

Electronic Supplementary Information for

Quantum anomalous Hall effect with high and tunable Chern number in monolayer NdN₂

Shengshi Li,^{a,†} Xinyang Li,^{b,†} Weixiao Ji,^{a,*} Ping Li,^a Shishen Yan,^a and Changwen Zhang^{a,*}

^a Spintronics Institute & School of Physics and Technology, University of Jinan, Jinan, Shandong, 250022, People's Republic of China. E-mail: hellojiweixiao@163.com; E-mail: ss_zhangchw@ujn.edu.cn

^b School of Physics, Southeast University, Nanjing, Jiangsu, 211100, People's Republic of China.

[†] These authors contributed equally to this work.

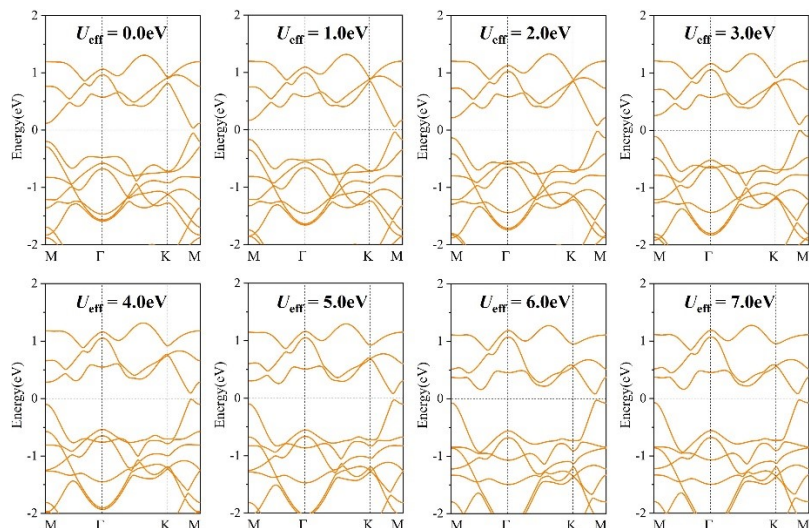


Fig. S1 Calculated SOC band structures of monolayer NdN₂ with different U_{eff} .

Table S1 Calculated lattice constant a , local magnetic moment of Nd atom M_{Nd} , sum of local magnetic moment of two N atoms M_{N_2} , and total magnetic moment per unit cell M_{tot} of monolayer NdN₂ using different pseudopotential for Nd atom.

	a (Å)	M_{Nd} (μ_{B})	M_{N_2} (μ_{B})	M_{tot} (μ_{B})
NdN ₂ ($5s^25p^65d^16s^2$)	3.93	0.06	2.94	3.00
NdN ₂ ($5s^26s^25p^65d^14f^3$)	3.92	0.13	2.87	3.00

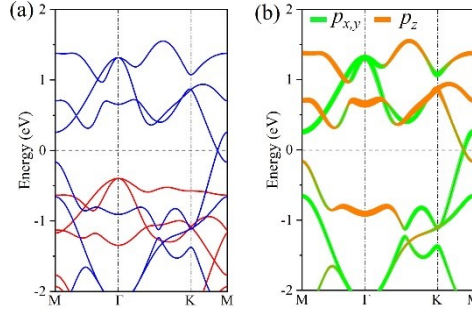


Fig. S2 Spin- and orbital-resolved band structures without SOC of monolayer NdN₂ obtained from GGA+*U* scheme with specific parameters of $U = 7.2$ eV and $J = 1.0$ eV implemented for the Nd-4*f* orbitals.

Table S2 Slater-Koster TB parameters for hexagonal lattice with lattice constant $a = 3.65\text{\AA}$ and interlayer distance $d = 1.77\text{\AA}$. The unprimed and primed parameters correspond to nearest and next-nearest neighbor hoppings, respectively.

ϵ_s	ϵ_{p_x, p_y}	ϵ_{p_z}	$V_{ss\sigma}$	$V_{sp\sigma}$	$V_{pp\sigma}$	$V_{pp\pi}$	$V'_{pp\sigma}$	$V'_{pp\pi}$
-2.8	-2.39	-2.41	-0.15	-1.94	-1.90	-0.55	-0.6	-0.2

