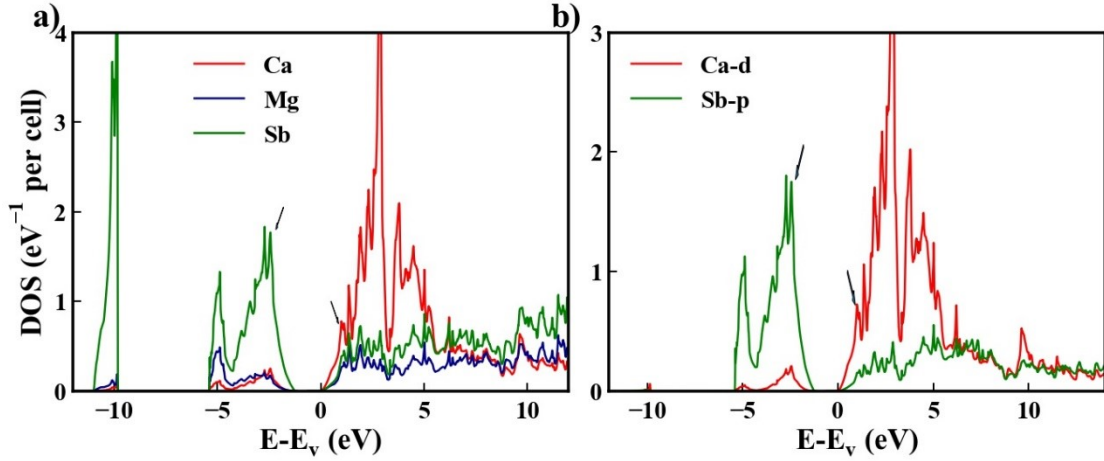
	P1(Sb-Mg)	P2 (Sb-Mg)	P3 (Sb-A)
$\text{MgMg}_2\text{Sb}_2$	2.32	2.18	0.59
$\text{CaMg}_2\text{Sb}_2$	2.60	2.16	1.13
$\text{BaMg}_2\text{Sb}_2$	2.45	2.21	1.11

**Table S2.** The electronic relaxation time extracted from the experimental electrical transport properties of  $\text{Mg}_3\text{Sb}_2$  single crystal.<sup>1</sup>

Temperature (K)	300	400	500	600	700	800
$\tau$ ( $10^{-14}$ s)	1.23	0.80	0.57	0.44	0.34	0.28

## Analysis of the $p$ - $d$ bonding in $\text{CaMg}_2\text{Sb}_2$

Figure S18 presents the PDOS of Ca- $d$  and Sb- $p$  orbitals in  $\text{CaMg}_2\text{Sb}_2$ . Clearly, the major peak of Ca- $d$  orbital in the valence band is around -2.5 eV, where Sb- $p$  orbital exhibits a higher peak with similar shape. The similar DOS peaks of Sb- $p$  and Ca- $d$  orbitals could also be observed in the conduction bands from 0 to 5 eV. The strong hybridization of the PDOS indicates a covalent-like bonding between Sb- $p$  and Ca- $d$  orbitals. Therefore, the DOS peak around -2.5 eV can be identified as a bonding state, and conduction band edge can be identified as an anti-bonding state.



**Fig. S18** Partial density of states of  $\text{CaMg}_2\text{Sb}_2$ . (a) Atomic PDOS of Ca, Mg and Sb atoms. (b) PDOS of Ca- $d$  and Sb- $p$  orbitals

## Calculation details for the energy change upon doping.

The formation energy ( $E_{\text{formation}}$ ) of a given compound can be written as:

$$E_{\text{formation}} = E_{\text{compound}} - \sum n_i E_{\text{atom},i} \quad (1)$$

where  $E_{\text{compound}}$  is the total energy of compound, and  $n_i$  and  $E_{\text{atom},i}$  are the number and energy of  $i^{\text{th}}$  component atom in free space, respectively. The difference in  $E_{\text{formation}}$  between two compounds can be calculated as the follows:

$$\Delta E_{\text{formation}} = \Delta E_{\text{compound}} - \sum \Delta n_i E_{\text{atom},i} \quad (2)$$

