

Supplementary Material for: Designing Barrier-Free Metal/MoS₂ Contacts through Electrene Insertion

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1 Metal/MoS₂ Contacts

To begin, several important features of metal/MoS₂ contacts are briefly discussed in order to understand the impact of electrene insertion at this interface. Table S1 summarizes the structural properties and exfoliation energies of the studied metal/MoS₂ interfaces. Fig. S1 shows the structures of these interfaces, as well as the computed band structures, MoS₂ densities of states, electrostatic potential plots, and Bader charge transfer.

The Schottky barrier height (SBH) for metal/MoS₂ contacts can be determined by locating the MoS₂ conduction band minimum with respect to the Fermi level of the interface. Despite the hybridization between the metal and MoS₂ states (particularly in the case of Cu), the MoS₂ band edge can be identified upon comparing the material-projected electronic

Table S 1: Number of atoms, \mathbf{k} -point mesh, in-plane lattice constant (a), strain on the metal slab (ϵ_m), and interlayer spacing between the metal and the semiconductor (\bar{d}) for the Au/MoS₂ and Cu/MoS₂ heterojunctions. Also shown are exfoliation energies computed using either unstrained (unstr) or strained (str) lattice constants for the isolated metal slab, the strain energies, and the decomposition of the strained exfoliation energies into base functional and D3 dispersion contributions. All energies are in meV/Å², using the area of the strained unit cell.

Interface	Au/MoS ₂	Cu/MoS ₂
# atoms	33	105
\mathbf{k} -point mesh	$7 \times 7 \times 1$	$4 \times 4 \times 1$
$a (= b)$ (Å)	5.46	9.47
ϵ_m (%)	3.26	2.45
\bar{d} (Å)	2.82	2.21
$E_{\text{exfo,unstr}}$	3.5	59.8
E_{strain}	48.2	36.7
$E_{\text{exfo,str}}$	51.7	96.5
$E_{\text{base,str}}$	-7.8	21.9
$E_{\text{disp,str}}$	59.5	74.6

bands of the interface with the electronic bands of individual layers (metal and MoS₂) in the same supercell configuration, as depicted in Fig. S2. From the electronic band energies shown in Fig. S1(b,f), the SBHs of Au/MoS₂ and Cu/MoS₂ interfaces are calculated to be ~ 0.5 eV and 0.4 eV, respectively, in reasonable agreement with literature values.¹⁻³

Pristine monolayer MoS₂ physisorbs on Au, as is evident from the dispersion contribution to the exfoliation energy (Table S1), which is slightly higher than the total due to non-bonded repulsion from the base density functional. As such, the top-contacted configuration of the Au/MoS₂ interface shown in Fig. S1(a) has a vdW gap, which causes a tunnelling barrier for electron transport.⁴ The electrostatic potential profile of this heterostructure is depicted in Fig. S1(d) and is calculated as the sum of ionic and Hartree contributions, averaged over the xy -plane perpendicular to the interface. The red bar in Fig. S1(d) shows the range of potential values at the point with the highest average potential within the vdW gap. As the minimum of this range lies above the Fermi level, this indicates the presence of a tunnelling barrier. Since the electrostatic potential (relative to the Fermi level) is positive across the whole contact region, no electrons pass through the contact without tunnelling; this is in

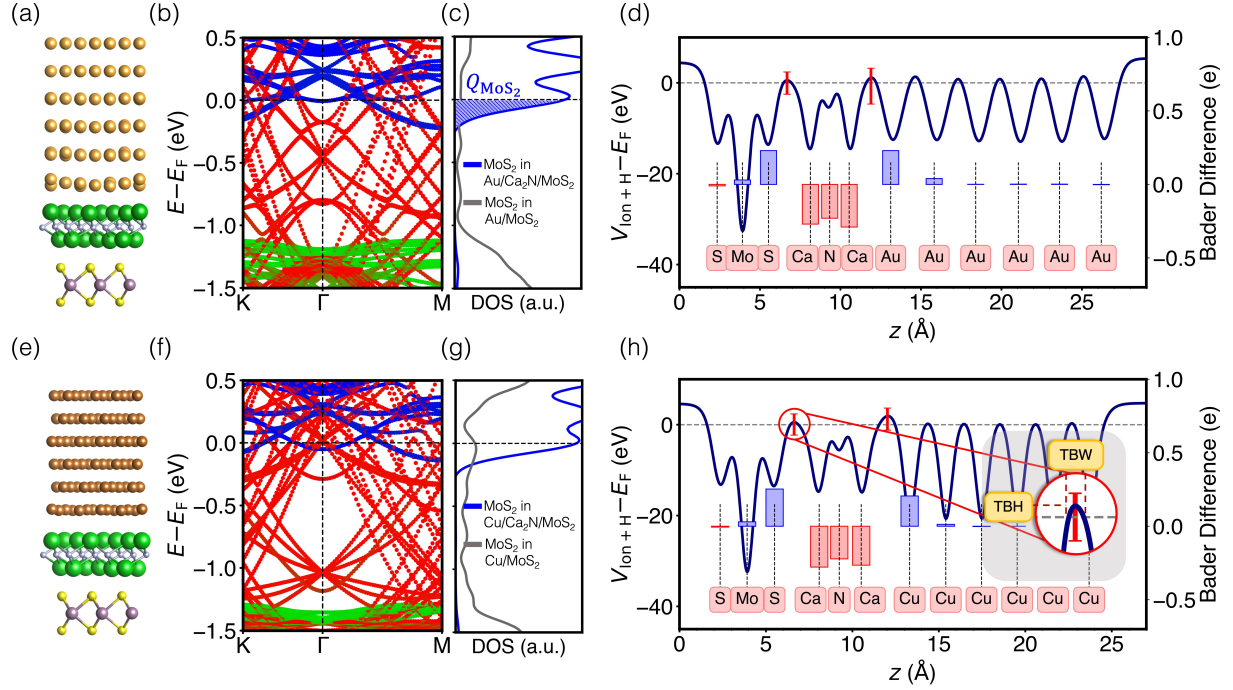


Fig. S1: (a) Optimized geometry, (b) material-projected band structure, (c) densities of MoS₂ states per formula unit, and (d) electrostatic potential for the Au/MoS₂ interface. Panels (e)-(h) present the same quantities for the Cu/MoS₂ interface. In the potential plots, the red bars indicate the range of values at the maximum point within the interface region. The results are overlaid with the differences in Bader atomic charges for the heterostructures, relative to the separated materials. Positive (negative) values indicate accumulation (depletion) of electron density.

addition to the already present Schottky barrier.

Turning to Cu/MoS₂, the range of potential values at the interface extends below the Fermi level, as shown in Fig. S1(h). Thus, while a Schottky barrier also exists for the Cu/MoS₂ heterostructure, our calculations indicate that electrons can travel through some fraction of the interface without tunnelling. Notably, Helfrecht *et al.*³ showed chemisorption of TMDCs on Cu. Similar to Ti/MoS₂ interfaces,^{5,6} partial covalent bonding characteristics can explain the absence of tunnelling barriers at Cu/MoS₂ contacts. We also predict some chemisorption of MoS₂ on Cu, since Table S1 shows a favourable bonding contribution from the base density functional to the exfoliation energy for the Cu/MoS₂ interface.

Finally, the secondary axes in Fig. S1(d,h) show the Bader charge transfer due to interface formation. For the Au/MoS₂ heterostructure in Fig. S1(d), there is effectively zero charge

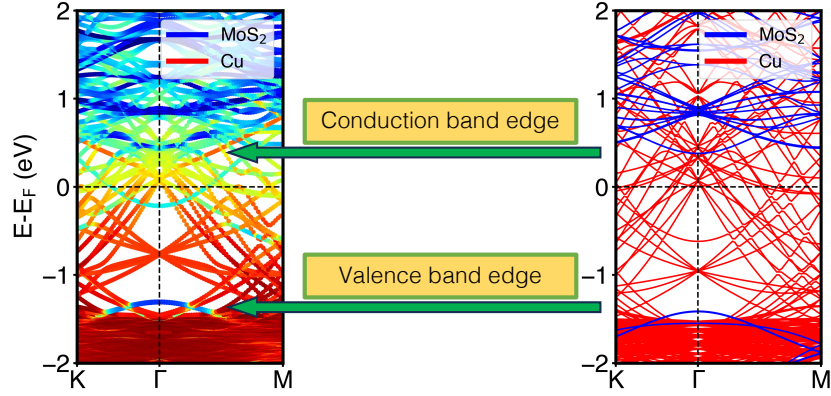


Fig. S2: Determination of the MoS₂ conduction band edge and SBH for the Cu/MoS₂ interface. Left: The electronic band structure of the Cu/MoS₂ interface; the colour scale indicates the contributions from the Cu slab (red) and MoS₂ (blue). Right: Overlaid electronic band structures of the isolated Cu slab (red) and MoS₂ monolayer (blue) using the same (distorted) geometries as for the full Cu/MoS₂ interface.

transfer, as expected for a vdW contact. Conversely, for the Cu/MoS₂ heterostructure in Fig. S1(f), the negative (red) bar for the Cu atoms nearest the interface indicates depletion of charge from the metal, while the positive (blue) bars illustrate the accumulated charge within the MoS₂. This charge transfer is also an indicator of chemisorption for the Cu/MoS₂ interface.

2 Hybrid Functional Tests on $\text{Ca}_2\text{N}/\text{MoS}_2$ Contacts

To investigate the effect of delocalization error on our results, calculations were performed for the $\text{Ca}_2\text{N}/\text{MoS}_2$ interface, using the same interface structure as in Ref. 7. The atomic positions were then optimized using the PBE and HSE06 functionals and the XDM dispersion correction⁸ as implemented in the FHI-aims program.⁹ The calculations additionally used ‘tier1’ or ‘light’ basis sets, ‘tier2’ or ‘tight’ integration grids, the ZORA scalar relativity correction,¹⁰ and a $9 \times 9 \times 1$ \mathbf{k} -point mesh. Figure S3 shows the band structures obtained with both functionals at the corresponding optimized geometries. The two band structures are quite similar overall, although the gap between the Ca_2N nitrogen states and surfaces states is larger with HSE06, as expected. The MoS_2 conduction band also reaches farther below the Fermi level with HSE06, indicating greater charge transfer from Ca_2N to MoS_2 .

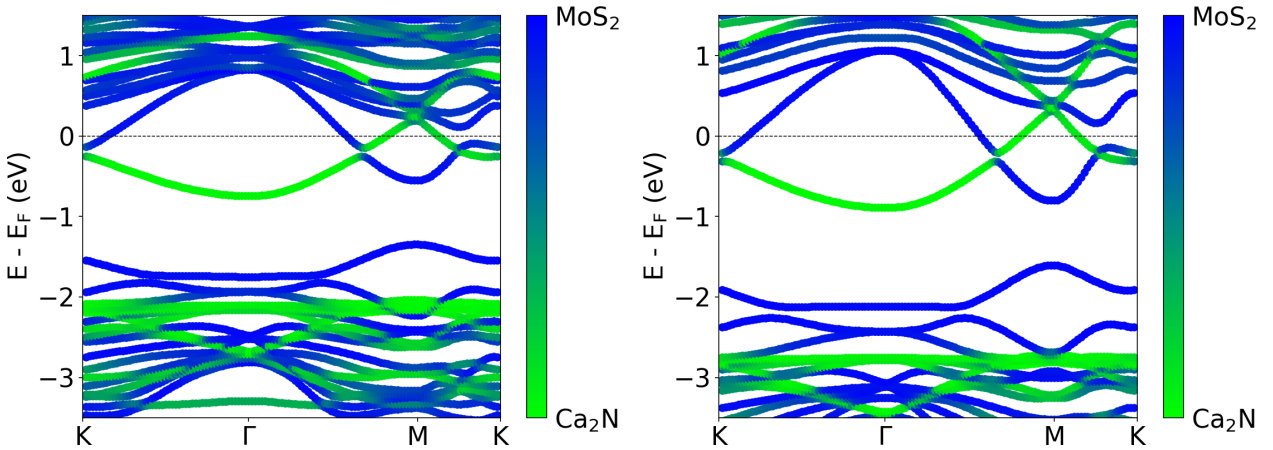


Fig. S 3: Computed materials-projected band structures for the $\text{Ca}_2\text{N}/\text{MoS}_2$ interface, obtained with the PBE (left) and HSE06 (right) density functionals using the FHI-aims program.

