## Supporting Information

## Janus CrSSe Monolayer with an Interesting Ferromagnetism

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## Computational Details

The particle swarm optimization (PSO) method within the evolutionary algorithm as implemented in the Crystal structure AnaLYsis by Particle Swarm Optimization (CALYPSO) code ${ }^{1,2}$ was applied to find new structures of CrSSe monolayers. Unit cells containing 1, and 2 formula units (f.u.) were considered. In the first step, random structures with certain symmetry are constructed in which atomic coordinates are generated by the crystallographic symmetry operations. Local optimizations using the VASP code ${ }^{3}$ were done with the conjugate gradients method and stopped when Gibbs free energy changes became smaller than $1 \times 10^{-5} \mathrm{eV}$ per cell. After processing the first-generation structures, $60 \%$ of them with lowest Gibbs free energies are selected to construct the next generation structures by PSO. $40 \%$ of the structures in the new generation are randomly generated. A structure fingerprinting technique of bond characterization matrix is applied to the generated structures, so that identical structures are strictly forbidden. These procedures significantly enhance the diversity of the structures, which is crucial for structural global search efficiency. In most cases, structural searching simulations for each calculation were stopped after generating 1000~1200 structures (e.g., about $20 \sim 30$ generations).

The local structural relaxations and electronic property calculations were performed in the framework of the density functional theory (DFT) ${ }^{4}$ within the generalized gradient approximation (GGA) ${ }^{5}$ as implemented in the VASP code. The projector augmented wave (PAW) method is used to treat the ion-electron interaction, in which the valence electrons of each atom are $\mathrm{Cr}: 3 d^{5} 4 s^{1}, \mathrm{~S}$ : $3 s^{2} 3 p^{4}, \mathrm{Se}: 4 s^{2} 4 p^{4}$.

We calculated the cohesive energy of the Janus CrSSe monolayer by using the following expression:

$$
E_{c o h}=\left(E_{C r}+E_{S}+E_{S e}-E_{C r S S e}\right) / 3,
$$

where $E_{\mathrm{Cr}}, E_{\mathrm{S}}, E_{\mathrm{Se}}$ and $E_{C r S S e}$ are the total energies of a single Cr atom, S atom, Se atom and the Janus CrSSe monolayer, respectively.

The Young's modulus $Y(\theta)$ and Poisson's ratio $v(\theta)$ along any direction $\theta(\theta$ is the angle with respect to the positive $x$-direction) are defined as

$$
\begin{array}{r}
\frac{C_{11} C_{22}-C_{12}^{2}}{C_{11} s^{4}+C_{22} c^{4}+\left(\frac{C_{11} C_{22}-C_{12}^{2}}{C_{66}}-2 C_{12}\right) c^{2} s^{2}} \\
C_{12}\left(c^{4}+s^{4}\right)-\left(C_{11}+C_{22}-\frac{C_{11} C_{22}-C_{12}^{2}}{C_{66}}\right) c^{2} s^{2} \\
C_{11} s^{4}+C_{22} c^{4}+\left(\frac{C_{11} C_{22}-C_{12}^{2}}{C_{66}}-2 C_{12}\right) c^{2} s^{2}
\end{array}
$$

were $\mathrm{c}=\cos \theta$ and $\mathrm{s}=\sin \theta$.
Based on the Heisenberg model, the different magnetic configurations for the Janus CrSSe monolayers are given by the following expressions:

$$
\begin{aligned}
& E_{F M}=E_{0}-4 J_{1} \vec{S}^{2}-4 J_{2} \vec{S}^{2} \\
& E_{A F M 1}=E_{0}-4 J_{1} \vec{S}^{2}+4 J_{2} \vec{S}^{2} \\
& E_{A F M 2}=E_{0}+4 J_{1} \vec{S}^{2}+4 J_{2} \vec{S}^{2}
\end{aligned}
$$

were $\mathrm{E}_{0}$ is the energy without magnetic coupling, and $S$ is set to 1.5 . Hence, we can estimate the magnetic coupling parameters $J_{1}$ and $J_{2}$ by solving these three equations.

