

Electronic Supplementary Information for

Optical generation and electric control of pure spin photocurrent in ferroelectric Ruddlesden-Popper perovskite $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer[†]

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Table S1: Comparison between the calculated (Cal) and experimentally measured (Exp¹) crystallographic parameters for bulk (MA)₂Pb(SCN)₂I₂ in Pmn2₁ symmetry, d_{Pb} indicates the polar displacement of Pb cation relative to Pb(SCN)₂I₄ octahedral center.

FE bulk (MA) ₂ Pb(SCN) ₂ I ₂						
Exp: a = 18.580 Å, b = 6.267 Å, c = 6.466 Å, d_{Pb} = 0.202 Å						
Cal: a = 17.849 Å, b = 6.178 Å, c = 6.398 Å, d_{Pb} = 0.374 Å						
Atom	Exp			Cal		
	x	y	z	x	y	z
Pb (2a)	0.50000	0.39288	0.23378	0.50000	0.37280	0.22410
I (2a)	0.50000	0.89281	0.27687	0.50000	0.86758	0.29457
I (2a)	0.50000	0.36268	-0.25772	0.50000	0.33813	-0.26116
S (4b)	0.34205	0.32704	0.22918	0.33470	0.31644	0.22633
C (4b)	0.31055	0.56112	0.26832	0.30834	0.56446	0.26784
C (4b)	0.35707	-0.14562	0.80520	0.36188	-0.15622	0.80944
N (4b)	0.28684	0.73142	0.29692	0.28881	0.74391	0.29686
N (4b)	0.33454	0.00682	0.65212	0.33271	-0.00206	0.65161
H (4b)	0.29860	0.08530	0.70240	0.28604	0.08255	0.70561
H (4b)	0.31960	-0.06180	0.53940	0.31799	-0.08284	0.51496
H (4b)	0.37130	0.09120	0.61970	0.37284	0.11081	0.61558
H (4b)	0.37630	-0.07170	0.92270	0.37064	-0.07151	0.95627
H (4b)	0.39320	-0.23720	0.74750	0.41466	-0.22228	0.75324
H (4b)	0.31640	-0.22980	0.84760	0.32112	-0.28525	0.83067

Table S2: Our calculated (Cal) energy gaps (E_g in eV) for Perovskite (MA)₂Pb(SCN)₂I₂ bulk and monolayer using the different α and μ parameter values. HSE+SOC calculations with $\alpha = 0.43$ and $\mu = 0.2$ can predict E_g closer to the experimental (Exp) results. The direct E_g of (MA)₂Pb(SCN)₂I₂ are provided in parentheses.

	Cal		Exp
	$\alpha = 0.3, \mu = 0$	$\alpha = 0.43, \mu = 0.2$	
bulk	2.29 (2.43)	2.15 (2.26)	2.04 (2.11)
mono	2.78 (2.86)	2.63 (2.68)	—

Table S3: Predicted crystallographic parameters for 2D FE $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer in $\text{Pmc}2_1$ symmetry.

FE monolayer $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$			
$x = 35.00 \text{ \AA}, y = 9.05 \text{ \AA}, z = 8.09 \text{ \AA}, d_{\text{Pb}} = 0.382 \text{ \AA}$			
Atom	x	y	z
Pb (2a)	0.00000	0.23910	0.85342
I (2a)	0.00000	0.45508	0.54644
I (2a)	0.00000	0.09107	0.20575
S (4c)	-0.08517	0.21377	0.87874
C (4c)	-0.09739	0.32902	0.73433
C (4c)	0.89544	0.81653	0.82168
N (4c)	0.89409	0.41380	0.63000
N (4c)	-0.08377	0.67441	0.81034
H (4c)	-0.08791	0.38260	0.42097
H (4c)	-0.09319	0.60684	0.71551
H (4c)	-0.05502	0.30827	0.29471
H (4c)	-0.09964	0.88066	0.71041
H (4c)	-0.09386	0.12331	0.42829