3 Metal/Electrene/MoS₂ Contacts

3.1 Construction of Interfaces

TableS 2: Number of atoms, \mathbf{k} -point mesh, in-plane lattice constants (a), strain on the electrene (ϵ_e), strain on the metal slab (ϵ_{Au}), the interlayer distance between the MoS₂ and electrene layer (\bar{d}_{MoS_2}) and the interlayer distance between the electrene and metal slab (\bar{d}_{metal}) for the metal/electrene/MoS₂ interfaces.

Au/electrene/MoS ₂							
Electrene	Ca ₂ N	Sr ₂ N	Sr ₂ P	Ba ₂ N	Ba ₂ P	Ba ₂ As	Ba ₂ Sb
# atoms	120	162	156	243	234	111	111
\mathbf{k} -point mesh	$4 \times 4 \times 1$	$3 \times 3 \times 1$	$3 \times 3 \times 1$	$3 \times 3 \times 1$	$3 \times 3 \times 1$	$4 \times 4 \times 1$	$4 \times 4 \times 1$
$a (= b)$ (Å)	9.47	13.75	11.38	13.75	11.38	9.47	9.47
ϵ_e (%)	0.16	0.34	2.34	0.39	0.28	1.05	4.28
ϵ_{Au} (%)	3.26	2.57	0.79	2.57	2.57	3.26	3.26
\bar{d}_{MoS_2} (Å)	2.59	2.70	2.71	2.94	2.86	2.93	2.83
\bar{d}_{Au} (Å)	2.68	2.85	2.83	2.95	2.92	2.81	2.81

Cu/electrene/MoS ₂							
Electrene	Ca ₂ N	Sr ₂ N	Sr ₂ P	Ba ₂ N	Ba ₂ P	Ba ₂ As	Ba ₂ Sb
# atoms	126	180	174	261	117	117	219
\mathbf{k} -point mesh	$4 \times 4 \times 1$	$3 \times 3 \times 1$	$3 \times 3 \times 1$	$3 \times 3 \times 1$	$4 \times 4 \times 1$	$4 \times 4 \times 1$	$3 \times 3 \times 1$
$a (= b)$ (Å)	9.47	11.38	11.38	13.75	9.47	9.47	12.62
ϵ_e (%)	0.16	0.93	2.34	0.39	3.41	1.05	3.46
ϵ_{Cu} (%)	2.46	1.87	1.87	1.46	2.45	2.45	1.45
\bar{d}_{MoS_2} (Å)	2.57	2.68	2.65	2.90	2.88	2.88	2.83
\bar{d}_{Cu} (Å)	2.78	2.85	2.82	2.99	2.93	2.95	2.97

3.2 Strain Energies

Table S 3: Percent strains on the metal slab ($\epsilon_{\text{Au/Cu}}$) and the electrene layer (ϵ_e) in metal/electrene/MoS₂ heterostructures, along with the component and total strain energies in meV/Å² using the area of the strained unit cell.

Au/electrene/MoS ₂					
Electrene	ϵ_{Au} (%)	ϵ_e (%)	$E_{\text{strain,Au}}$	$E_{\text{strain,e}}$	E_{strain}
Ca ₂ N	3.26	0.16	49.1	0.6	49.7
Sr ₂ N	2.57	0.34	46.4	1.3	47.7
Sr ₂ P	0.79	2.34	1.6	1.4	3.0
Ba ₂ N	2.57	0.39	46.5	0.1	46.6
Ba ₂ P	2.57	2.28	46.6	0.1	46.7
Ba ₂ As	3.26	1.05	49.3	0.2	49.5
Ba ₂ Sb	3.26	4.28	49.7	2.8	52.5

Cu/electrene/MoS ₂					
Electrene	ϵ_{Cu} (%)	ϵ_e (%)	$E_{\text{strain,Cu}}$	$E_{\text{strain,e}}$	E_{strain}
Ca ₂ N	2.45	0.16	16.4	0.6	17.0
Sr ₂ N	1.87	0.93	8.1	0.7	8.8
Sr ₂ P	1.87	2.34	22.1	1.4	23.5
Ba ₂ N	1.46	0.39	16.2	0.1	16.3
Ba ₂ P	2.45	3.40	33.7	2.1	35.8
Ba ₂ As	2.45	1.05	33.7	0.2	33.9
Ba ₂ Sb	1.42	3.46	5.8	1.7	7.5

3.3 Bader Charge Transfer

Table S 4: Bader charge transfer values of metal/electrene/MoS₂ interfaces. $Q_{\text{Electrene}}^{\text{Bader}}$ values are the amount of charge transferred from the electrene layer, while $Q_{\text{MoS}_2}^{\text{Bader}}$ and $Q_{\text{Au/Cu}}^{\text{Bader}}$ are charges received by the semiconductor and the metal slabs, respectively. The values are reported in units of electrons per formula unit of the material. Note that the charges do not sum to one because there are different numbers of formula units of each material present in the unit cell.

Au/electrene/MoS ₂											
	$Q_{\text{MoS}_2}^{\text{Bader}} (e^-/\text{f.u.})$				$Q_{\text{Electrene}}^{\text{Bader}} (e^-/\text{f.u.})$				$Q_{\text{Au}}^{\text{Bader}} (e^-/\text{f.u.})$		
	Ca	Sr	Ba		Ca	Sr	Ba		Ca	Sr	Ba
N	0.26	0.25	0.19	N	-0.80	-0.80	-0.75	N	0.05	0.04	0.04
P	-	0.22	0.15	P	-	-0.85	-0.85	P	-	0.04	0.04
As	-	-	0.14	As	-	-	-0.88	As	-	-	0.03
Sb	-	-	0.18	Sb	-	-	-1.00	Sb	-	-	0.04

Cu/electrene/MoS ₂											
	$Q_{\text{MoS}_2}^{\text{Bader}} (e^-/\text{f.u.})$				$Q_{\text{Electrene}}^{\text{Bader}} (e^-/\text{f.u.})$				$Q_{\text{Cu}}^{\text{Bader}} (e^-/\text{f.u.})$		
	Ca	Sr	Ba		Ca	Sr	Ba		Ca	Sr	Ba
N	0.28	0.26	0.23	N	-0.77	-0.77	-0.76	N	0.04	0.03	0.03
P	-	0.22	0.16	P	-	-0.85	-0.79	P	-	0.03	0.02
As	-	-	0.16	As	-	-	-0.80	As	-	-	0.02
Sb	-	-	0.18	Sb	-	-	-0.88	Sb	-	-	0.02

3.4 MoS₂ Band Gap Variation in Metal/Electrene/MoS₂ Contacts

Table S 5: MoS₂ valence and conduction band edges and band gaps (eV) in metal/MoS₂ and metal/electrene/MoS₂ interfaces. For comparison, the band gap of pristine MoS₂ is calculated to be 1.77 eV and, as this is a semiconductor, the Fermi level is taken to lie in the middle of the gap.

Electrene	Au			Cu		
	Valence	Conduction	Gap	Valence	Conduction	Gap
None	-1.27	0.49	1.76	-1.30	0.41*	1.71*
Ca ₂ N	-1.75	-0.14	1.61	-1.84	-0.13	1.71
Sr ₂ N	-1.89	-0.10	1.79	-1.96	-0.13	1.83
Sr ₂ P	-2.02	-0.16	1.86	-1.83	-0.08	1.75
Ba ₂ N	-1.93	-0.08	1.85	-1.95	-0.09	1.86
Ba ₂ P	-1.96	-0.15	1.81	-1.80	-0.09	1.71
Ba ₂ As	-1.87	-0.08	1.79	-1.81	-0.09	1.72
Ba ₂ Sb	-1.85	-0.11	1.74	-1.74	-0.08	1.66

*These values are estimates only due to the high degree of state mixing near the conduction band edge.

3.5 Effect of Sliding the MoS₂ Layer

Table S 6: Changes in various properties of the Cu/Ca₂N/MoS₂ interface as a function of differing initial geometries, with the MoS₂ translated in the a, b lattice directions by a particular fraction of these cell lengths before relaxation. Shown are: the relative energy, ΔE ; the Fermi energy, E_{Fermi} ; the charge transfer to MoS₂, Q_{MoS_2} ; the fraction of the contact area that is Ohmic, F_{Ohmic} ; the tunneling barrier height, TBH; and tunneling barrier width, TBW.

Sliding Fraction	ΔE (meV/cell)	E_{Fermi} (eV)	Q_{MoS_2} ($10^{14}e^-/\text{cm}^2$)	F_{Ohmic} (%)	TBH (eV)	TBW (Å)
0.0	0.0	4.61	3.00	32.9	0.496	0.46
0.2	20.9	4.62	2.95	30.9	0.546	0.48
0.4	11.4	4.62	2.94	31.5	0.539	0.48
0.6	18.4	4.62	2.92	31.2	0.547	0.48
0.8	21.7	4.62	2.94	30.7	0.554	0.49

3.6 Summary of Interface Properties

Table S 7: Figures of merit computed for various the metal/electrene/MoS₂ contacts: the charge transfer to MoS₂, Q_{MoS_2} ; the fraction of the contact area that is Ohmic, F_{Ohmic} ; the tunneling barrier height, TBH; and tunneling barrier width, TBW. Results for the metal/MoS₂ interfaces with no electrene are shown for comparison; note that the Q_{MoS_2} for these two cases results from mid-gap states rather than from the MoS₂ conduction band edge.

Interface	Q_{MoS_2} ($10^{14}e^-/\text{cm}^2$)	F_{Ohmic} (%)	TBH (eV)	TBW (Å)
Cu/MoS ₂	1.92	30.8	0.72	0.40
Cu/Ca ₂ N/MoS ₂	3.09	32.5	0.54	0.47
Cu/Sr ₂ N/MoS ₂	2.98	27.3	0.72	0.55
Cu/Sr ₂ P/MoS ₂	1.91	21.7	1.03	0.68
Cu/Ba ₂ N/MoS ₂	2.32	17.8	1.07	0.71
Cu/Ba ₂ P/MoS ₂	1.49	10.1	1.37	0.88
Cu/Ba ₂ As/MoS ₂	1.57	12.3	1.41	0.89
Cu/Ba ₂ Sb/MoS ₂	1.55	13.8	1.37	0.85
Au/MoS ₂	0.77	0.0	3.01	1.08
Au/Ca ₂ N/MoS ₂	2.28	31.7	0.56	0.49
Au/Sr ₂ N/MoS ₂	2.23	26.4	0.70	0.70
Au/Sr ₂ P/MoS ₂	1.74	19.7	1.05	0.74
Au/Ba ₂ N/MoS ₂	1.88	10.1	1.21	0.78
Au/Ba ₂ P/MoS ₂	1.72	8.8	1.54	0.95
Au/Ba ₂ As/MoS ₂	1.44	9.8	1.62	1.00
Au/Ba ₂ Sb/MoS ₂	1.34	13.6	1.38	0.85

