

## Supplementary Information

**Morphology, dynamic disorder, and charge transport in an indoloindole-based hole-transporting material from a multi-level theoretical approach**

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## Kinetic Monte Carlo

The electronic mobility has been evaluated using the kinetic Monte Carlo approach<sup>1</sup> implemented in a home-made code. In the kinetic Monte Carlo, we perform a stochastic dynamic simulation of the hopping of a single charge carrier. We use the direct method based on Gillespie algorithm,<sup>2,3</sup> where an adaptative timestep ( $\tau$ ) is calculated as:

$$\tau = -\frac{\ln(r_1)}{a_0} \quad (1)$$

where  $r_1$  is a random number and  $a_0$  is the summation of the rate constants corresponding to all possible charge hopping paths,  $a_0 = \sum_i k_i$ . On the other hand, the direction of each individual hop is decided by using a second random number ( $r_2$ ) according to the following condition:

$$\frac{\sum_{i=1}^{i-1} k_i}{a_0} < r_2 < \frac{\sum_{i=1}^i k_i}{a_0} \quad (2)$$

where at each hop, the rate constants are computed according to the classical Marcus rate expression.<sup>4</sup> Additionally, in order to take into account the instantaneous fluctuations due to the dynamical disorder,  $V$  and  $\Delta E$  are obtained for each hop according to:

$$V = \text{gauss\_random}(r_3, \langle V \rangle, \sigma_V) \quad (3)$$

$$\Delta E = \text{gauss\_random}(r_4, \langle \Delta E \rangle, \sigma_{\Delta E}) + q\mathbf{F}\mathbf{r} \quad (4)$$

where `gauss_random` is a function that computes a random number ( $r_3$  and  $r_4$ ) according to a normal distribution with given parameters (mean value and standard deviation) previously obtained,  $q$  is the charge,  $\mathbf{r}$  is the displacement vector for the hopping path, and  $\mathbf{F}$  is the applied electric field along the measured direction.

The position of the charge carrier is tracked along the hopping dynamics and the final charge mobility ( $\mu$ ) is computed with a similar expression to those used in master equation approaches:<sup>5,6</sup>

$$\mu = \frac{\mathbf{r}\mathbf{F}}{t|\mathbf{F}|^2} \quad (5)$$

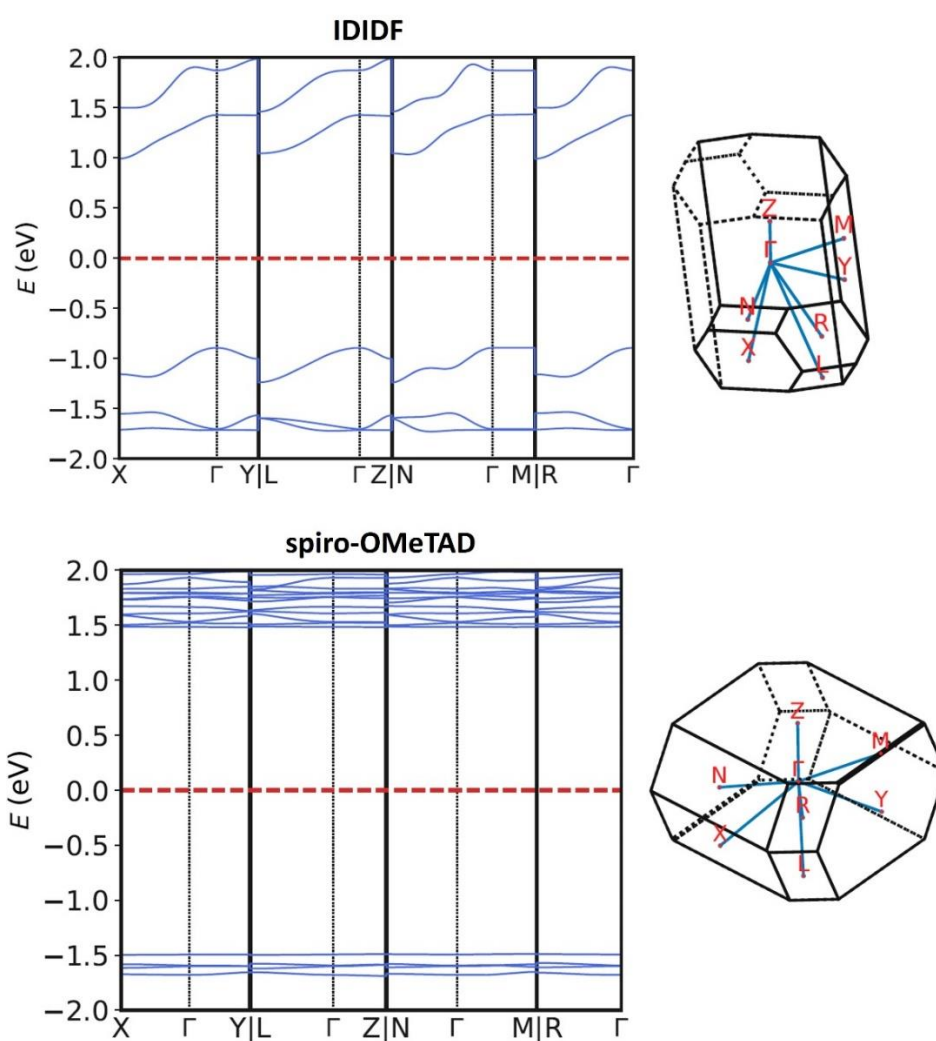
where  $\mathbf{r}$  is the total displacement vector from the initial to the last charge carrier position at the end of the kinetic simulation,  $t$  is the time at the end of the simulation and  $\mathbf{F}$  is the applied electric field.