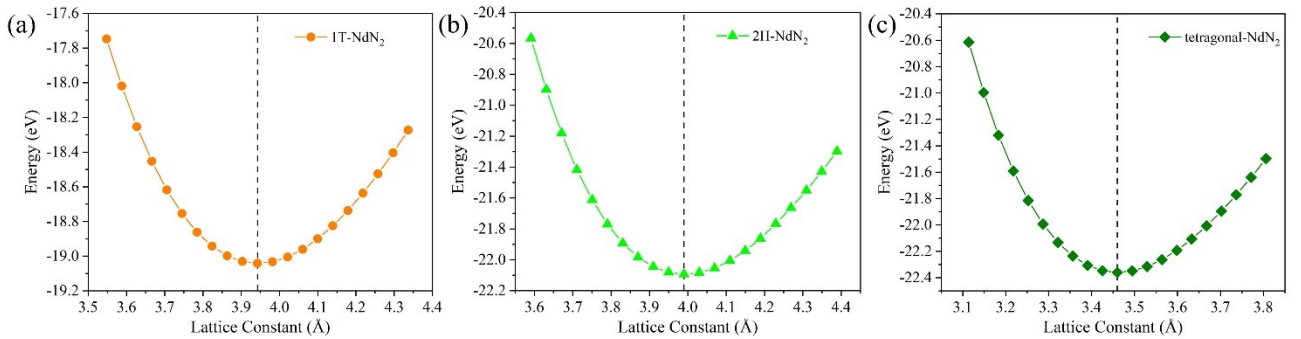


Fig. S3 Variation of energy as a function of lattice constant for monolayers 1T-NdN₂ (a), 2H-NdN₂ (b), and tetragonal-NdN₂ (c).

Part I: The MAE of monolayer NdN₂

Here, a detailed analysis is conducted on the magnetic anisotropy energy (MAE) of monolayer NdN₂. The angular dependence of MAE can be expressed as follows:

$$MAE(\theta, \varphi) = K_1 \cos^2 \theta + K_2 \cos^4 \theta + K_3 \cos^4 \varphi$$

where K_1 and K_2 are anisotropy constants, while θ and φ are polar and azimuth angles in the spherical coordinates, respectively. By fitting the MAE results shown in Fig. S3(a), we can obtain the anisotropy constants, *i.e.*, $K_1 = 143 \mu\text{eV}$, $K_2 = 35 \mu\text{eV}$, and $K_3 = 0 \mu\text{eV}$. The positive value of K_1 indicates that monolayer NdN₂ exhibits an out-of-plane magnetic anisotropy with its easy magnetization axis along the c axis. While $K_3 = 0$ implies isotropic magnetization in the xy plane. Figure S3(b) displays a spherical plot of the MAE arising from magnetization rotation, in which the out-of-plane magnetization is more energetically favorable than the in-plane magnetization by $171 \mu\text{eV/u.c.}$

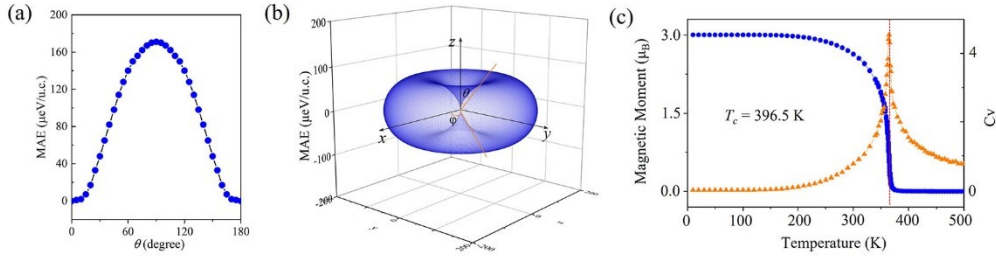


Fig. S4 (a) Polar angular dependence of MAE for monolayer NdN₂. (b) Magnetic anisotropy of NdN₂ through the whole space. (c) Variations of moment and C_v as a function of temperature.

Part II: The Curie temperature of monolayer NdN₂

Here, the Curie temperature (T_C) of monolayer NdN₂ is evaluated in detail. We first calculate exchange coupling parameters between nearest, next-nearest, and next-next-nearest neighbors (J_1 , J_2 , and J_3) using different magnetic configurations in relation to the Ising model:

$$H = -\sum_{i,j} J_1 M_i M_j - \sum_{k,l} J_2 M_k M_l - \sum_{m,n} J_3 M_m M_n \quad (\text{S2})$$

where M is the net magnetic moment at the N site i , while (i, j) , (k, l) , and (m, n) run over the nearest, next-nearest and next-next-nearest N atoms. The obtained J_1 and J_2 are 8.4 meV and 1.4 meV, respectively. However, the value of J_3 is only 0.13 meV, which is negligible and can be disregarded. The T_C is estimated through Monte-Carlo (MC) simulation on a 200×200 supercell for 1×10^9 loops, based on the Ising Hamiltonian. Figure S3(c) presents the evolution of magnetic moment per unit cell with increasing temperature. It should be noted that the magnetic moment decreases to $0.0 \mu_B$ when the temperature exceeds 396 K. Furthermore, we also calculate the heat capacity (C_v) expressed as:

$$C_v = \lim_{\Delta T \rightarrow 0} \frac{\Delta E_T}{\Delta T} \quad (\text{S3})$$

where ΔE_T is the variation of total energy as the temperature increases from T to $T + \Delta T$. According to the simulated C_v curve shown in Fig. S3(c), it can be inferred that the 2D NdN₂ exhibits a T_C of 396.5 K, indicating its potential application

in spintronic devices operating at room temperature.