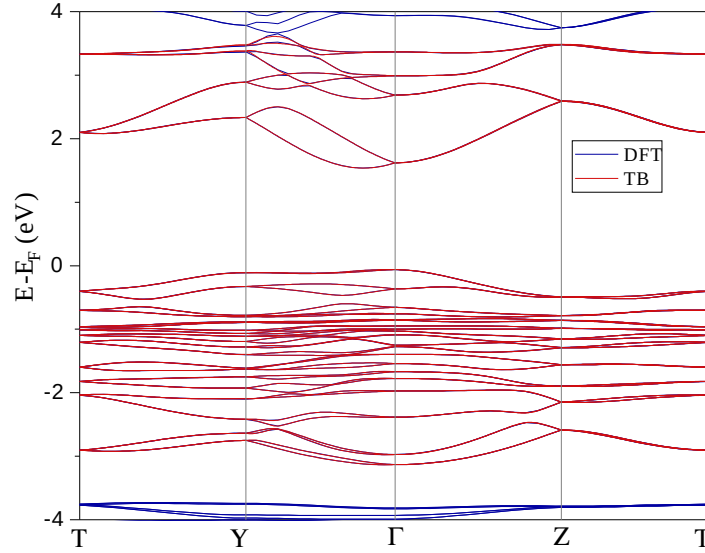


Figure S1: Comparison of band structure for $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer predicted by DFT(PBE functional) and tight binding (TB) fitting performed using the Wannier90 package. SOC effect is included self-consistently. Owing to the underestimation of band gap from PBE functional, scissor operator is applied to correct PBE band gap to the target value predicted HSE06 hybrid functional.

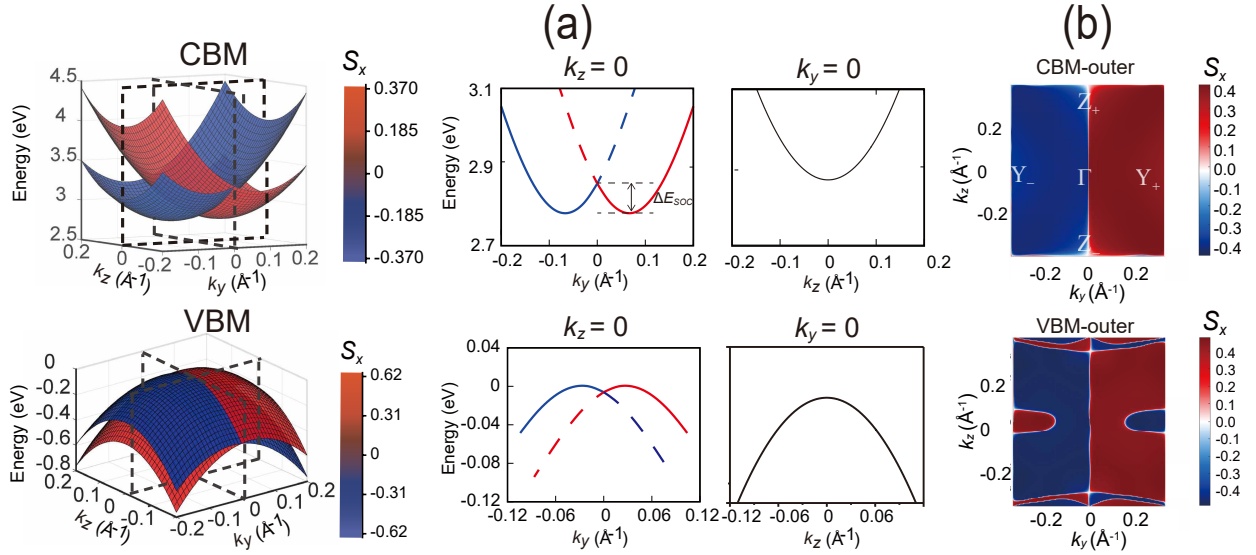


Figure S2: Spin-split band structure and spin texture obtained from $-P$ state of $\text{FE}-(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer. (a) 3D contour plots for spin-split top of the valence and bottom of conduction bands among 2D reciprocal space. The dispersion of spin-resolved and spin-degenerate energy bands along k_y and k_z directions are obtained, after projection 3D bands into $k_z = 0$ and $k_y = 0$ reciprocal planes, respectively. (b) Spin texture for outer valence and conduction band branches among 2D Brillouin zone of $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer, predicted by TB Hamiltonian constructed using Wannier90 package. Compared with the results from Fig. 3 of main text, switching of in-plane ferroelectricity between $+/-P$ states of $\text{FE}-(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer can lead to the reversal of out-of-plane spin orientations.