3.7 Electronic Properties of Au/Electrene/MoS₂ Contacts

3.7.1 Au/Ca₂N/MoS₂

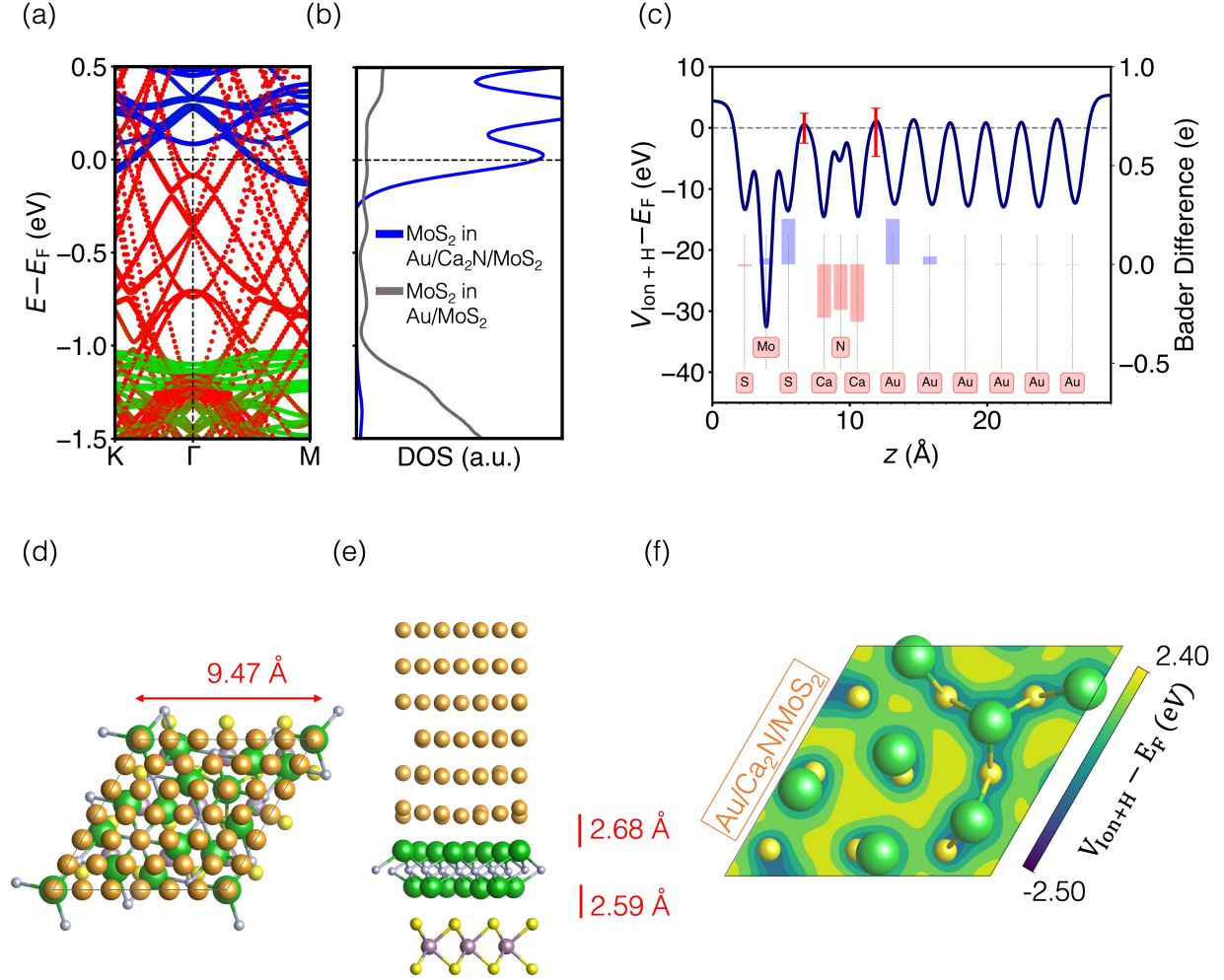


Fig. S4: Properties of the Au/Ca₂N/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.7.2 Au/Sr₂N/MoS₂

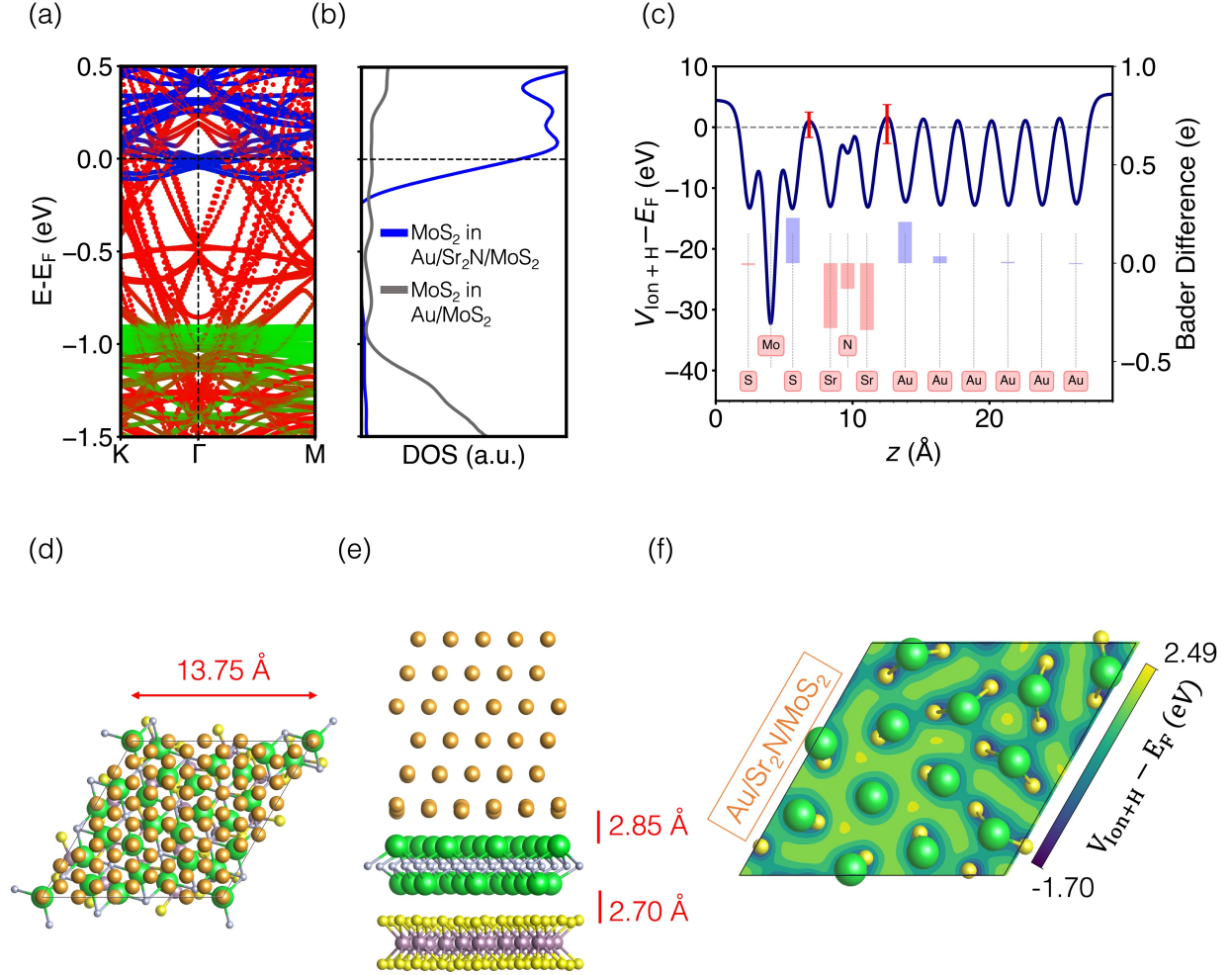


Fig. S5: Properties of the Au/Sr₂N/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.7.3 Au/Sr₂P/MoS₂

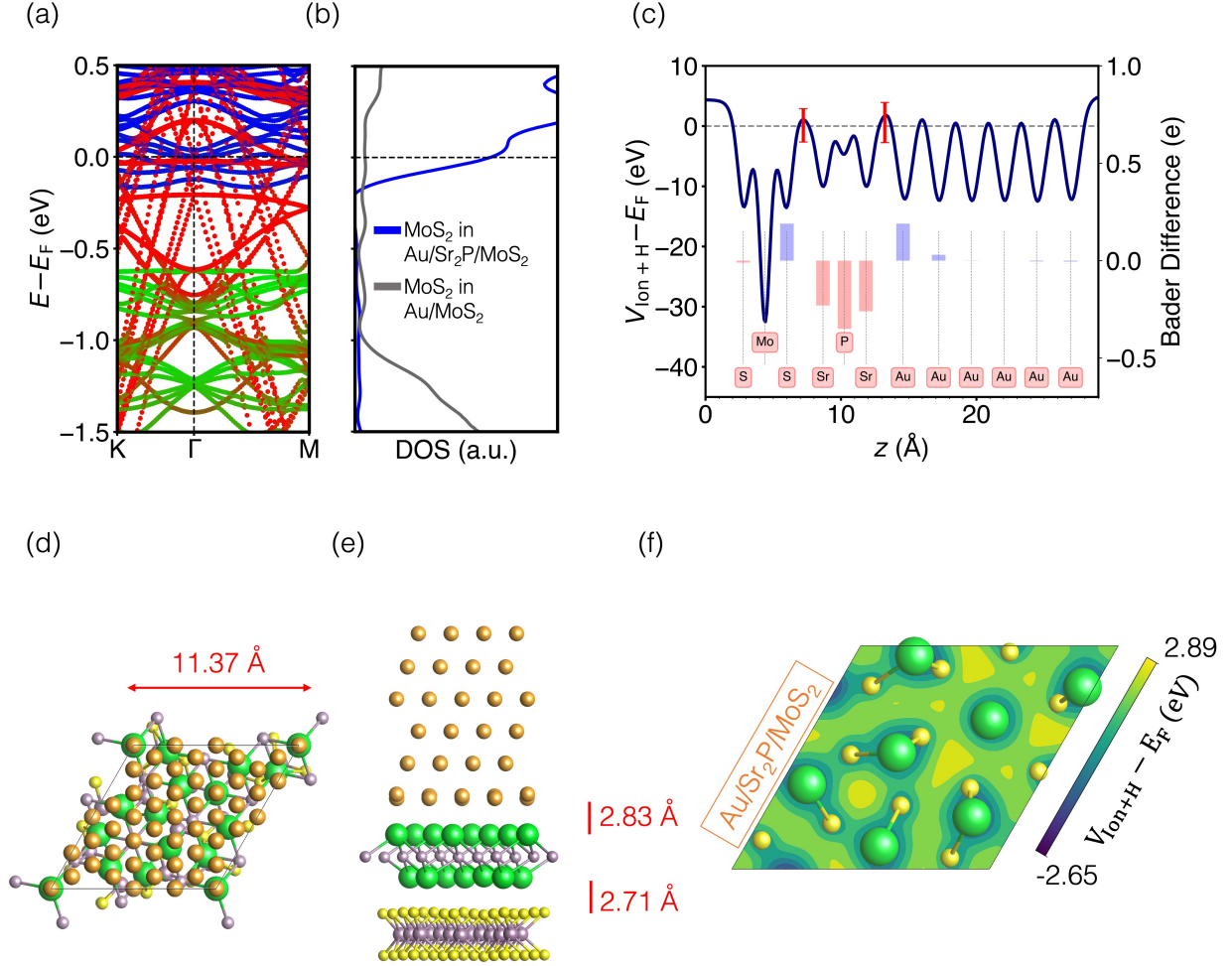


Fig. S6: Properties of the Au/Sr₂P/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.7.4 Au/Ba₂N/MoS₂

