

Electronic Supplementary Information

Temperature effects on the internal conversion of excited adenine and adenosine

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SI-1. Excited-state characterization

Table S1. Excitation energies, oscillator strengths (f), and state characterization given by natural transition orbitals and their respective contributions (in percentage) calculated for 9H-adenine and *anti*-adenosine at the S_0 minimum at ADC(2)/SV(P) level.

State	Adenine			Adenosine		
	Energy (eV)	f	State characterization	Energy (eV)	f	State characterization
S ₁	5.16	0.001	$n\pi^*$ (0.96)	5.13	0.000	$n\pi^*$ (0.95)
S ₂	5.34	0.060	$\pi\pi^*$ (0.71) $\pi\pi^*$ (0.26)	5.25	0.091	$\pi\pi^*$ (0.74) $\pi\pi^*$ (0.21)
S ₃	5.52	0.269	$\pi\pi^*$ (0.89)	5.45	0.283	$\pi\pi^*$ (0.86)

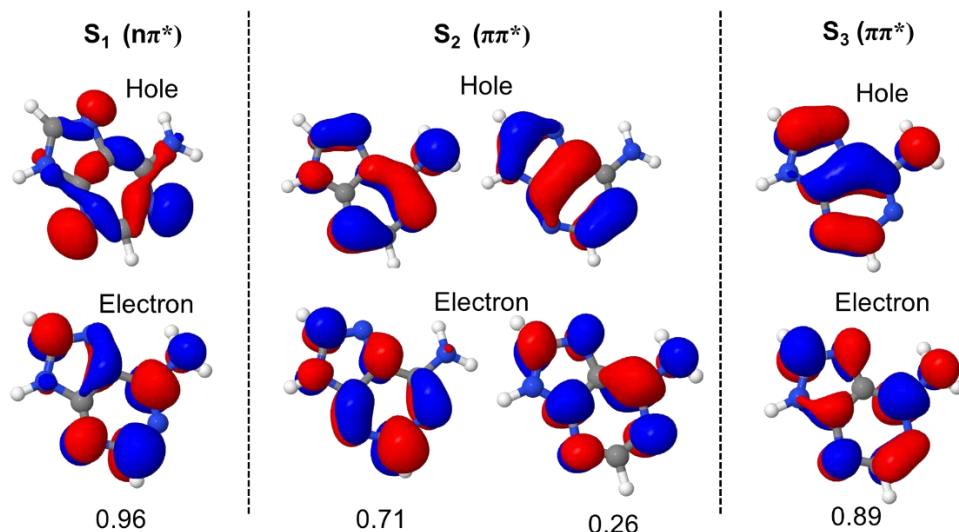


Figure S1. Natural transition orbitals (NTOs) describing the first singlet excitations on 9H-adenine at the ground state geometry. NTOs computed at ADC(2)/SV(P) level.

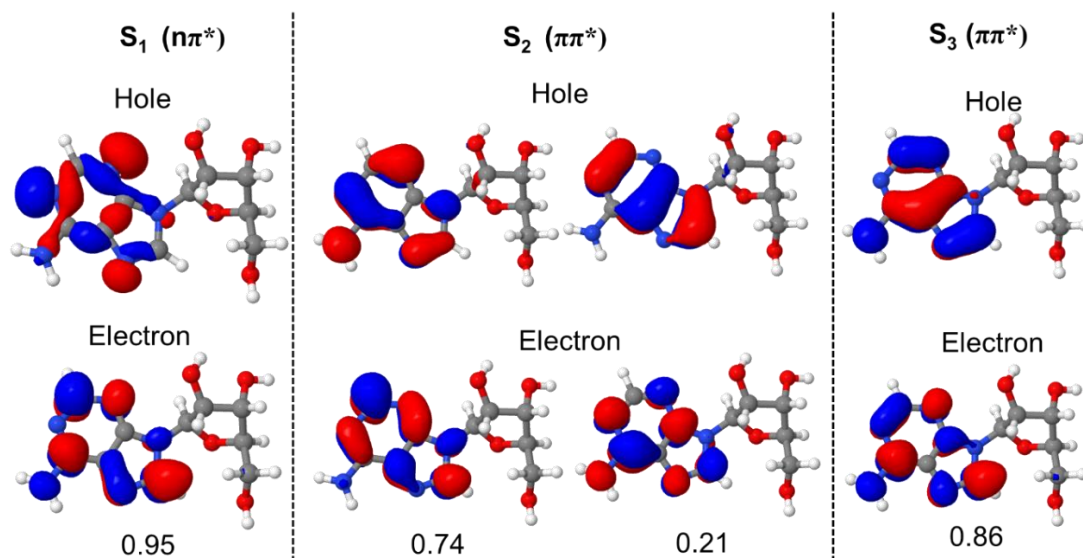


Figure S2. Natural transition orbitals (NTOs) describing the first singlet excitations on *anti*-adenosine at the ground state geometry. NTOs computed at ADC(2)/SV(P) level.

