

Supplementary Information

# Electronic and Structural Properties of Mixed-Cation Hybrid Perovskites Studied by Efficient Spin-Orbit Included DFT-1/2 Approach

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## S1. DOS

The total densities of states (DOS) for single-cation perovskites calculated by PBE\_non-SOC (blue), PBE (red), and PBE-1/2 (black) are shown in [Figure S1](#).

## S2. Tetragonal Phase of **mapi**

In addition to the pseudo cubic phase of **mapi**, we also studied its tetragonal phase, which is stable at room temperature. [Figure S2](#) shows structure of the tetragonal phase of **mapi** as well as its band structure calculated using PBE-1/2. Each unit cell contains two

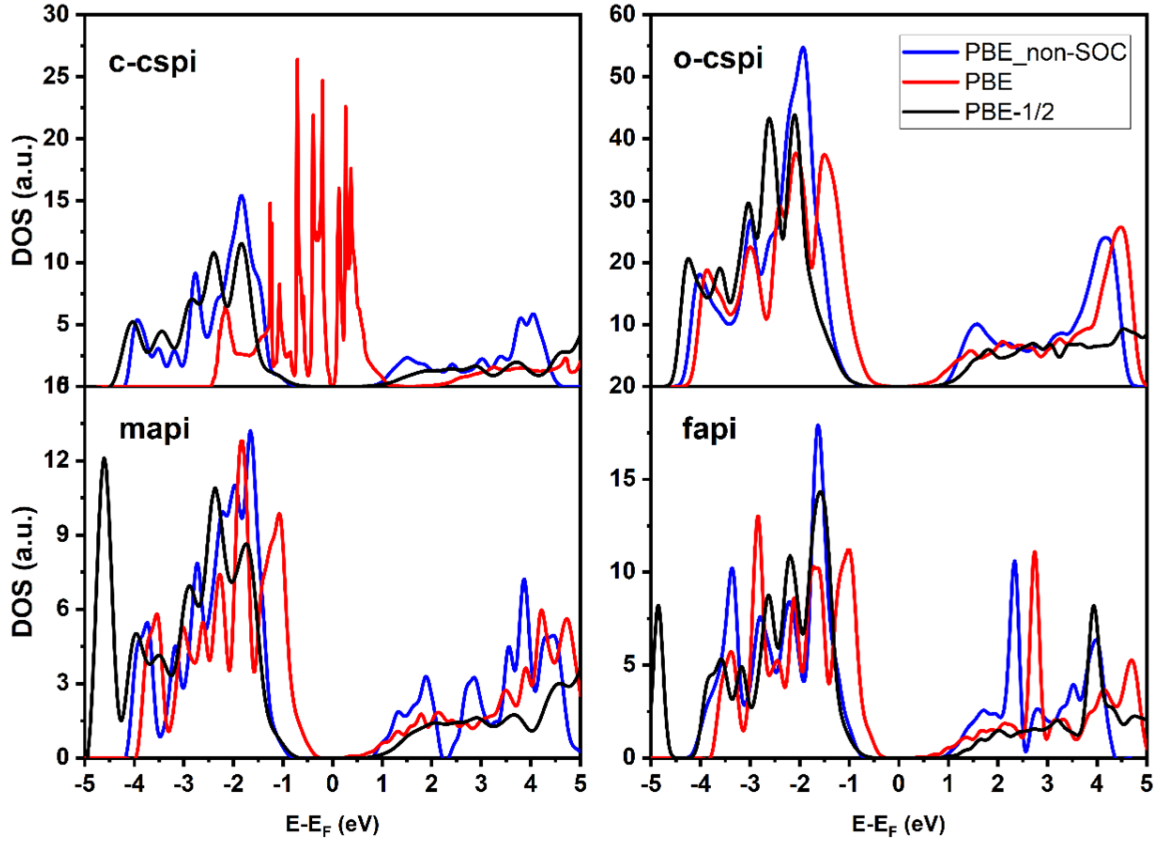


Figure S1: Total DOS for c-cspi, o-cspi, mapi, and fapi using PBE functional without SOC (PBE\_non-SOC, blue), with SOC (PBE, red) and PBE with SOC (PBE-1/2, black.)