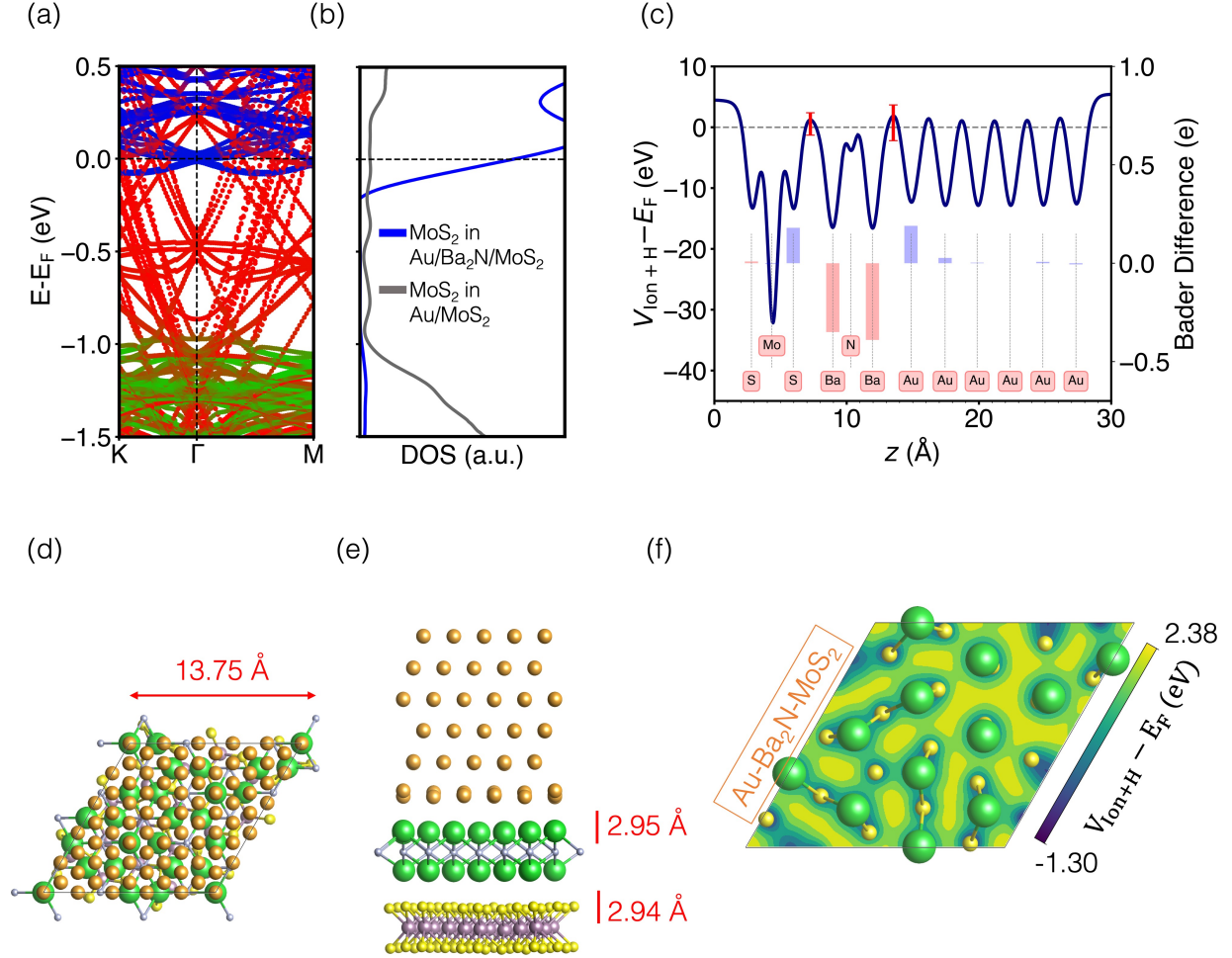


Fig. S7: Properties of the Au/Ba₂N/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.7.5 Au/Ba₂P/MoS₂

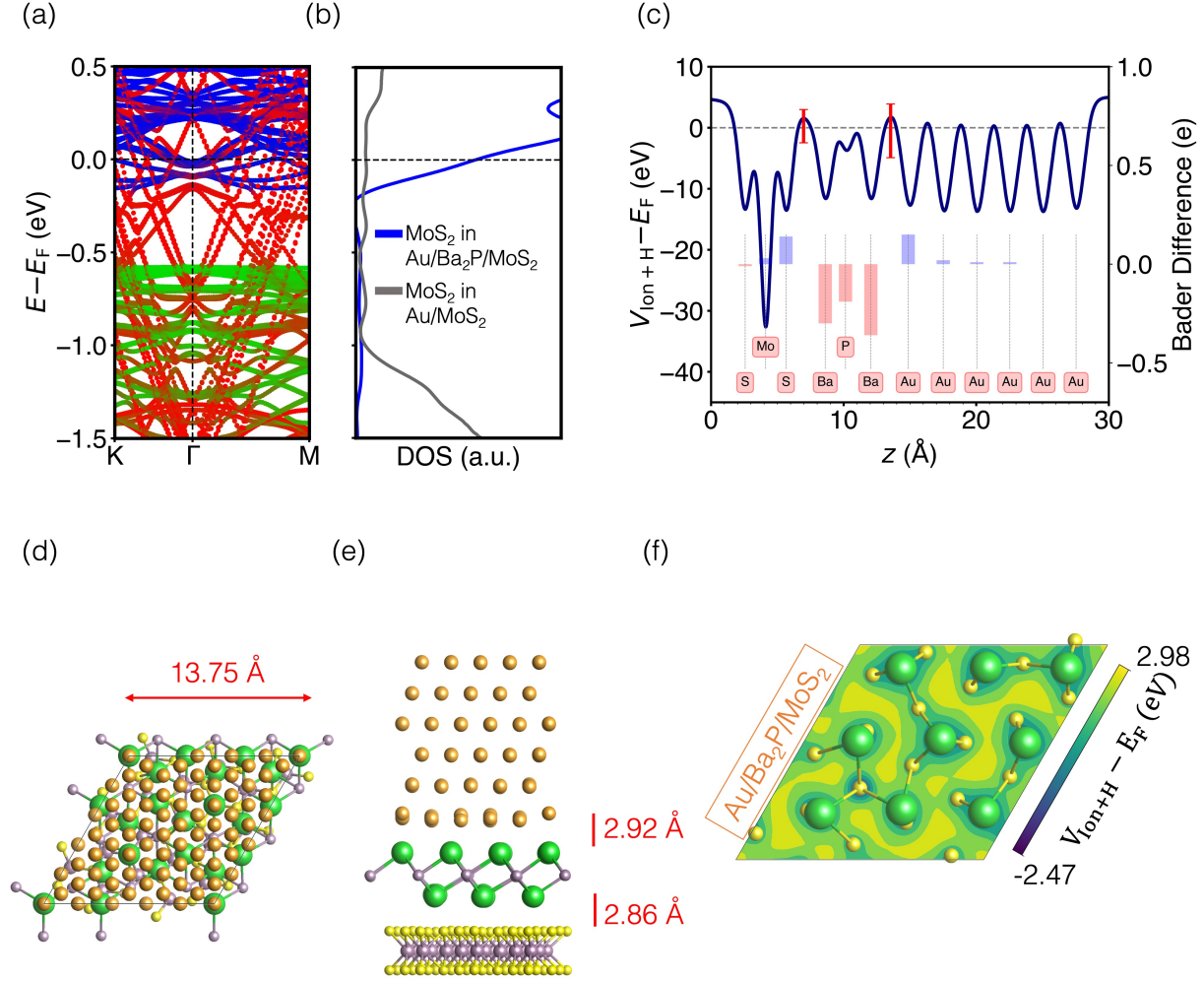


Fig. S8: Properties of the Au/Ba₂P/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.7.6 Au/Ba₂As/MoS₂

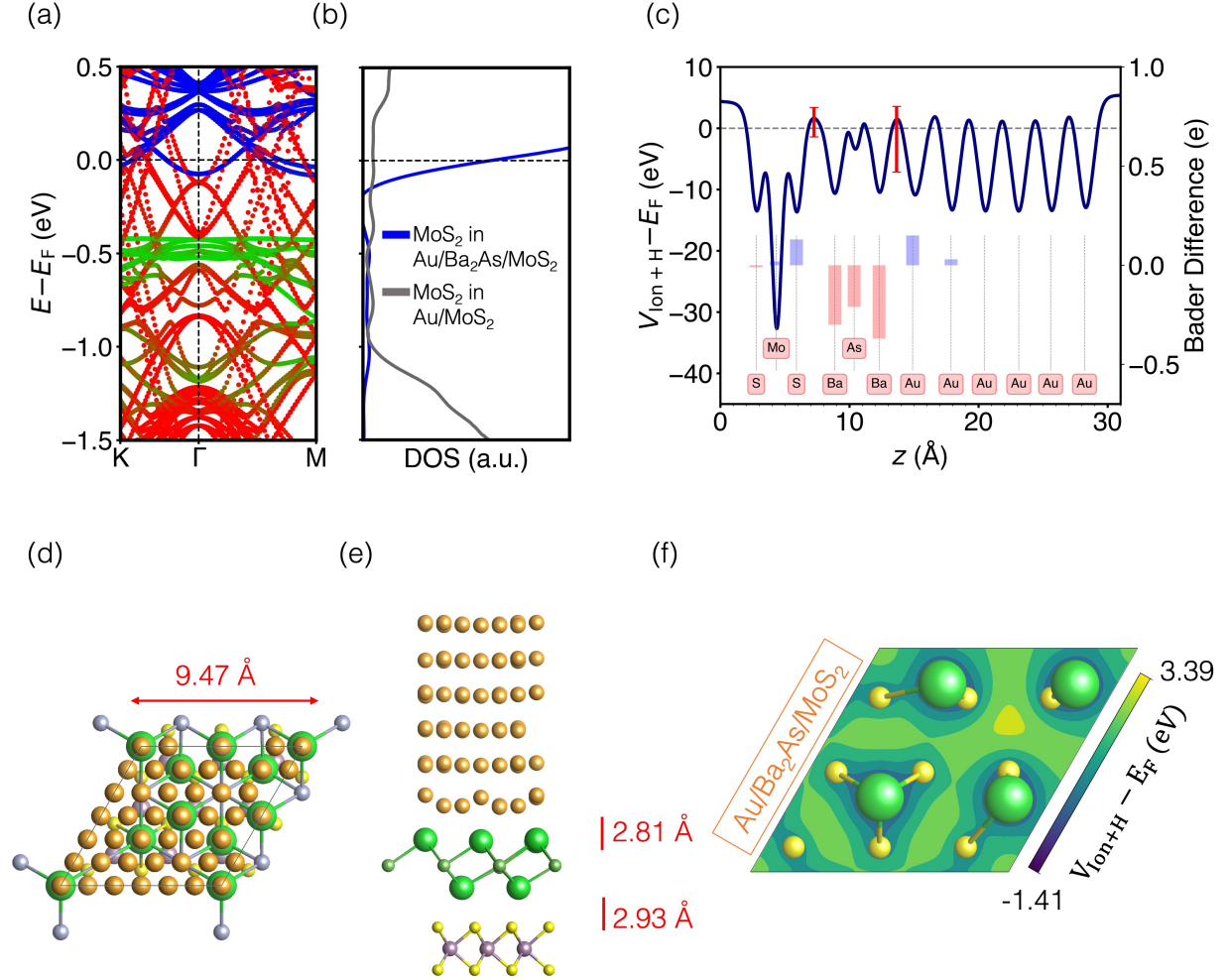


Fig. S 9: Properties of the Au/Ba₂As/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.7.7 Au/Ba₂Sb/MoS₂