The Curie temperature was calculated by the Monte Carlo method with the Metropolis algorithm based on the Heisenberg model. All of them were implemented in MCSOLVER. ${ }^{6}$

## Supporting Figures and Tables

Table S1. Structural parameters of the predicted Janus CrSSe monolayer: space group (SG), lattice constants $a$ and $b$, layer thickness $h$, bond lengths of $\mathrm{Cr}-\mathrm{S} d_{1}$ and $\mathrm{Cr}-\mathrm{Se} d_{2}$, and bond angles of $\mathrm{Cr}-\mathrm{S}-$ $\operatorname{Cr} \theta_{1}$ and $\mathrm{Cr}-\mathrm{Se}-\mathrm{Cr} \theta_{2}$.

|  |  |  |  |  |  | $d_{1}$ | $d_{2}$ | $\theta_{1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\theta_{2}$ |  |  |  |  |  |  |  |  |
|  | SG | $a$ <br> $(\AA)$ | $b$ <br> $(\AA)$ | $(\AA)$ <br> $(\AA)$ | $\AA)$ |  <br> $\left({ }^{\circ}\right)$ | $\left({ }^{\circ}\right)$ |  |
| CrSSe | Pmm 2 | 3.58 | 3.65 | 2.66 | 2.17 | 2.32 | 111.0 | 104.0 |



Fig. S1. ELF maps of different interfaces in the Janus CrSSe monolayer.


Fig. S2. (a) Phonon dispersion curves, (b) total energies and the snapshots of the final frame, and (c, d) orientation-dependent in plane Young's modulus $Y(\theta)$ and Poisson's ratio $v(\theta)$ of the Janus CrSSe monolayer.


Fig. S3. Calculated exfoliation energy of the Janus CrSSe monolayer deposited on the $\mathrm{Cu}(001)$ surface.

To find out whether the Janus CrSSe monolayer can be synthesized experimentally on a substrate, we have examined its stability on several substrates, such as $\mathrm{Cu}(001), \mathrm{Ag}(001), \mathrm{Au}(001)$ and graphene, which are widely used for surface growth in experiments. Finally, our calculations show that the Janus CrSSe monolayer remains stable on the $\mathrm{Cu}(001)$ substrate. The mismatch between the Janus $\operatorname{CrSSe}(a=7.165 \AA)$ and the $\mathrm{Cu}(001)(a=7.213 \AA)$ is $0.7 \%$, which is suitable for the formation of the heterostructure made by 2D Janus CrSSe and $\mathrm{Cu}(001)$ surface. Subsequently, to explore the possibility of fabricating the Janus CrSSe monolayer deposited on the Cu ( 001 ) surface, we have simulated the exfoliation process and predicted exfoliation energy with respect to separation. The calculated exfoliation energy of the Janus CrSSe monolayer is $0.7 \mathrm{~J} / \mathrm{m}^{2}$, which is comparable to the $\delta-\mathrm{InP}_{3}\left(0.827 \mathrm{~J} / \mathrm{m}^{2}\right)^{7}$ and $\mathrm{SnP}_{3}\left(0.71 \mathrm{~J} / \mathrm{m}^{2}\right),{ }^{8}$ indicating the Janus CrSSe monolayer could be prepared experimentally from $\mathrm{Cu}(001)$ surface from using similar approaches as graphene by mechanical cleavage or liquid phase exfoliation.
(a)

FM
(b)

AFM1
(c)

AFM2

Fig. S4. Considered different magnetic configurations: FM, AFM1, and AFM2 of the Janus CrSSe monolayer in a $2 \times 2 \times 1$ supercell. Red and blue arrows denote spin-up and spin-down states, respectively.

Table S2. Total energies of the Janus CrSSe monolayer under different magnetic configurations. (Considering the SOC selects the direction of the spin state along out-of-plane.)

|  |  | FM | AFM1 | AFM2 |
| :---: | :---: | :---: | :---: | :---: |
| CrSSe | GGA+U | $\mathbf{- 6 1 . 3 3 5 0 1 4}$ | -59.533385 | -60.667903 |
|  | GGA+U+SOC | $\mathbf{- 6 1 . 4 5 4 1 2 0}$ | -59.651125 | -60.788104 |



Fig. S5. PDOS onto five partial $3 d$ orbitals of Cr atoms of the FM Janus CrSSe monolayer.