Vertical excitations and potential energy surfaces of adenine were explored at different levels of theory (see Table S2 and supplementary Ref. 1). Reaction pathways connecting ground state minimum and conical intersections for adenine at the multiconfigurational level are also explored in previous works.²⁻⁴ It has been shown that ADC(2) can satisfactorily reproduce both vertically excited states and reaction pathways for adenine.

Table S2. Comparison of excitation energies and oscillator strengths (f) obtained at different levels of theory for 9H-adenine. L_b and L_a states denote, respectively, the weakly absorbing $\pi\pi^*$ state and the bright $\pi\pi^*$ state.

9H-Adenine				
Level of theory	$n\pi^*$	$\pi\pi^*$ (L_b)	$\pi\pi^*$ (L_a)	Reference
RI-ADC(2)/SV(P)	5.16 (0.001)	5.34 (0.060)	5.52 (0.269)	This work.
RI-ADC(2)/aug-cc-pVDZ ^a	5.00 (0.03)	5.07 (0.10)	5.06 (0.17)	Aquino <i>et al.</i> ⁵
CC2/aug-cc-pVTZ ^b	5.12 (0.01)	5.25 (0.02)	5.25 (0.30)	Fleig <i>et al.</i> ⁶
MS-CASPT2(12,10)/ANO-L ^c	4.85	5.09	5.44 (0.18)	Conti <i>et al.</i> ⁷
MS-CASPT2/aug-cc-pVTZ ^d	5.04 (0.02)	5.02 (0.02)	5.12 (0.40)	Silva-Junior <i>et al.</i> ⁸
MS-CASPT2(16,12)/6-31G(d) ^e	5.29	5.43	5.56	Barbatti <i>et al.</i> ⁹
MS-CASPT2(16,12)/6-31G(d) ^f	4.96	5.03	5.17	Barbatti <i>et al.</i> ⁹
MS-CASPT2(16,13)/6-31G(d) ^g	4.96	5.16	5.35	Serrano-Andrés <i>et al.</i> ¹⁰
MS-CASPT2/6-31+G(d) ^h	5.31 (0.010)	5.42 (0.039)	5.61 (0.152)	Barbatti <i>et al.</i> ⁴
MS6-CASPT2(14,10)/VDZ ⁱ	4.86	5.13	5.28	Marchetti <i>et al.</i> ³
Experimental (gas phase) ^j			5.16 ± 0.07 (0.24)	Barbatti <i>et al.</i> ⁴
Experimental (H ₂ O) ^l			4.76	Voet <i>et al.</i> ¹¹

^aGS geometry optimized at MP2/TZVP level in C1 point-group symmetry. ^bGS geometry optimized at CC2/aug-cc-pVTZ level in C1 point group symmetry. ^cCASPT2//CASSCF(12,10)/ANO-L state averaged (SA) on six roots. GS geometry optimized at CASSCF(8,8)/6-31G* level of theory. ^dGS geometry optimized at MP2-6-31G* in C_s point group symmetry, SA-8 for $\pi\pi^*$ states, and SA-3 for $n\pi^*$ states. Active space (0 13)12 for $\pi\pi^*$ and (3 13)18 for $n\pi^*$ (active space orbitals of A' and A'' symmetry given in parentheses followed by the number of electrons). ^eCASPT2//SA5-CASSCF(16,12)/6-31G* method used for single-point calculations on geometries optimized at the MR-CIS(6,5)/SA-3-CAS(12,10)/6-31G* level; IPEA shift of 0.25. ^fCASPT2//SA5-CASSCF(16,12)/6-31G* method used for single-point calculations on geometries optimized at the MR-CIS(6,5)/SA-3-CAS(12,10)/6-31G* level; IPEA 0.00. ^gCASPT2//CASSCF(16,13)/6-31G* method. ^hCAS(18,13); Level shift 0.2 a.u.; IPEA 0.25 a.u. ⁱGeometry optimization with CASSCF(10,8)/VDZ and energies calculated with MS6-CASPT(14,10)/VDZ. ^jEstimated from gas phase band maximum data. ^lData from UV spectra in water, pH 7.0.