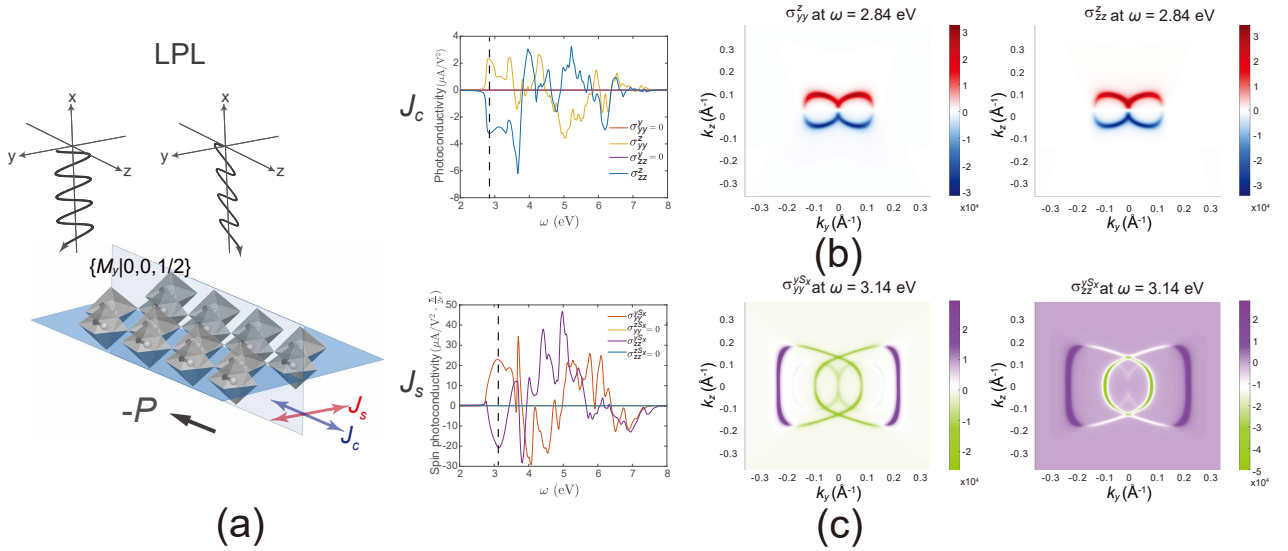


Figure S3: (a) Schematic diagram for generation of charge and spin photocurrent in $-P$ state of $\text{FE}-(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer, under the illumination of LPL propagating along x axis, polarized along y and z axes, respectively. The spatially separated pure spin and charge photocurrent (J_s and J_c) can be obtained along planar y and z axes, respectively. (b) Calculated J_C conductivity, and the distribution of non-zero conductivity components $\sigma_{yy}^z(k)$ and $\sigma_{zz}^z(k)$ among 2D reciprocal space of $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer, under the incident LPL with photon energy $\omega = 2.84$ eV. It is noted the color bar in different scales are marked for $\sigma_{yy}^z(k)$ and $\sigma_{zz}^z(k)$ components. (c) Calculated J_S conductivity with out-of-plane spin component S_x , and the distributions of non-zero spin current conductivity components $\sigma_{yy}^{Sx}(k)$ and $\sigma_{zz}^{Sx}(k)$ among 2D reciprocal space, under the incident LPL with photon energy $\omega = 3.14$ eV. It is clearly shown that LPL induced J_C and J_S will have their current directions reversed, when in-plane polarization is switched from $+P$ to $-P$ state of $\text{FE}-(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer.