Fig. S6. Band structures and spin-resolved PDOS of the FM Janus CrSSe monolayer at the hybrid HSE06 level. The red and blue lines indicate the spin-up and spin-down states, respectively.


Fig. S7. Geometry of structures of $\mathrm{CrS}_{2}$ and $\mathrm{CrSe}_{2}$ monolayers.

Table S3. Structural parameters of the $\mathrm{CrS}_{2}$ and $\mathrm{CrSe}_{2}$ monolayers: space group (SG), lattice constants $a$ and $b$, layer thickness $h$, bond lengths of $\mathrm{Cr}-\mathrm{S} d_{1}$ and $\mathrm{Cr}-\mathrm{Se} d_{2}$, bond angles of $\mathrm{Cr}-\mathrm{S}-\mathrm{Cr} \theta_{1}$, $\mathrm{Cr}-\mathrm{Se}-\mathrm{Cr} \theta_{2}$ and cohesive energy $E_{\text {coh }}$.

| Systems | SG | $\begin{gathered} a \\ (\AA) \end{gathered}$ | $\begin{gathered} b \\ (\AA) \end{gathered}$ | $\begin{gathered} h \\ (\AA \AA) \end{gathered}$ | $\begin{gathered} d_{1} \\ (\AA) \end{gathered}$ | $\begin{gathered} d_{2} \\ (\AA) \end{gathered}$ | $\begin{aligned} & \theta_{1} \\ & \left({ }^{\circ}\right) \end{aligned}$ | $\begin{aligned} & \theta_{2} \\ & \left({ }^{\circ}\right) \end{aligned}$ | $\begin{aligned} & E_{\text {coh }} \\ & (\mathrm{eV}) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{CrS}_{2}$ | $P-4 m 2$ | 3.53 | 3.53 | 2.56 | 2.18 | / | 108.2 | 1 | 3.92 |
| $\mathrm{CrSe}_{2}$ | $P-4 m 2$ | 3.70 | 3.70 | 2.77 | / | 2.31 | 1 | 106.3 | 3.46 |



Fig. S8. Calculated structural stabilities of (a, c) $\mathrm{CrS}_{2}$ and $(\mathrm{b}, \mathrm{d}) \mathrm{CrSe}_{2}$ monolayers. (e, f$) Y(\theta)$ and $v$ $(\theta)$.

Table S4. Elastic constants ( $C_{11}, C_{12}, C_{66}$, in $\mathrm{N} / \mathrm{m}$ ), Young's modulus and Poisson's ratio of the $\mathrm{CrS}_{2}$
and $\mathrm{CrSe}_{2}$ monolayers.

|  | $C_{11}$ <br> $(\mathrm{~N} / \mathrm{m})$ | $C_{12}$ <br> $(\mathrm{~N} / \mathrm{m})$ | $C_{66}$ <br> $(\mathrm{~N} / \mathrm{m})$ | $Y_{\max }$ | $Y_{\text {min }}$ | $V_{\text {max }}$ | $V_{\text {min }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{CrS}_{2}$ | 45.39 | 10.14 | 10.71 | 43.12 | 30.91 | 0.44 | 0.22 |
| $\mathrm{CrSe}_{2}$ | 34.74 | 7.83 | 8.82 | 32.97 | 24.93 | 0.41 | 0.23 |



Fig. S9. Angular dependence of the magnetic anisotropy energy (MAE) of the (a) $\mathrm{CrS}_{2}$ and (b) $\mathrm{CrSe}_{2}$ monolayers.


Fig. S10. Projected band structures with spin-up and spin-down states for (a, c) $\mathrm{CrS}_{2}$, and (b, d) $\mathrm{CrSe}_{2}$ monolayers. The Fermi level is set to 0 eV . The 2.31 and 2.22 eV direct band-gap in spindown channel are illustrated.


Fig. S11. (a) Calculated average magnetic moment of Cr atoms and (b) magnetic susceptibility with respect to temperature of the $\mathrm{CrS}_{2}$ and $\mathrm{CrSe}_{2}$ monolayers based on the Heisenberg model from MC simulations.