Adenosine vertical excitation energies resemble the ones from adenine, as shown in our ADC(2) calculations and also experimentally¹¹ (see Tables S2 and S3). Similarly to adenine, the S₁ state of adenosine has a nπ* character at the Franck-Condon region. For the syn-isomer, CASPT2 calculations also identify a ππ*/S₀ MECI with a puckered geometry lying lower than the S₁ state at the Franck-Condon region. EDPT reaction pathway from the bright ππ* (L_a) state to this intersection is barrierless, but from the S₁ (nπ*) to the ππ*/S₀ MECI, a barrier of 0.3 eV is observed.³ Reaction pathways computed at both ADC(2) and CASPT(2) levels for adenosine (Figure 15 and Figure S34 of Ref. 3) show that ADC(2) can reproduce the main features.

Table S3. Comparison of excitation energies and oscillator strengths (*f*) obtained at different levels of theory for adenosine. L_b and L_a states denote, respectively, the weakly absorbing ππ* state and the bright ππ* state.

Adenosine				
Level of theory	nπ*	ππ* (L _b)	ππ* (L _a)	Reference
RI-ADC(2)/SV(P)	5.13 (0.001)	5.25 (0.060)	5.45 (0.269)	This work
ADC(2)/cc-pVDZ	5.20	5.18	5.30	Tuna et al. ¹²
MS6-CASPT2(10,10)/VDZ ^a	5.0	5.55	5.65	Marchetti <i>et al.</i> ³
Experimental (H ₂ O) ^j			4.78	Voet <i>et al.</i> ¹¹

^a Optimizations with MP2/6-31G(d) and vertical energies calculated with MS6-CASPT2(10,10)/VDZ. ^j Data from UV spectra in water, pH 7.9.