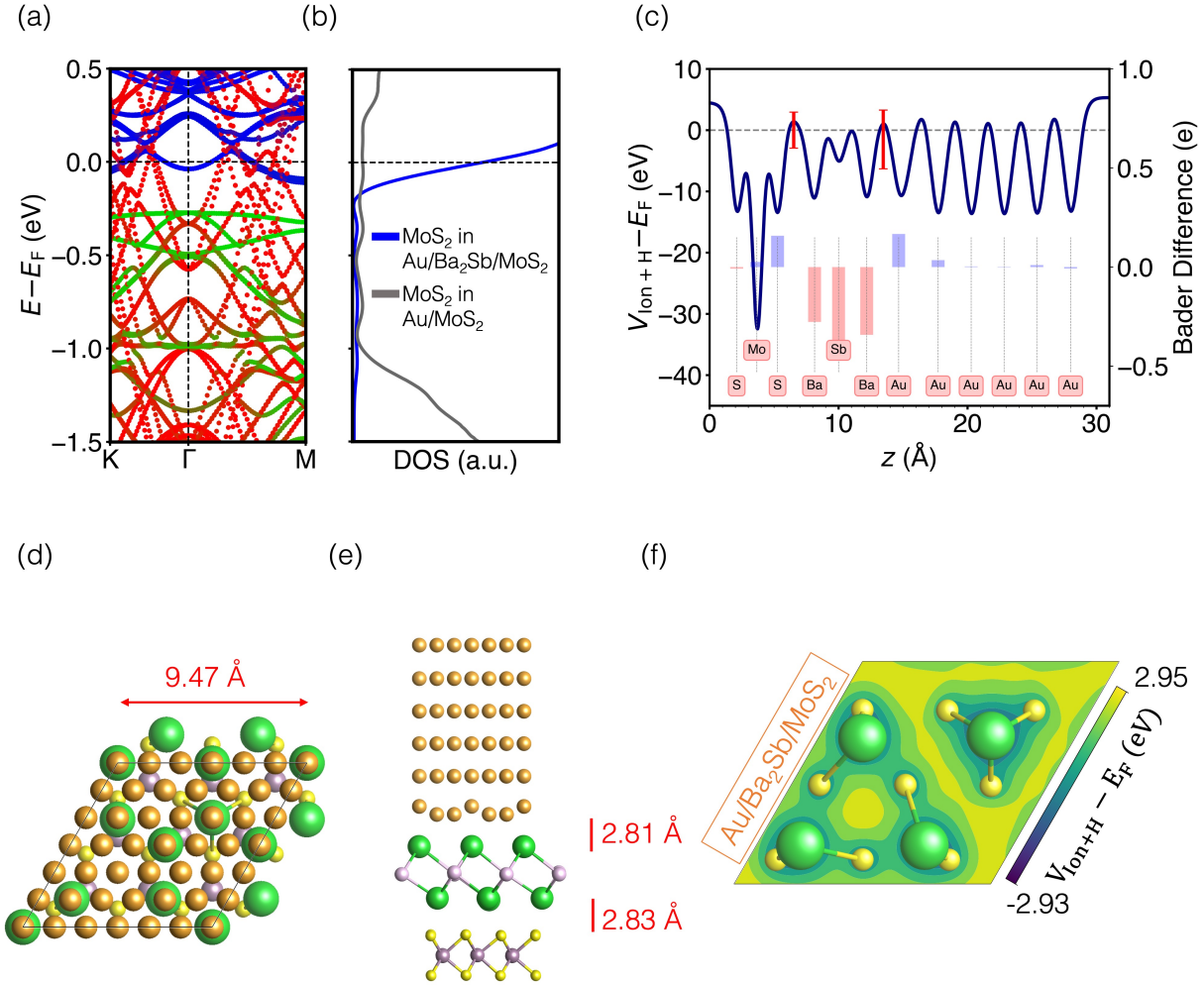


Fig. S 10: Properties of the Au/Ba₂Sb/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8 Electronic Properties of Cu/electrene/MoS₂ Contacts

3.8.1 Cu/Ca₂N/MoS₂

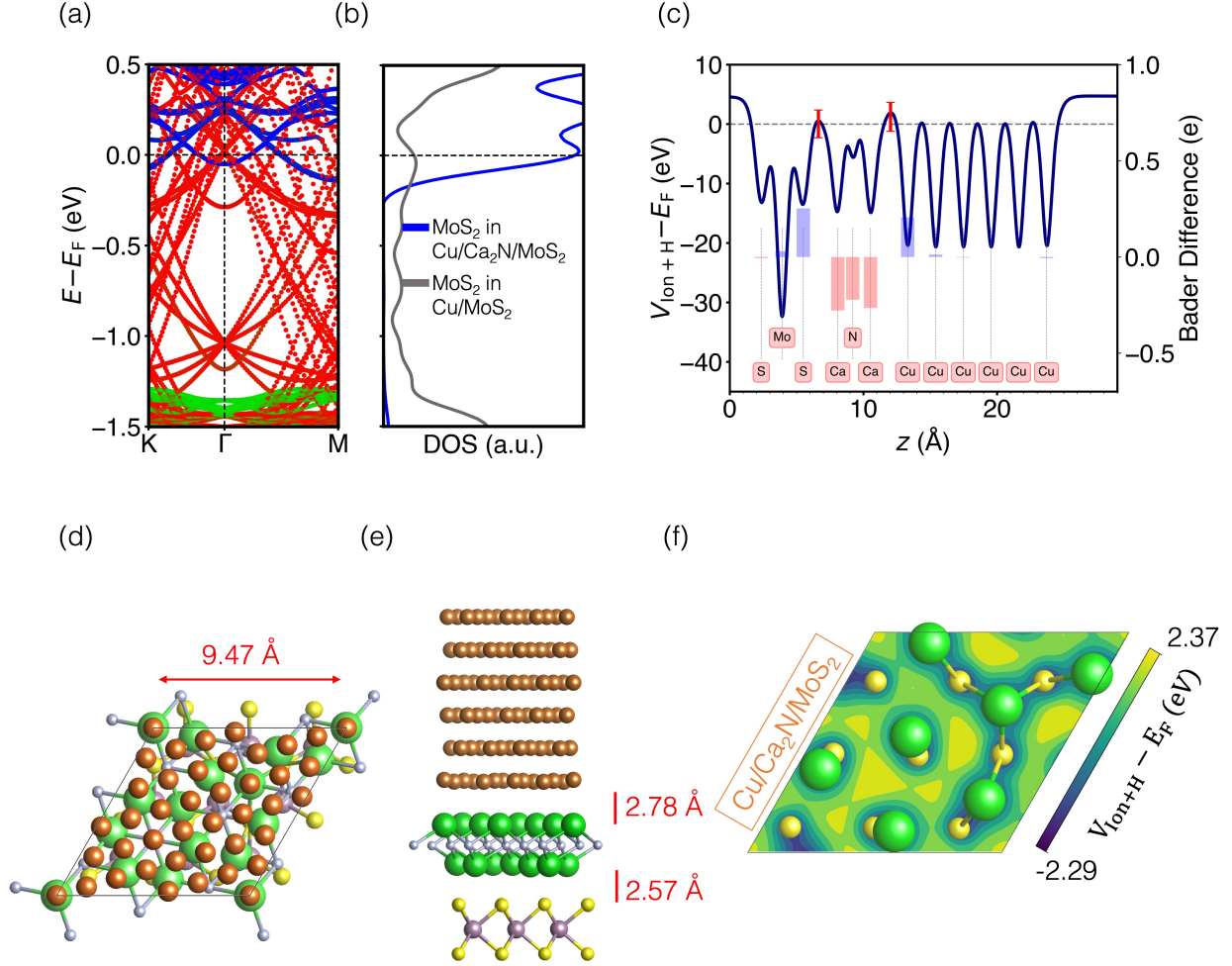


Fig. S 11: Properties of the Cu/Ca₂N/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8.2 Cu/Sr₂N/MoS₂

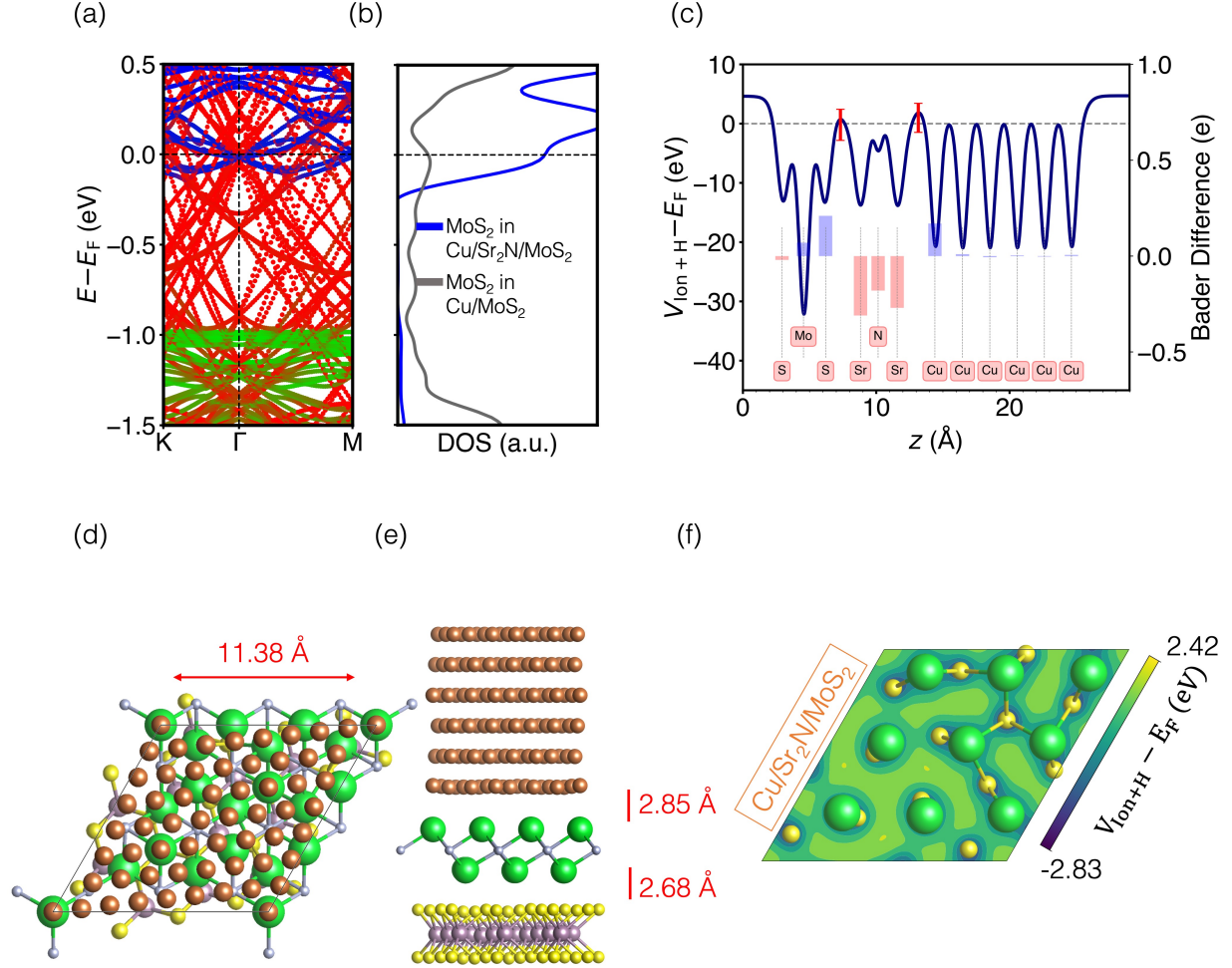


Fig. S12: Properties of the Cu/Sr₂N/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8.3 Cu/Sr₂P/MoS₂