Table S5. Data related to the $\mathrm{CrS}_{2}$ and $\mathrm{CrSe}_{2}$ monolayers, including local magnetic moment on $\mathrm{Cr}, \mathrm{S}$ and Se atoms, magnetic exchange constants $J_{1}$ and $J_{2}$, evaluated Curie temperatures ( $T_{C}$ ), maximum magnetic anisotropy energy ( $\mathrm{MAE}_{\text {max }}$ ) and the band-gaps $\left(E_{g}\right)$.

| systems | $\mathrm{M}_{\mathrm{Cr}}$ <br> $\left(\mu_{\mathrm{B}}\right)$ | $\mathrm{M}_{\mathrm{S}}$ <br> $\left(\mu_{\mathrm{B}}\right)$ | $\mathrm{M}_{\mathrm{Se}}$ <br> $\left(\mu_{\mathrm{B}}\right)$ | $J_{I}$ <br> $(\mathrm{meV})$ | $J_{2}$ <br> $(\mathrm{meV})$ | $T_{C}$ <br> $(\mathrm{~K})$ | $\mathrm{MAE}_{\text {max }}$ <br> $(\mu \mathrm{eV} / \mathrm{Cr})$ | $E_{g}$ <br> $(\mathrm{eV})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{CrS}_{2}$ | 2.93 | -0.36 | $/$ | 49.92 | $/$ | 820 | 7.07 | $0 / 2.31$ |
| $\mathrm{CrSe}_{2}$ | 3.24 | $/$ | -0.46 | 50.47 | $/$ | 835 | 411.65 | $0 / 2.22$ |

Table S6. Total energies of the $\mathrm{CrS}_{2}$ and $\mathrm{CrSe}_{2}$ monolayers under different magnetic configurations. (Considering the SOC selects the direction of spin state along the out-of-plane.)

|  |  | FM | AFM1 | AFM2 |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{CrS}_{2}$ | GGA+U | $\mathbf{- 6 4 . 2 8 2 3 3 2}$ | -62.493469 | -63.293510 |
|  | GGA+U+SOC | $\mathbf{- 6 4 . 3 1 0 9 2 3}$ | -62.513796 | -63.317936 |
| $\mathrm{CrSe}_{2}$ | GGA+U | $\mathbf{- 5 8 . 3 8 0 7 0 5}$ | -56.542859 | -57.526540 |
|  | GGA+U+SOC | $\mathbf{- 5 8 . 5 8 7 8 8 4}$ | -56.771095 | -57.739471 |

Table S7. Structural information of the predicted $2 \mathrm{D} \mathrm{CrSSe}, \mathrm{CrS}_{2}$, and $\mathrm{CrSe}_{2}$ monolayers.

|  | Space Group | Lattice Parameters$\left(\AA,{ }^{\circ}\right)$ | Wyckoff <br> Positions <br> (fractional) |  | $y$ | $z$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | Atoms | $x$ |  |  |
| CrSSe | Pmm 2 | $a=3.5872$ | $\mathrm{Cr}(1 \mathrm{a})$ | 0.00000 | 0.00000 | 0.49820 |
|  |  | $b=3.6538$ | $\mathrm{Se}(1 \mathrm{~b})$ | 0.00000 | 0.50000 | 0.53705 |
|  |  | $\alpha=\beta=\gamma=90.00$ | S (1c) | 0.50000 | 0.00000 | 0.46475 |
| $\mathrm{CrS}_{2}$ | $P-4 m 2$ | $a=b=3.5293$ | $\mathrm{Cr}(1 \mathrm{~d})$ | 0.00000 | 0.00000 | 0.50000 |
|  |  | $\alpha=\beta=\gamma=90.00$ | S(2g) | 0.00000 | 0.50000 | 0.53546 |
| $\mathrm{CrSe}_{2}$ | $P-4 m 2$ | $a=b=3.7006$ | $\mathrm{Cr}(1 \mathrm{~d})$ | 0.00000 | 0.00000 | 0.50000 |
|  |  | $\alpha=\beta=\gamma=90.00$ | $\mathrm{Se}(2 \mathrm{~g})$ | 0.00000 | 0.50000 | 0.53724 |

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