When an atom in some pure substance is substituted for a dopant atom, the doping energy  $E_{\text{doping}}$  can be defined as following

$$E_{\text{doping}} = (E_{\text{doped}} - E_{\text{pure}}) - (E_{\text{atom,dopant}} - E_{\text{atom,a}}) \quad (3)$$

The  $E_{\text{doping}}$  may indicate the level of difficulty in doping. A larger  $E_{\text{doping}}$  suggest the

doping be more difficult and *vice versa*. Though  $E_{atom}$  is not directly accessible within the framework of VASP package which has an arbitrary zero-point energy, we can still compare the doping of two different compound, provided that both the dopant element and the doping site are the same in the two compounds. Therefore, the relative doping energy of two compound, e.g., the Te doping at the Sb site for  $Mg_3Sb_2$  and  $CaMg_2Sb_2$ ,  $\Delta E_{doping}$ , can be written as

$$\Delta E_{doping} = (E_{Te,Sb,CaMg_2Sb_2} - E_{pure,CaMg_2Sb_2}) - (E_{Te,Sb,Mg_3Sb_2} - E_{pure,Mg_3Sb_2}) \quad (4)$$

All terms in (4) are calculable with the VASP package.

For the specific calculation of  $\Delta E_{doping}$  carried out in this work, the Te or Se dopant replace X site atom (Sb or Bi), whereas the Sc dopant substitutes the tetrahedrally coordinated Mg in  $AMg_2X_2$ .

## Calculation details for the electrical conductivity $\sigma$ and electrical thermal conductivity $\kappa_e$

The electronic transport properties of all phases are calculated using the Boltzmann transport theory with the rigid band approximation. The electrical conductivity is calculated by the equation  $\sigma = (\sigma/\tau)\tau$ , where  $\tau$  denotes the relaxation time. The  $\sigma/\tau$  can be calculated by solving Boltzmann transport equation (BTE), whereas  $\tau$  is extracted from experimental measurement. Since there is no available experimental data for the n-type materials except  $Mg_3(Sb, Bi)_2$ , the relaxation time ( $\tau$ ) is extrapolated from the experimental value of the n-type  $Mg_3Sb_2$  single crystal (Table S2) and applied to all other phases.<sup>1</sup> For the calculation of electronic thermal conductivity ( $\kappa_l$ ), the Widemann-Franz relation ( $\kappa_l = L\sigma T$ ) is used as an approximation, with the  $L$  chosen as the Sommerfeld value. Due to the *p-d* bonding, the n-type  $CaMg_2Sb_2$  and  $BaMg_2Sb_2$  exhibits anisotropic n-type electrical conductivity, and hence the anisotropic power factors, with the *in-plane* direction being more favored.

## References

- (1) Imasato, K.; Fu, C.; Pan, Y.; Wood, M.; Kuo, J. J.; Felser, C.; Snyder, G. J.

- Metallic N-Type  $\text{Mg}_3\text{Sb}_2$  Single Crystals Demonstrate the Absence of Ionized Impurity Scattering and Enhanced Thermoelectric Performance. *Adv. Mater.* **2020**, *32* (16), 1908218. <https://doi.org/10.1002/adma.201908218>.
- (2) Tamura, S. Isotope Scattering of Dispersive Phonons in Ge. *Phys. Rev. B* **1983**, *27* (2), 858–866. <https://doi.org/10.1103/PhysRevB.27.858>.
- (3) Pei, Y.; Gibbs, Z. M.; Gloskovskii, A.; Balke, B.; Zeier, W. G.; Snyder, G. J. Optimum Carrier Concentration in N-Type PbTe Thermoelectrics. *Adv. Energy Mater.* **2014**, *4* (13), 1400486. <https://doi.org/10.1002/aenm.201400486>.