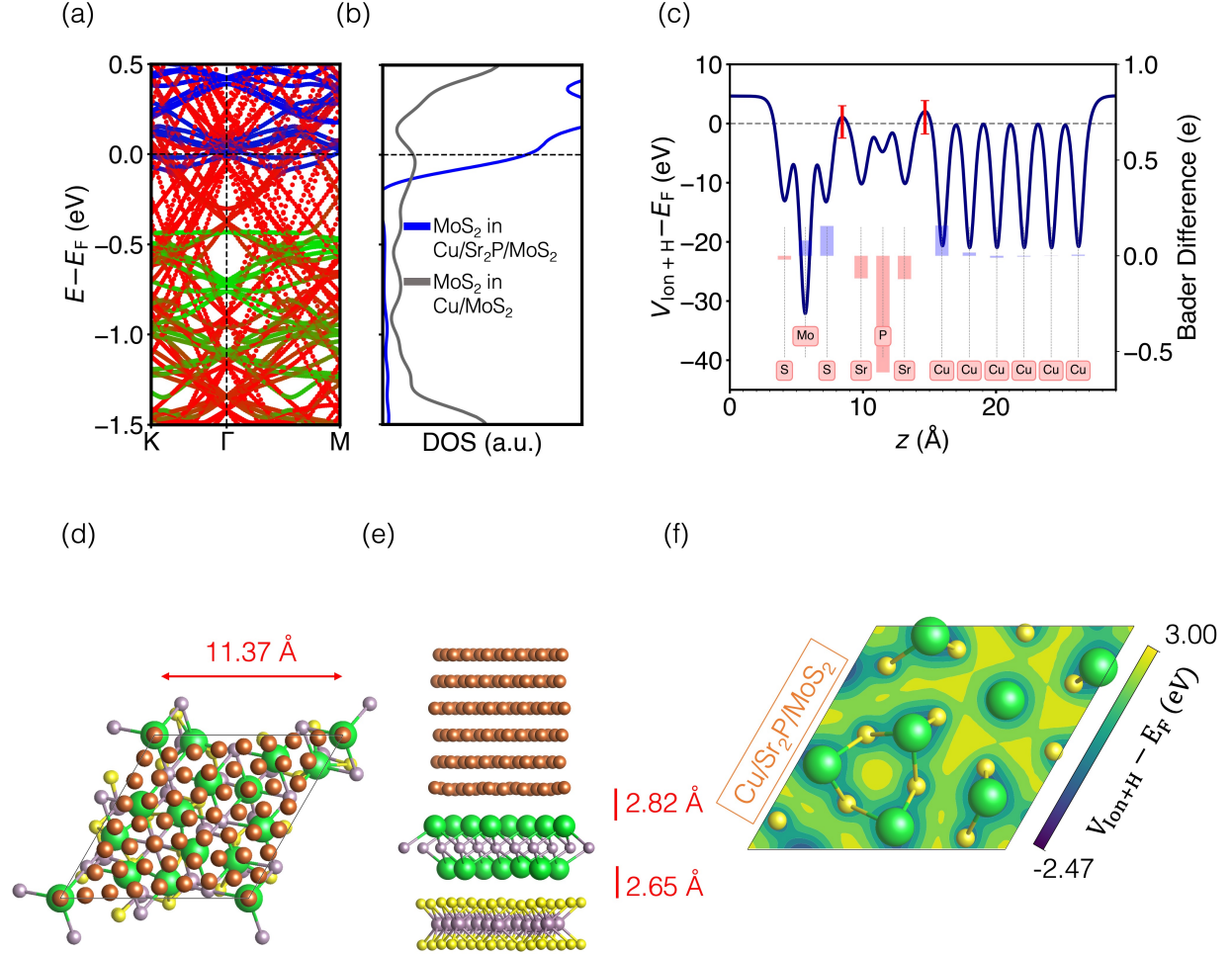


Fig. S13: Properties of the Cu/Sr₂P/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8.4 Cu/Ba₂N/MoS₂

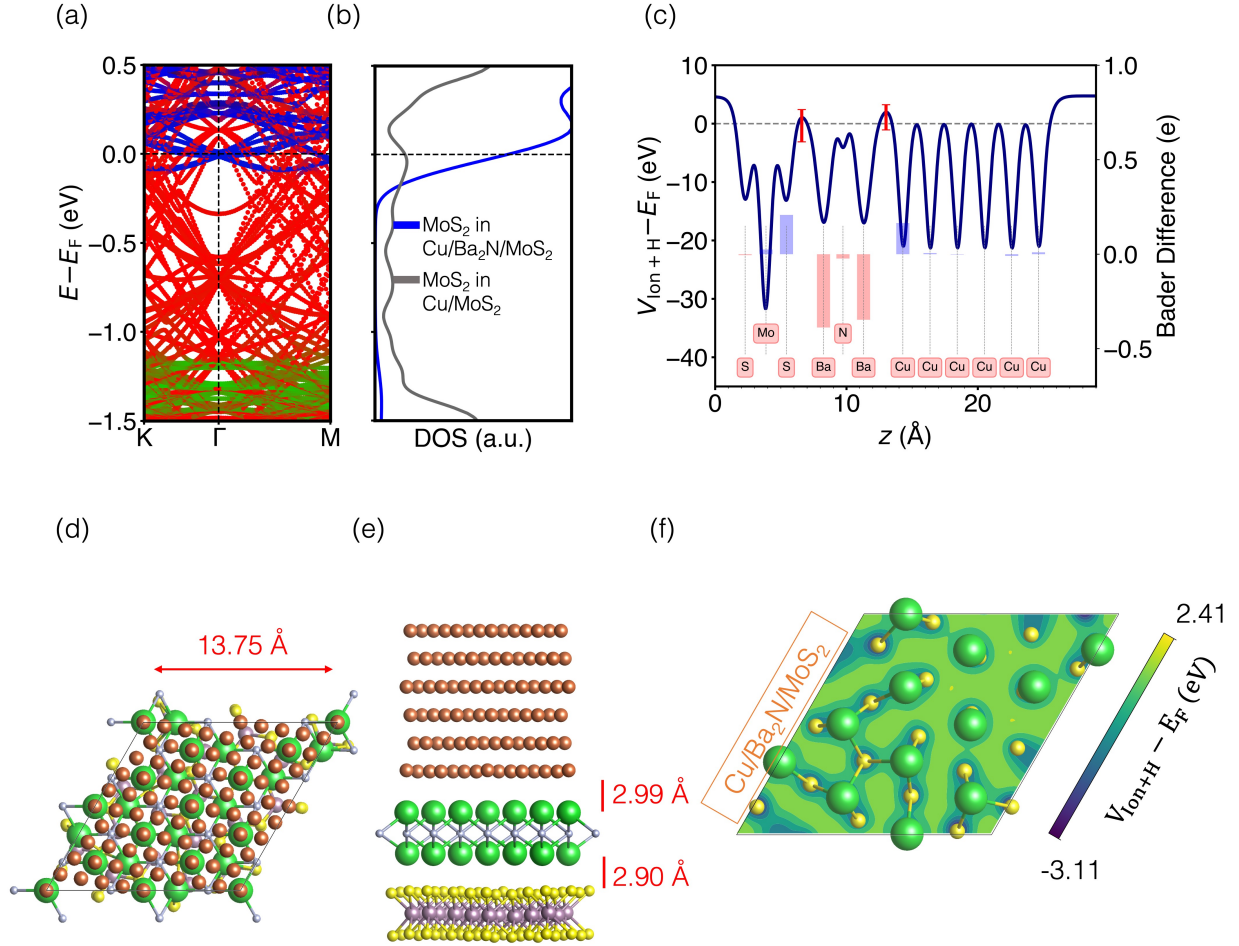


Fig. S 14: Properties of the Cu/Ba₂N/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8.5 Cu/Ba₂P/MoS₂

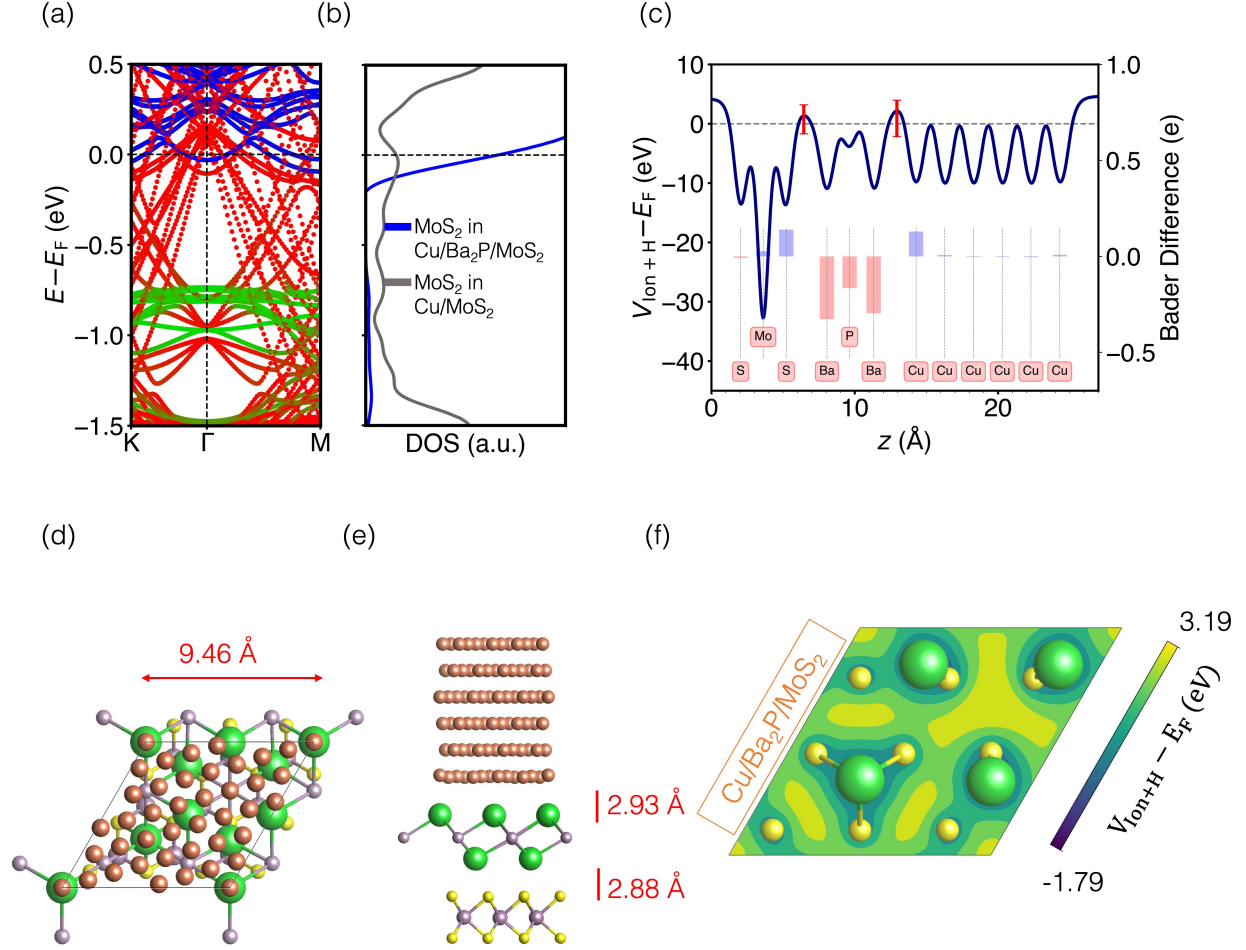


Fig. S 15: Properties of the Cu/Ba₂P/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8.6 Cu/Ba₂As/MoS₂