MA cation which are oriented antiparallel to each other. Contrary to the cubic phase, there is a double degeneracy at the  $\Gamma$  point and no Rashba band splitting is present. The direct band gap at the  $\Gamma$  point is 1.7 eV.

### S3. Rashba Band Splitting

The Rashba interaction coefficient,  $\alpha$  is given as:

$$\alpha = \frac{2\epsilon^\#}{\Delta k}$$

where  $\epsilon^\#$  is the difference between the split curves at the extreme points (see [Figure S3](#)) and  $\Delta k$  is the respective position of the  $\epsilon^\#$  alongside the  $k$  axis in reciprocal space. Here,  $\epsilon^\#$  and thus  $\alpha$  differ for VB and CB and we define  $\alpha_+$ ( $\alpha_-$ ) for left (right) hand side of the point

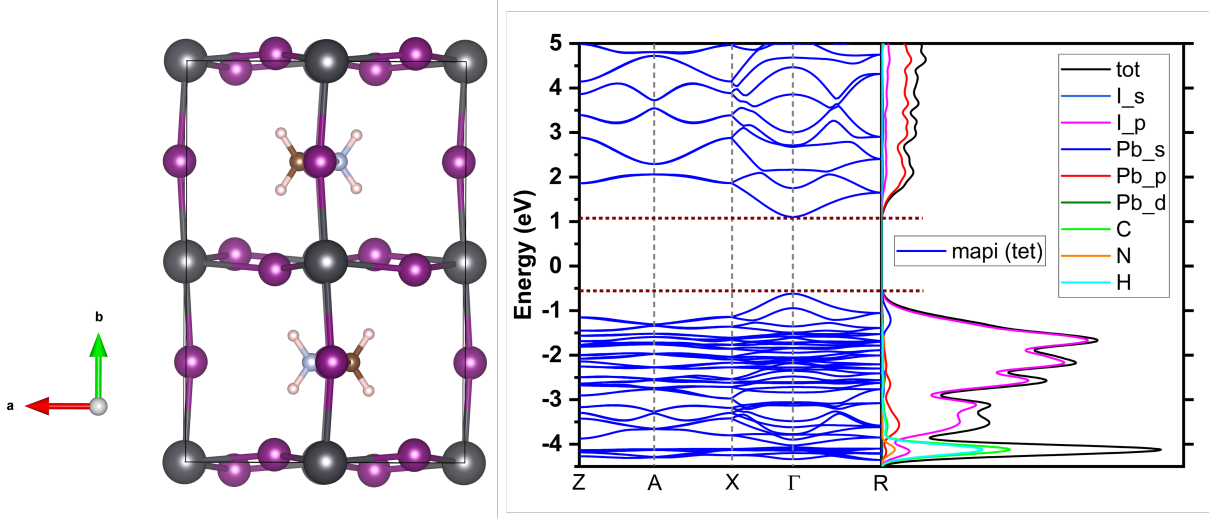


Figure S2: Left: Structure of the tetragonal phase of **mapi**. Each unit cell contains two MA cations which are aligned antiparallel to each other. Right: Band structure and PDOS for tetragonal **mapi**, using PBE-1/2. The Fermi energy is set to zero.

that band gap is appeared alongside. Therefore, a large (small) difference between  $\alpha_+$  and  $\alpha_-$  indicates a one-sided (two-sided) Rashba splitting.