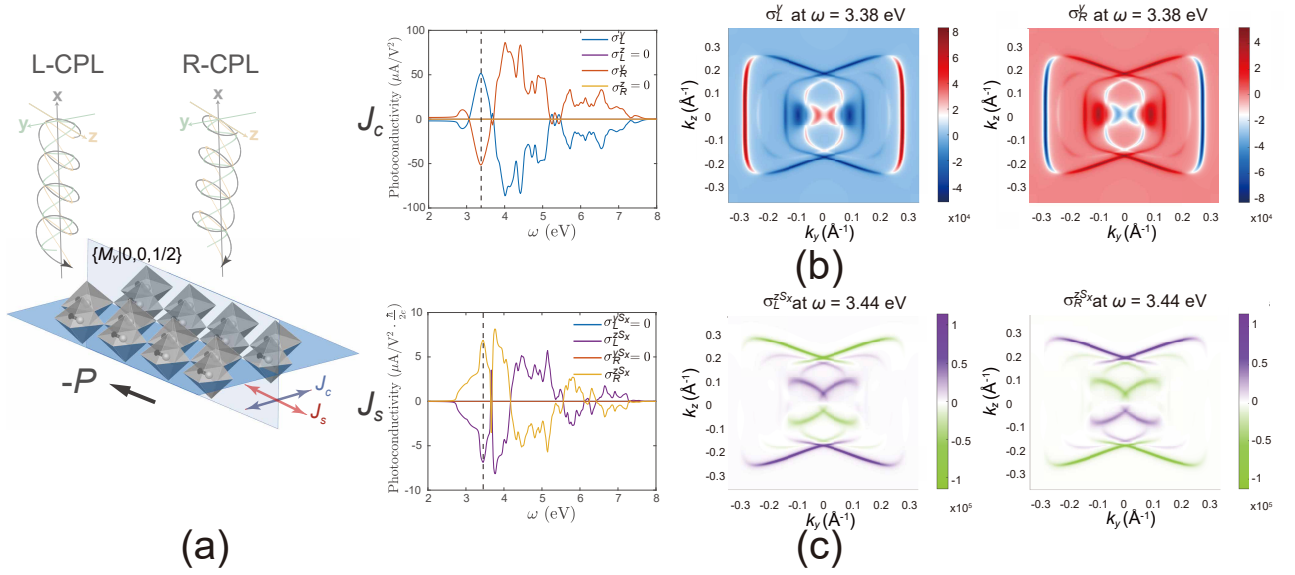


Figure S4: (a) Schematic diagram for generation of charge and spin photocurrent in $-P$ state of $\text{FE}-(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer, under the illumination of L-CPL and R-CPL propagating along x axis, respectively. The spatially separated pure charge and spin photocurrent (J_c and J_s) can be obtained along planar y and z axes, respectively. (b) Calculated J_c , and the distribution of non-zero conductivity components $\sigma_L^y(k)$ and $\sigma_R^y(k)$ among 2D reciprocal space of $(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer, under the incident L-CPL and R-CPL with photon energy $\omega = 3.38$ eV, respectively. (c) Calculated J_s conductivity with out-of-plane spin component S_x , and the distributions of non-zero spin current conductivity components $\sigma_L^{zS_x}(k)$ and $\sigma_R^{zS_x}(k)$ among 2D reciprocal space, under the incident L-CPL and R-CPL with photon energy $\omega = 3.44$ eV, respectively. It is clearly shown that CPL induced J_c and J_s will have their current directions reversed, when in-plane polarization is switched from $+P$ to $-P$ state of $\text{FE}-(\text{MA})_2\text{Pb}(\text{SCN})_2\text{I}_2$ monolayer.

References

- [1] Z. Xiao, W. Meng, B. Saparov, H.-S. Duan, C. Wang, C. Feng, W. Liao, W. Ke, D. Zhao, J. Wang *et al.*, *J. Phy. Chem. Lett.*, 2016, **7**, 1213–1218.