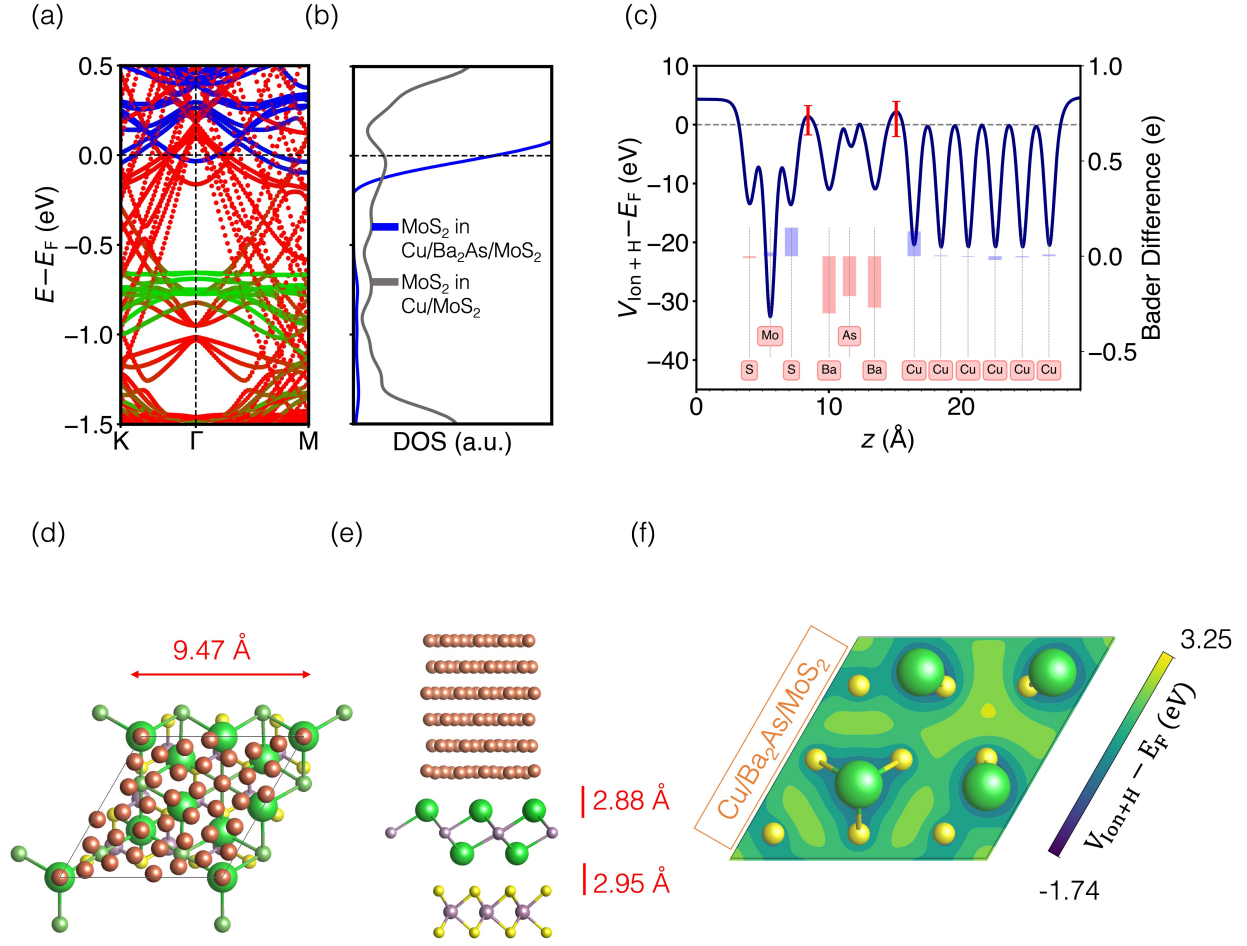


Fig. S 16: Properties of the Cu/Ba₂As/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.8.7 Cu/Ba₂Sb/MoS₂

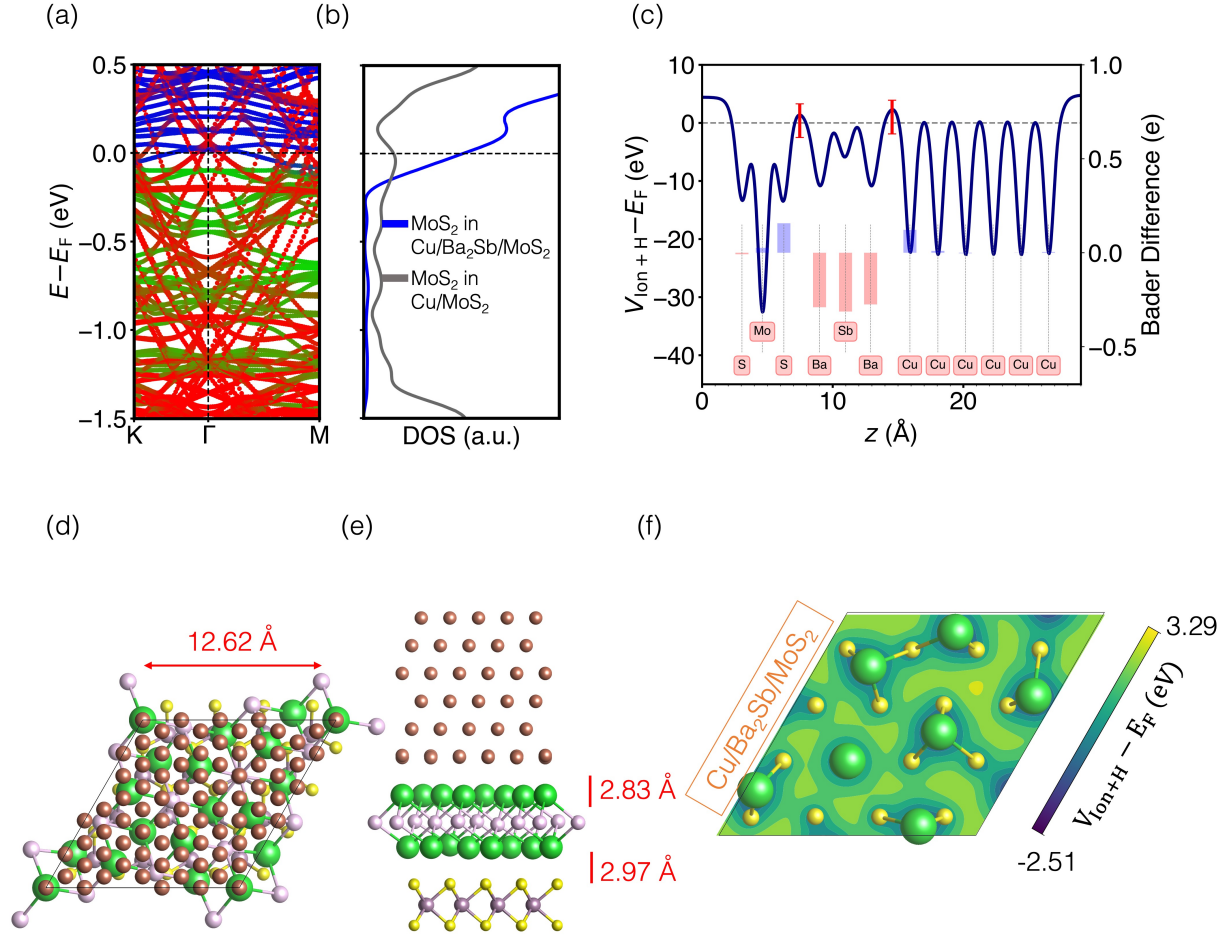


Fig. S 17: Properties of the Cu/Ba₂Sb/MoS₂ heterojunction: (a) materials-projected electronic band structure, (b) densities of MoS₂ states per formula unit, (c) averaged electrostatic potential profile, (d) and (e) top and side views of the interface, and (f) 2D projection of the electrostatic potential at the highest potential point of the electrene/MoS₂ interface.

3.9 Electronic Structure of Isolated Individual Layers

3.9.1 Band Structures of Distorted Layers from Au/Electrene/MoS₂ Interfaces

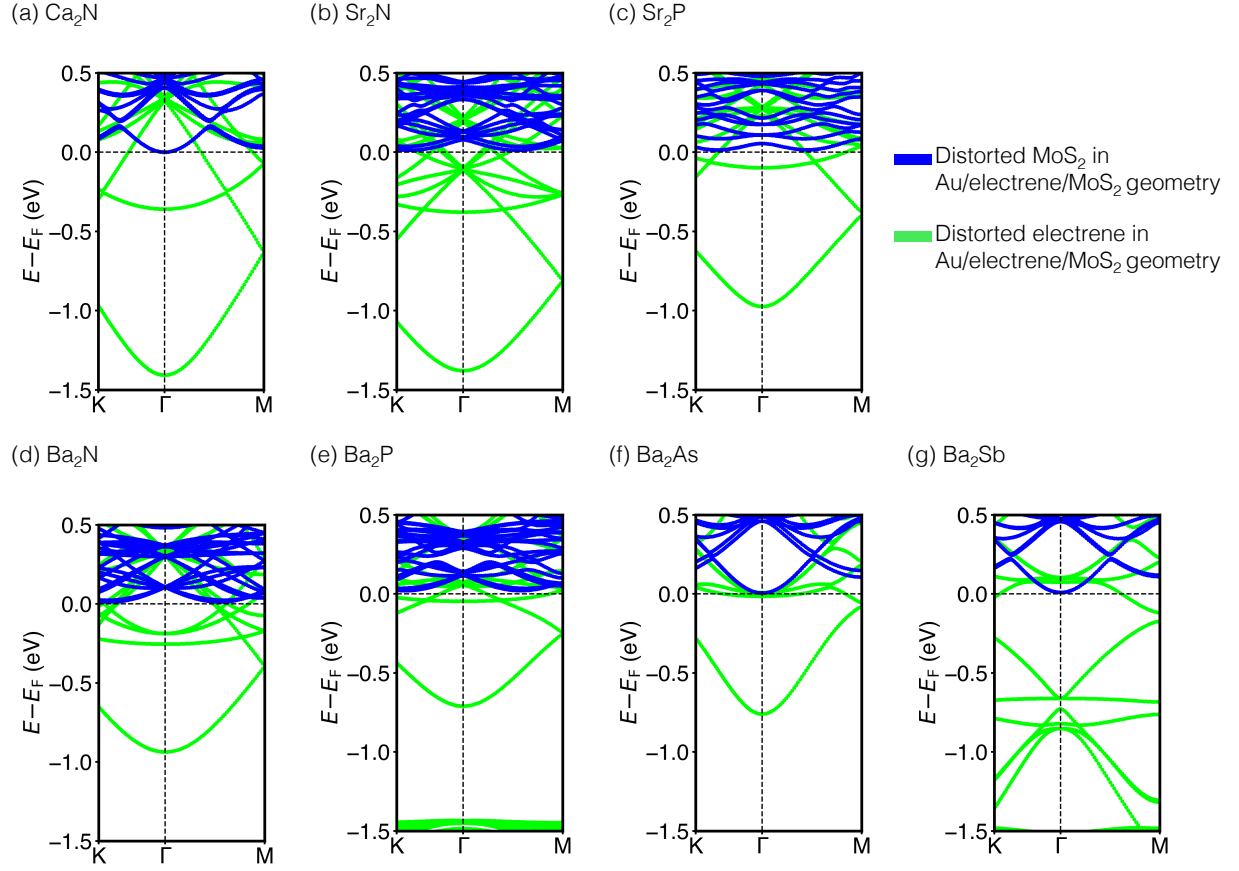


Fig. S 18: Band structures of isolated Ca_2N and MoS_2 layers at rigid, distorted geometries taken from the Au/electrene/ MoS_2 interfaces.

3.9.2 Band Structures of Relaxed Layers from Au/Electrene/MoS₂ Interfaces

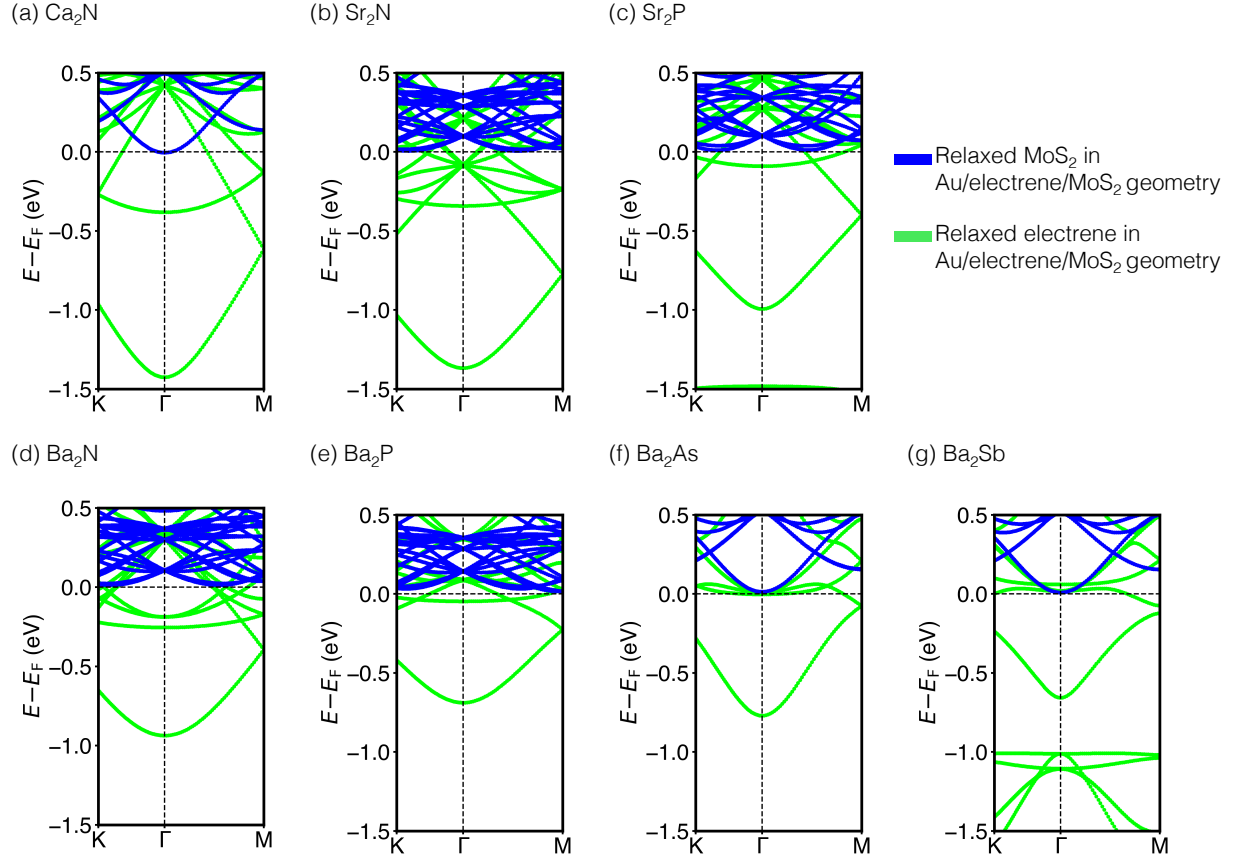


Fig. S19: Band structures of isolated Ca_2N and MoS_2 layers at relaxed geometries with fixed latticed constants matching those of the Au/electrene/ MoS_2 interfaces.