Since the Rashba parameter ( $\alpha$ ) determines the derivation of the band curvature from a perfect parabola and the indirect nature of the band gap, we focus on the difference between the Rashba parameter in VBM and CBM ( $\Delta\alpha = \alpha_C - \alpha_V$ ). [Figure S4](#) shows  $\Delta\alpha_+$  ( $\alpha_{+C} - \alpha_{+V}$ ) and  $\Delta\alpha_-$  ( $\alpha_{-C} - \alpha_{-V}$ ), which are denoted by L and R, respectively. Both PBE and PBE-1/2 show the same trends for **mapi**. On the quantitative level, if PBE-1/2 corrections are included,  $\alpha$  is reduced from 3.2 to 2.4. An increase of about 1.4 is obtained for the right hand side of **fapi**. This large increase is related to the modification of CBM; as it is seen in [Figure 4](#) (single-cation band structures), the PBE-1/2 mostly modifies the CBM while VBM does not bear a sensible splitting. One might expect that the correction of the self-energy in PBE-1/2 also improves the quantitative description of  $\alpha$ .

Parameters of Rashba band splitting for the mixed systems are listed in [Table 3](#). The Rashba parameter decreases systematically from **fapi** to **fascpi** to **famascpi**. This is in full agreement with the discussed changes to the Pb-I bond distribution. The more uniform the bonds, the smaller the Rashba split. [Figure S5a](#) shows the left/right  $\alpha$  parameters for all systems under study. Disregarding the organic or inorganic nature of an additional cation,

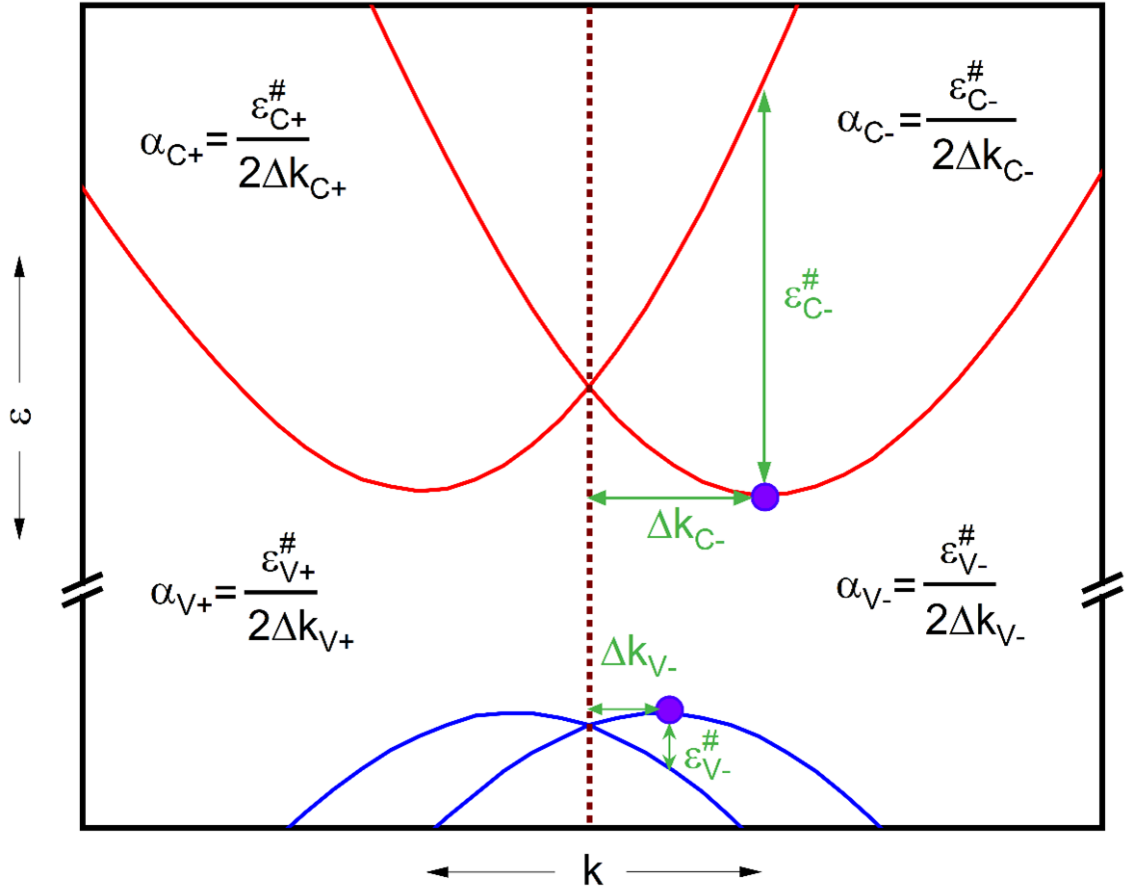


Figure S3: Rashba parameter,  $\alpha$ , in VB and CB according to derivation of energy bands with respect to an ideal parabolic band. V and C indices devote to valence and conduction band, respectively.