SI-2. Nuclear-ensemble spectra

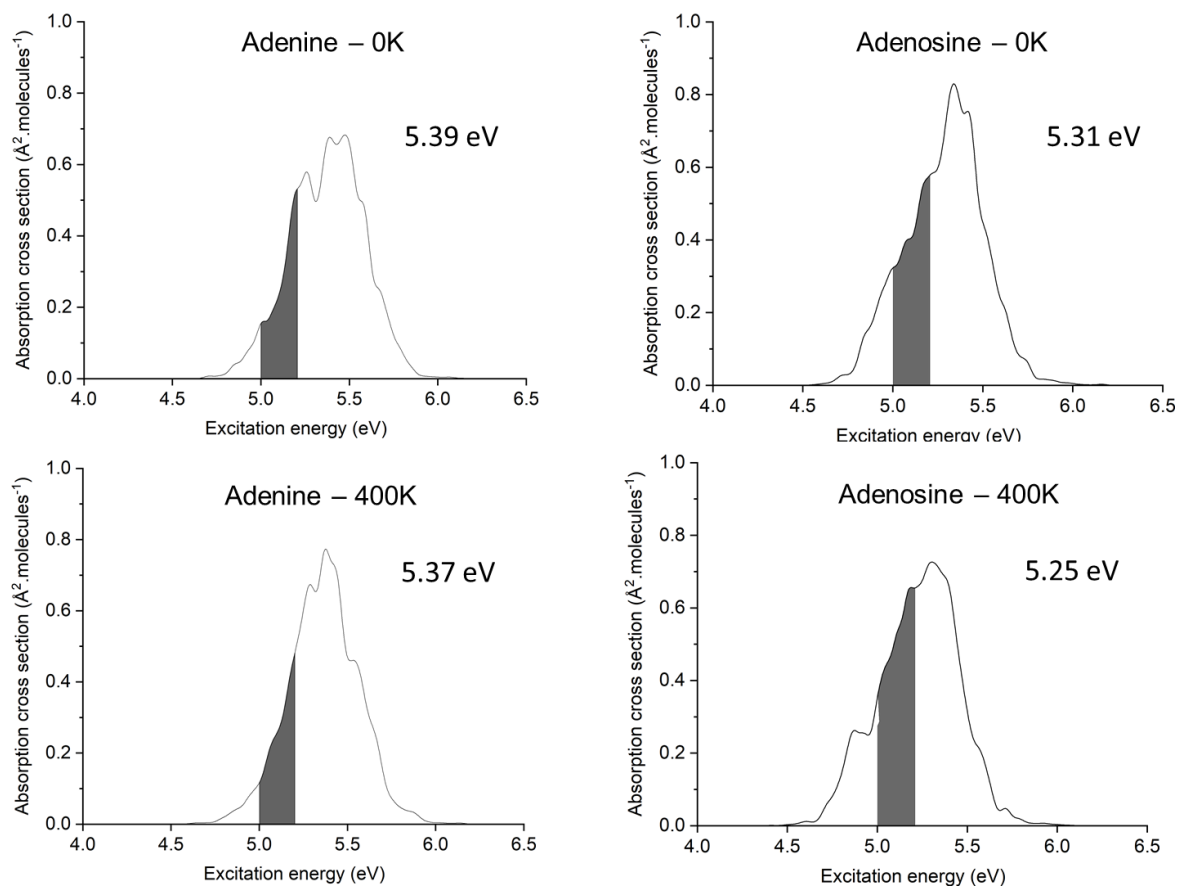


Figure S3. Computed absorption spectra based on 500 points nuclear ensemble for adenine and adenosine at 0 and 400 K calculated at ADC(2)/SV(P) level. The grey area represents the excitation window (5.1 ± 0.1 eV in all cases).

SI-3. Excited-state occupation

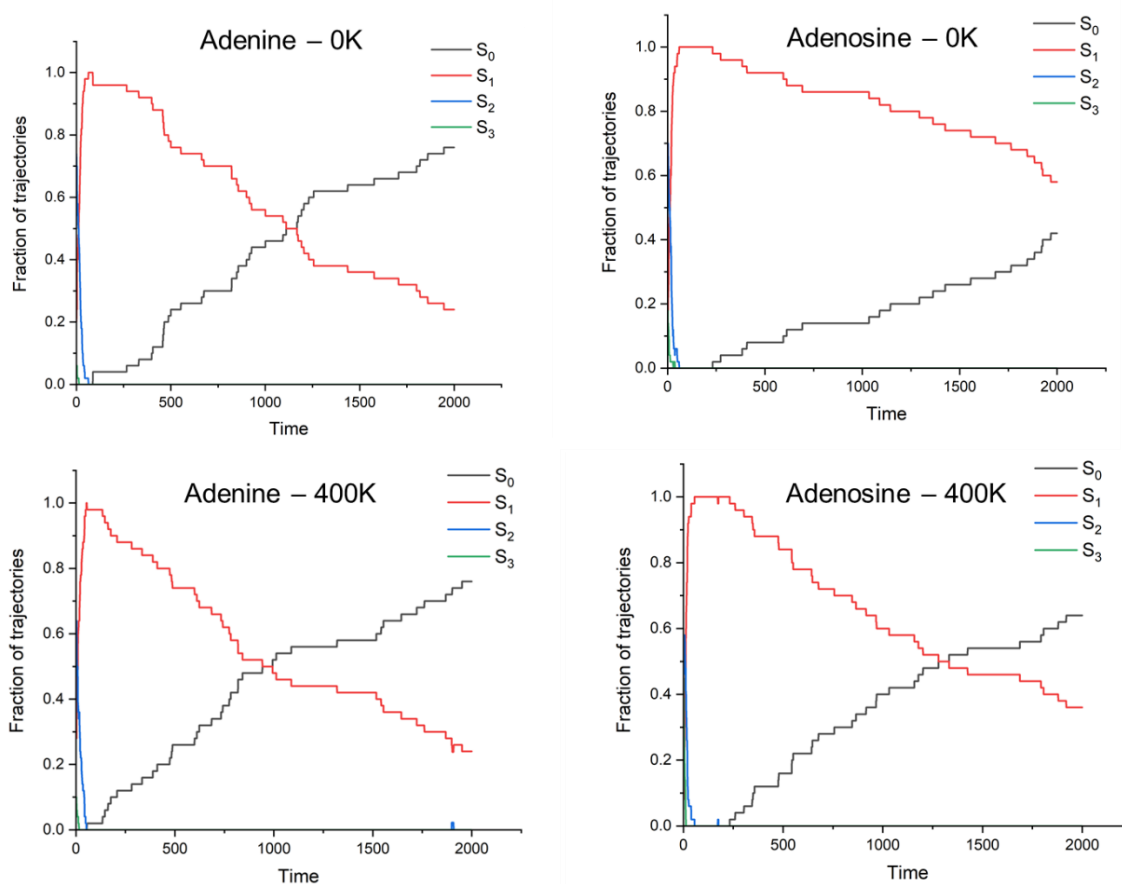


Figure S4. Adiabatic state occupation evolution for adenine and adenosine at 0 and 400 K calculated at ADC(2)/SV(P) level.