3.9.3 Band Structures of Distorted Layers from Cu/Electrene/MoS₂ Interfaces

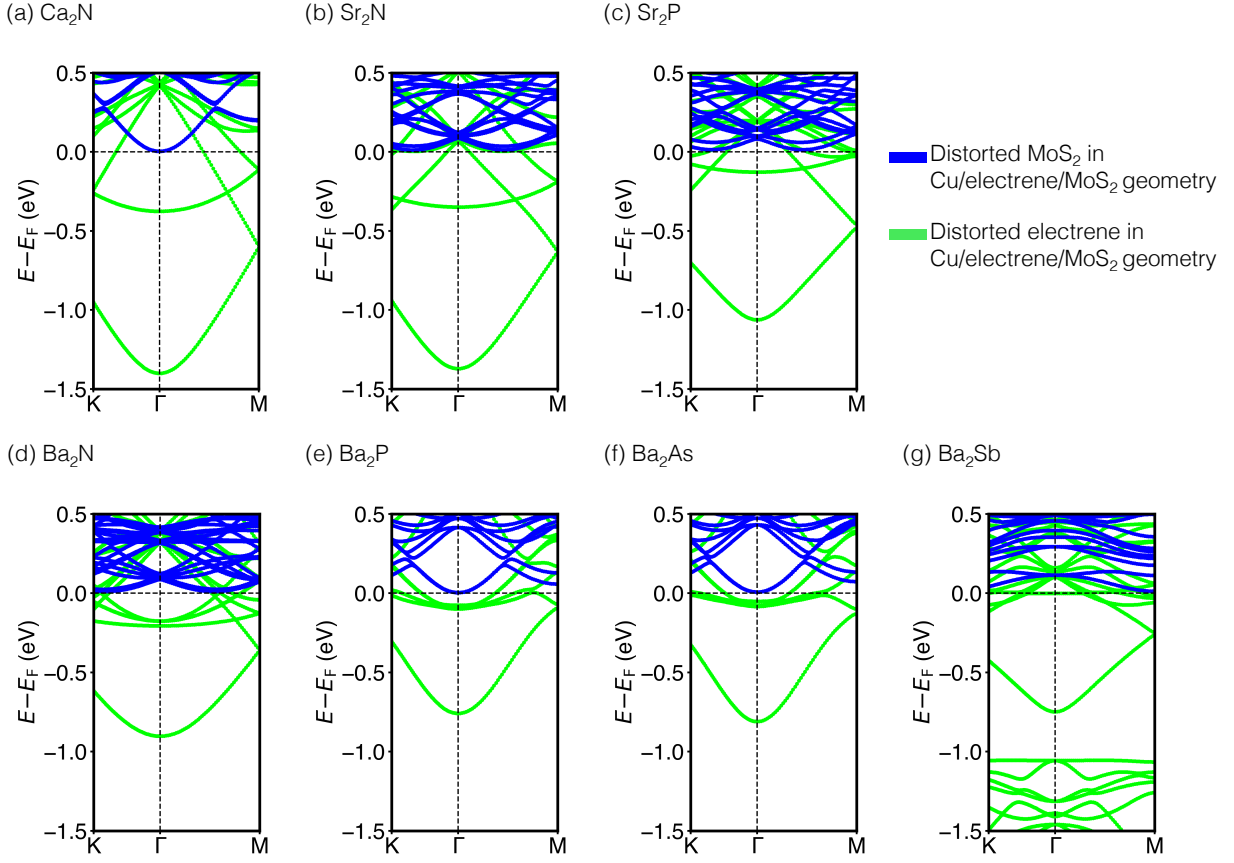


Fig. S 20: Band structures of isolated Ca_2N and MoS_2 layers at rigid, distorted geometries taken from the Cu/electrene/ MoS_2 interfaces.

3.9.4 Band Structures of Relaxed Layers from Cu/Electrene/MoS₂ Interfaces

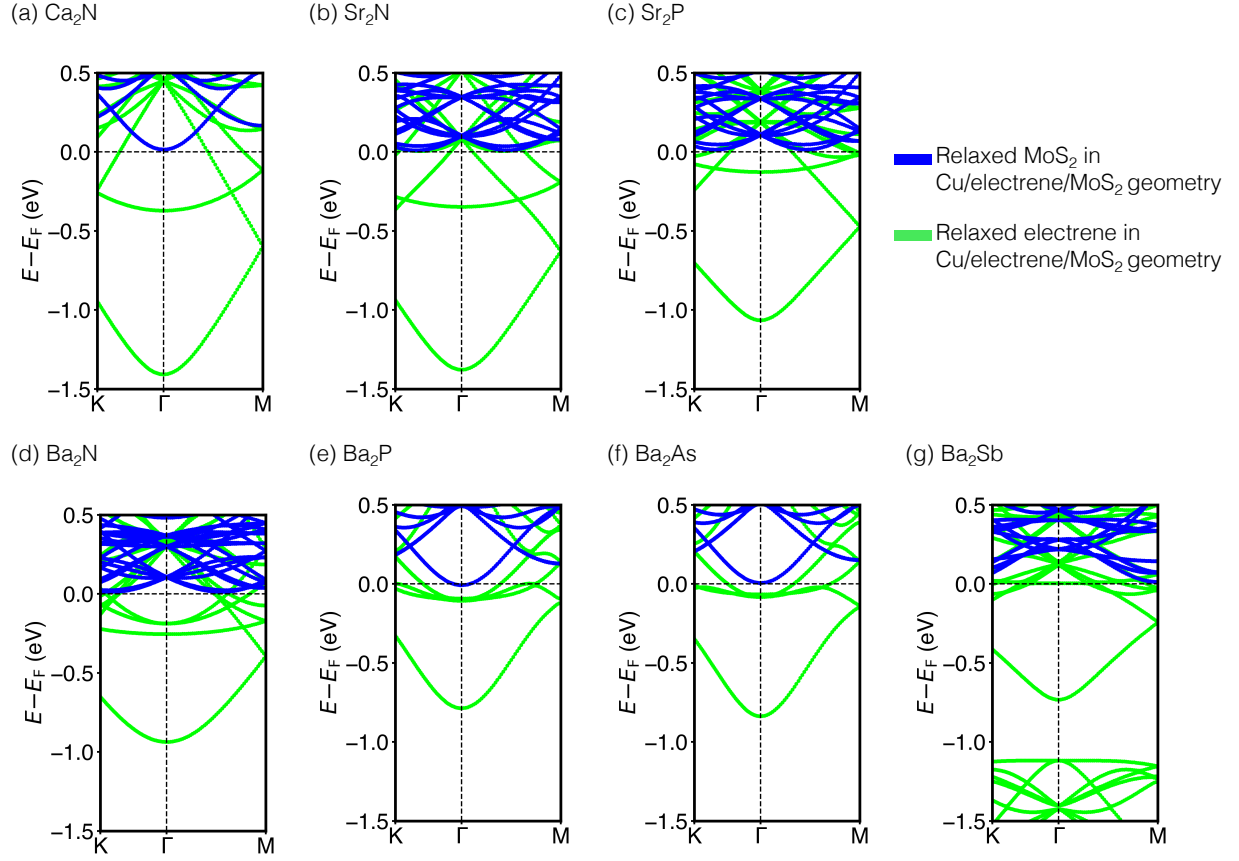


Fig. S21: Band structures of isolated Ca_2N and MoS_2 layers at relaxed geometries with fixed lattice constants matching those of the Cu/electrene/ MoS_2 interfaces.

4 Tunnelling Characteristics from the WKB Approximation

Here we discuss the tunnelling probability and tunnelling-specific resistivity based on the WKB approximation.^{11,12} Note that the WKB approximation is heavily reliant on the shape of the potential barrier. In its general form in one dimension, the transmission probability, \mathcal{T}_{WKB} , can be written in terms of the relative amplitudes of the transmitted (F) and incident (A) waves,

$$\mathcal{T}_{\text{WKB}} = \frac{|F|^2}{|A|^2} \approx e^{-2\gamma},$$

which can be approximated in terms of the γ parameter obtained by integrating the area under the potential curve, $V(x)$:

$$\gamma = \frac{\sqrt{2m}}{\hbar} \int_{x_1}^{x_2} \sqrt{V(x) - E} dx,$$

where $x_{1,2}$ are the classical turning points at which $E = V(x_{1,2})$.

Assuming a square-shaped barrier, the tunnelling probability can be evaluated in terms of the tunneling barrier height (TBH) and tunneling barrier width (TBW) as

$$\mathcal{T}_{\text{WKB}} = \exp\left(-\frac{4\pi}{h} \text{TBW} \sqrt{2m_e \text{TBH}}\right),$$

while the tunneling resistivity, ρ_t , is

$$\rho_t = \frac{2h^2 \text{TBW}}{3e^2 \sqrt{2m_e} \text{TBH}} \exp\left(\frac{4\pi}{h} \text{TBW} \sqrt{2m_e \text{TBH}}\right).$$

The computed values of \mathcal{T}_{WKB} and ρ_t for the various metal/electrene/MoS₂ heterojunctions are shown in Figure S22.

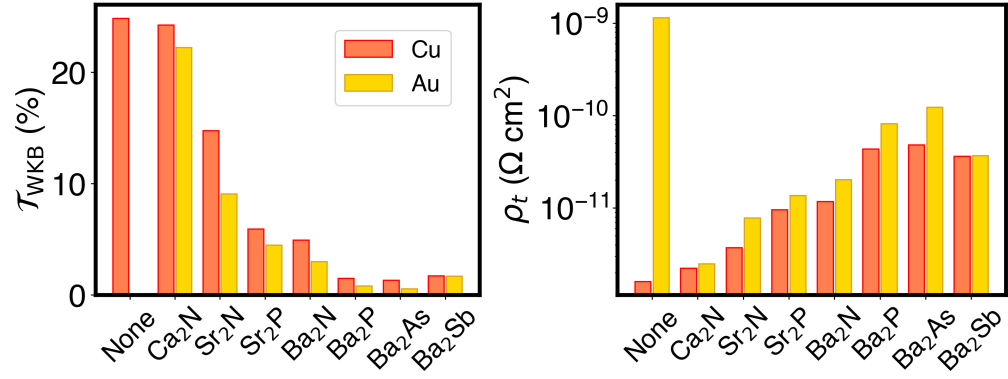


Fig. S22: The WKB tunnelling probability (\mathcal{T}_{WKB} , left) and the tunnelling specific resistivity (ρ_t , right) of the metal/electrene/MoS₂ interfaces, calculated by assuming a square potential.

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