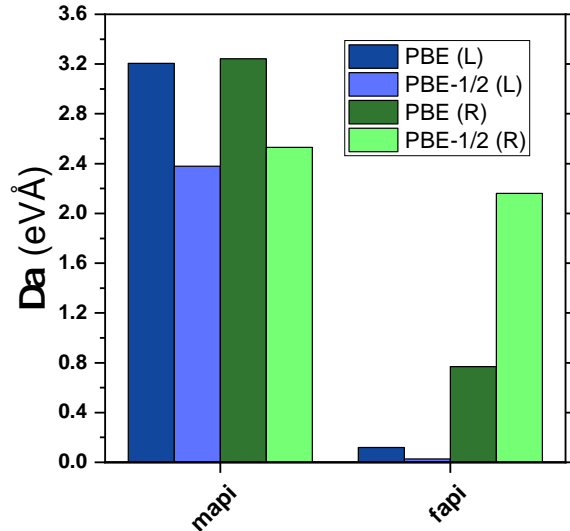


Figure S4: Difference between Rashba parameters in VB and CB ( $\Delta\alpha$ ) for **mapi** and **fapi** using PBE and PBE-1/2.

it reduces  $\alpha$  significantly. In the case of **facspi**, the nature of the one-sided splitting changes almost to a two-sided form. The reduction of Rashba band splitting intensity is shown in [Figure S5b](#). The triple **facspi** case has the lowest  $\Delta\alpha$  which means its band nature has the potential for a higher carrier recombination rate. This finding is in agreement with the experimental observations reported by Huang et. al. [1].

#### S4. Born effective charge

[Table S1](#) lists the Born effective charges for the studied single-cation systems. For the Born charges of iodine it is convenient to distinguish between the Born charge parallel to the Pb-I bond ( $Z_{||}$ ) and the two directions perpendicular to it ( $Z_{\perp}$ ). Compare solid and dashed arrows in [Figure S6](#). Note that [Table S1](#) lists the projections of  $Z$  on Cartesian directions (a,b,c). For **o-cspi**, the  $\text{PbI}_3$  scaffold is not aligned along these directions (see [Figure S6](#)) and thus the projection of  $Z$  is reduced compared to **c-cspi**.