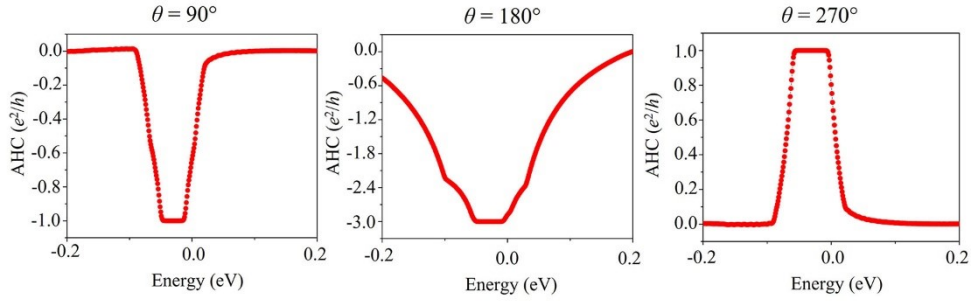


Fig. S5 Calculated AHC σ_{xy} as a function of Fermi level for the cases of $\theta = 90^\circ$, $\theta = 180^\circ$, and $\theta = 270^\circ$, respectively. The corresponding Chern numbers are $C = -1$, $C = -3$, and $C = 1$.

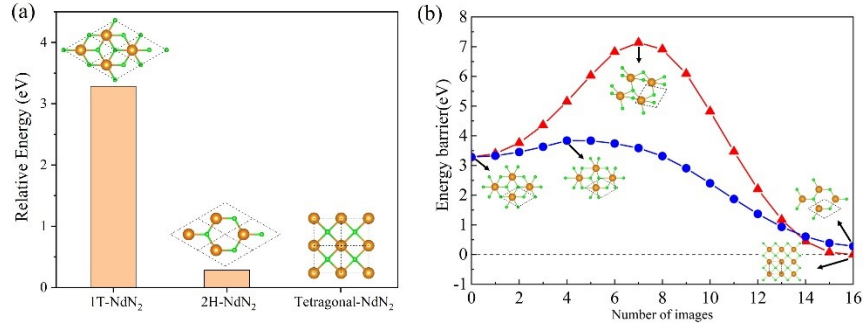


Fig. S6 (a) Relative energy of 2D allotropes of NdN_2 . (b) Energy curve along the structural phase transition paths of 1T-2H and 1T-tetragonal predicted by the c-NEB method.

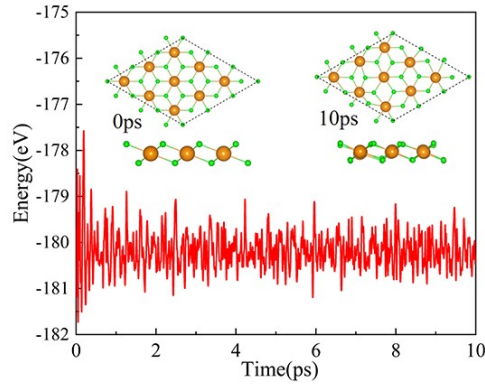


Fig. S7 Variation of total energy during 10ps AIMD simulation of monolayer 1T- NdN_2 . Inserts show the initial and final structures of monolayer 1T- NdN_2 after 10ps at 300 K.

Table S3 The lattice constant a , bond length of Nd-N d , N-N distance along the z axis h , and total magnetic moment per unit cell M_{tot} of monolayer NdN₂ with 1T, 2H, and tetragonal phases.

	a (Å)	d (Å)	h (Å)	M_{tot} (μ_B)
1T-NdN ₂	3.93	2.43	1.78	3.00
2H-NdN ₂	3.99	2.41	1.70	0.15
Tetragonal-NdN ₂	3.46	2.53	1.37	0.00

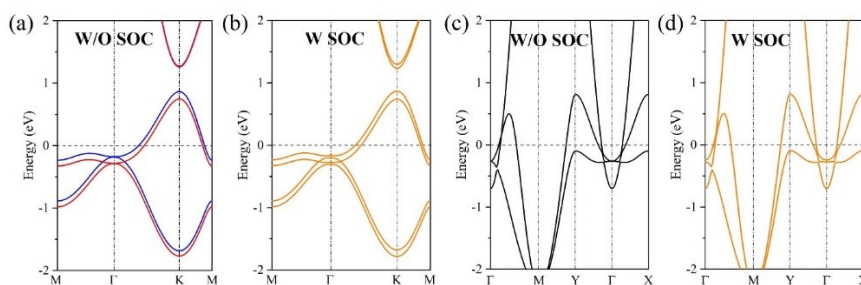


Fig. S8 Band structures of 1T-NdN₂ (a-b) and tetragonal-NdN₂ (c-d) monolayers without and with SOC. The red and blue lines represent the spin-up and spin-down states, respectively.