SI-4. Electronic density differences

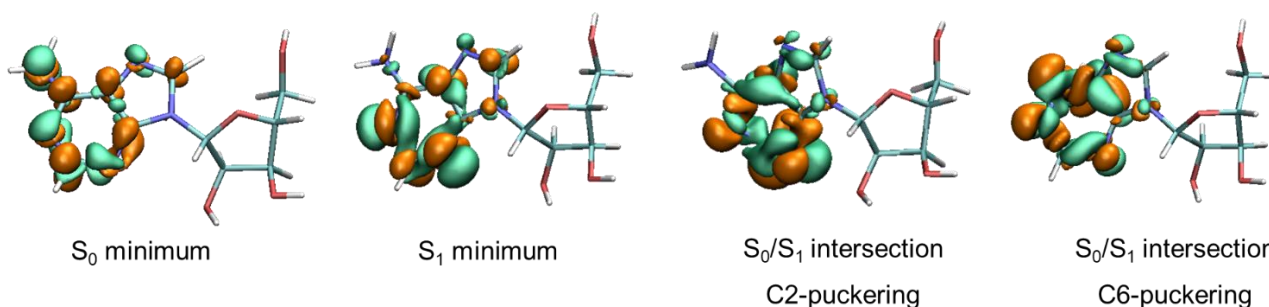


Figure S5. Electronic density difference between S_1 and S_0 states at the S_0 minimum, S_1 minimum, C2-puckered and C6-puckered intersections. Green and orange regions represent a decrease and an increase in electron density, respectively.

SI-5. Normal mode analysis

Normal modes are linear combinations of cartesian coordinates, typically obtained as the eigenvectors of the mass-weighted Hessian matrix of the energy. They are handy in identifying relevant nuclear vibrational motions of polyatomic molecules. The normal mode analysis is a way to project the molecular motion of a trajectory in terms of the corresponding normal modes displacement of the molecule.^{13, 14}

Let $\mathbf{Q} = (\mathbf{q}_1 | \mathbf{q}_2 | \dots | \mathbf{q}_{3N-6})$ be the matrix containing the $3N - 6$ normal modes at a reference geometry \mathbf{R}_0 in Cartesian coordinates, where N is the number of atoms in a molecule. $\mathbf{R}(i, t)$ is the Cartesian coordinates of the i^{th} trajectory at time t . The displacement vector in normal mode coordinates (i.e., normal mode amplitudes) with respect to the reference geometry can be written as

$$\mathbf{\Omega}(i, t) = \mathbf{Q}^{-1}[\mathbf{R}(i, t) - \mathbf{R}_0]$$

Taking an average over all trajectories, the average displacement $\bar{\mathbf{\Omega}}(t)$ for each normal mode can be obtained to monitor the average motion (nuclear vibrations) of the molecule as a function of time. Moreover, the standard deviation of this average displacement over time boils down to a single number per normal mode to represent its coherent activity:

$$\hat{\mathbf{\Omega}}_{coh}^2 = \frac{1}{k_{fin} - k_{ini}} \sum_{k=k_{ini}}^{k_{fin}} \bar{\mathbf{\Omega}}(k\Delta t)^2 - \left(\frac{1}{k_{fin} - k_{ini}} \sum_{k=k_{ini}}^{k_{fin}} \bar{\mathbf{\Omega}}(k\Delta t) \right)^2$$

In this equation, $\hat{\mathbf{\Omega}}_{coh}^2$ shows the coherent activity of a normal mode during the dynamics. k_{ini} and k_{fin} are the first and last time steps considered, and Δt is the constant time interval of the dynamics.