#### S5. References

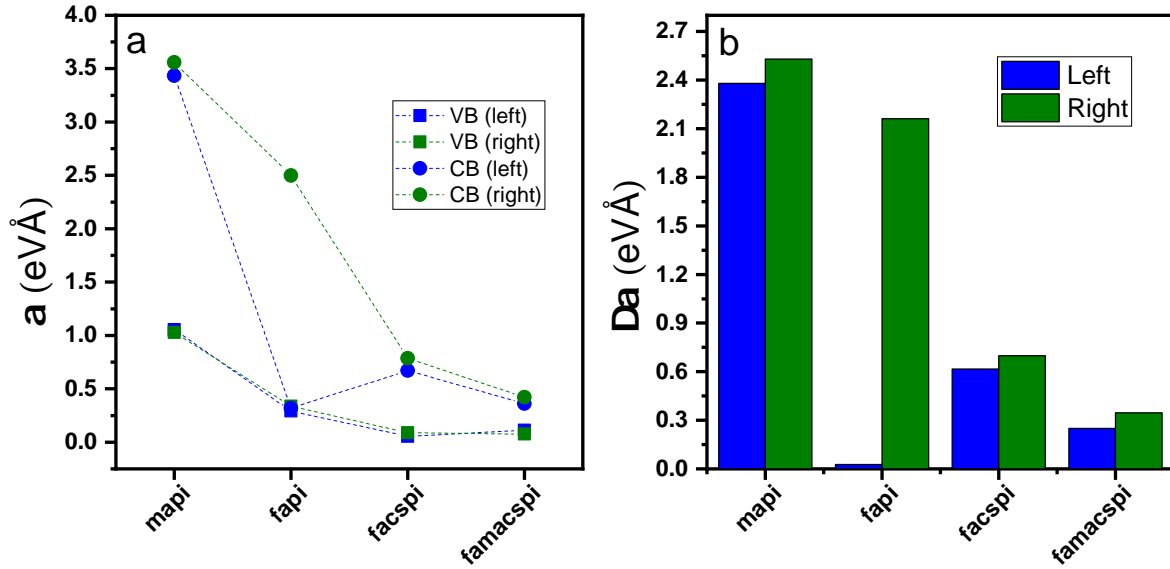


Figure S5: Rashba parameter for **mapi**, **fapi**, **facspi**, and **famacspi** using PBE-1/2. (a):  $\alpha$  in left and right hand side of VB and CB with respect to  $R$  and  $\Gamma$  points for single- and mixed-cation perovskites, respectively. (b):  $\Delta\alpha$  as the difference of VB and CB in left and right hand side.

## References

- [1] Huang, Z. *et al.* Band Engineering via Gradient Molecular Dopants for CsFA Perovskite Solar Cells. *Advanced Functional Materials* **31**, 2010572 (2021). URL <https://onlinelibrary.wiley.com/doi/10.1002/adfm.202010572>. 5

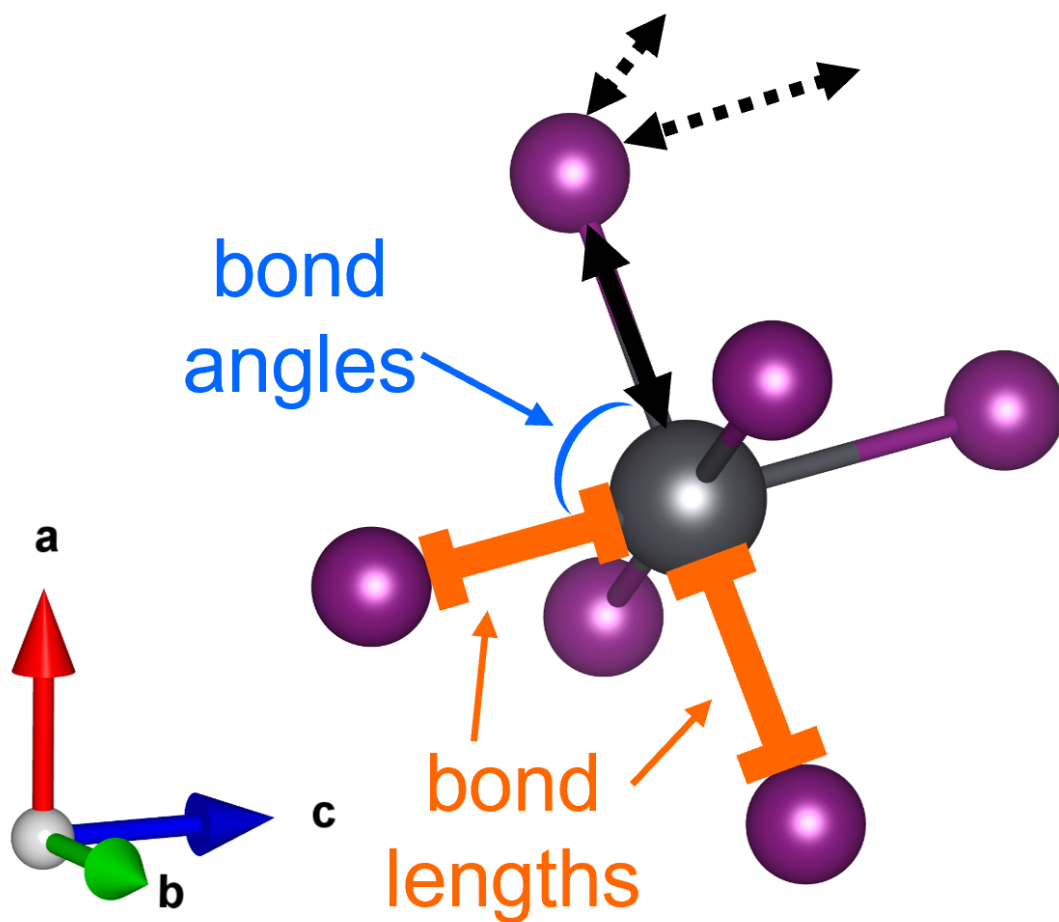


Figure S6: Illustration of one  $\text{Pb-I}_3$  scaffold in **o-cspi**. Orange lines mark the different Pb-I bond length in that scaffold and one I-Pb-I angle is exemplary marked. For each iodine, one can distinguish one direction parallel to the Pb-I bond ( $\parallel$ ) and two perpendicular directions ( $\perp$ ) as marked by solid and dashed arrows for the top I. While  $\parallel$  is approximately along the (pseudo)-cubic directions in **c-cspi** and **mapi**, the bonds deviate from (100), (010), and (001) for **o-cspi**.

Table S1: Impact of the chosen potential on the Born effective charges  $Z_{ii}^*$  for  $i : x, y, z$ . Note that  $Z_{xx}^* = Z_{yy}^* = Z_{zz}^*$  for cubic **cspi**, only. In case of pseudo-cubic **mapi**, the mean value of  $Z_{ii}^*$  is given. Upper part: Pb and I present in all systems. For I, it is convenient to define  $Z_{||}$ , the Born charge along the Cartesian direction with the smallest angle relative to the Pb-I bond and  $Z_{\perp}$ , the average of both other directions (refer to new figure with scaffold). Lower part: Total values for A-ions.

		PBE_non-SOC	PBE	PBE-1/2	PBE_non-SOC		PBE		PBE-1/2	
		Pb			I <sub>  </sub>	I <sub>⊥</sub>	I <sub>  </sub>	I <sub>⊥</sub>	I <sub>  </sub>	I <sub>⊥</sub>
<b>c-cspi</b>	$Z_{ii}$	4.83	5.49	4.19	-4.68	-0.76	-6.27	-0.29	-3.94	-0.78
	$Z_{xx}$	4.22	4.61	3.84	-2.35	-0.82	-2.55	-0.81	-2.13	-0.85
	$Z_{yy}$	4.35	4.75	3.92	-2.44	-0.75	-2.64	-0.75	-2.20	-0.78
<b>o-cspi</b>	$Z_{zz}$	4.53	5.03	4.01	-4.31	-0.80	-4.86	-0.77	-3.70	-0.84
	$Z_{ii}$	4.40	4.78	3.93	-4.25	-0.72	-4.64	-0.71	-3.64	-0.78
	$Z_{xx}$	5.18	6.00	4.48	-5.08	-0.64	-6.12	-0.56	-4.22	-0.72
<b>fapi</b>	$Z_{yy}$	3.96	4.20	3.54	-3.67	-0.80	-4.01	-0.74	-3.17	-0.82
	$Z_{zz}$	4.77	5.25	4.22	-4.67	-0.79	-5.28	-0.74	-3.93	-0.86
			PBE_non-SOC	PBE	PBE-1/2	PBE_non-SOC		PBE		PBE-1/2
		Cs in <b>c-cspi</b>			CH <sub>3</sub> NH <sub>3</sub> in <b>mapi</b>					
$Z_{ii}$		1.38	1.37	1.31	1.28		1.28		1.18	
		Cs in <b>o-cspi</b>			CH(NH <sub>2</sub> ) <sub>2</sub> in <b>fapi</b>					
$Z_{xx}$		1.30	1.30	1.28	1.18		1.22		1.17	
$Z_{yy}$		1.28	1.28	1.26	1.30		1.30		1.26	
$Z_{zz}$		1.38	1.38	1.35	1.48		1.51		1.42	