Alternatively, one may compute the total standard deviation over time steps and trajectories without prior averaging over the trajectories. Hence, it will quantify the total motion observed along the normal modes.

$$\hat{\mathbf{\Omega}}_{i,total}^2 = \frac{1}{k_{fin} - k_{ini}} \sum_{k=k_{ini}}^{k_{fin}} \mathbf{\Omega}(i, k\Delta t)^2 - \left(\frac{1}{k_{fin} - k_{ini}} \sum_{k=k_{ini}}^{k_{fin}} \mathbf{\Omega}(i, k\Delta t) \right)^2$$

Because we aim to describe the total motion, our results discussed in the paper correspond to this second quantity.

SI-6. Active normal modes

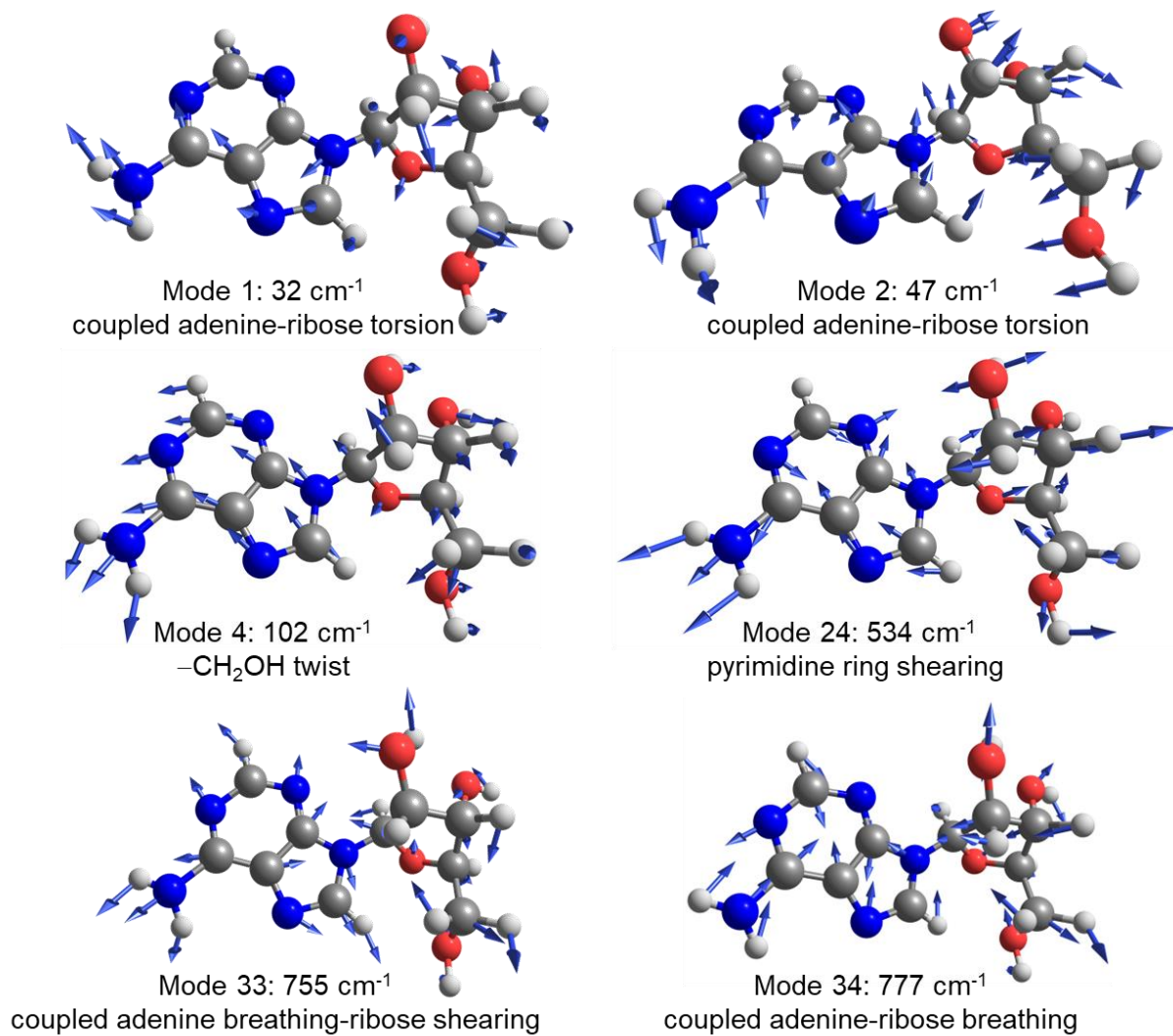


Figure S6. Main normal modes contributing to vibrational energy transfer in adenosine.

SI-7. Supplementary references

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