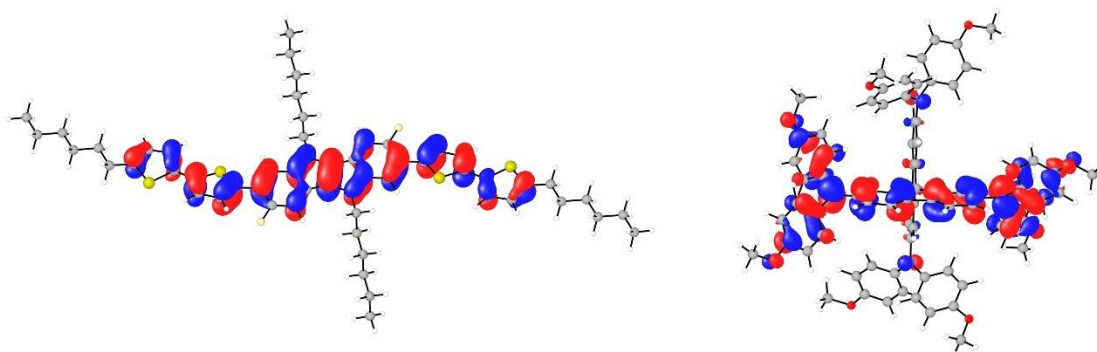
**Table S1.** Crystal cell parameters for IDIDF.

Structure	a (Å)	b (Å)	c (Å)	$\alpha$ (°)	$\beta$ (°)	$\gamma$ (°)
Experimental	5.535	13.714	17.791	68.19	82.84	80.89
DFT	5.549	13.943	17.746	67.13	81.42	79.53
DFT+vdW	5.011	13.077	18.790	66.51	80.73	81.95

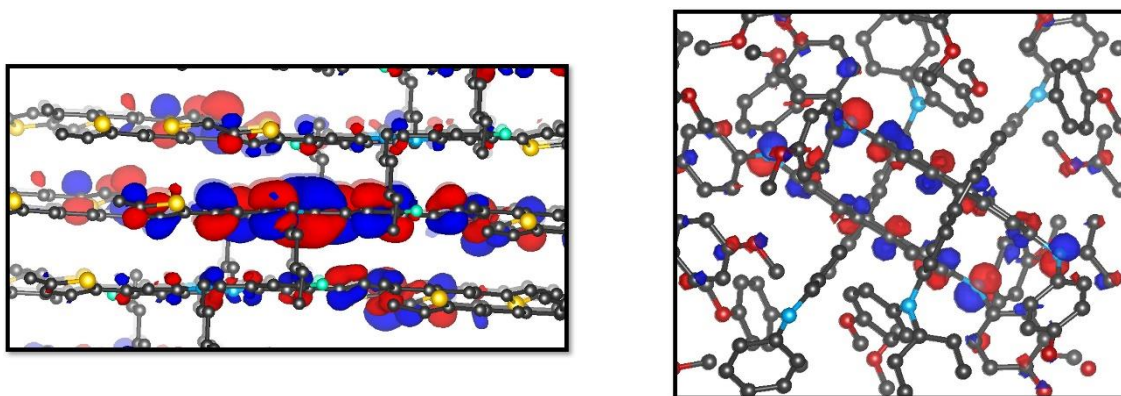
**Table S2.** Crystal cell parameters for spiro-OMeTAD.

Structure	a (Å)	b (Å)	c (Å)	$\alpha$ (°)	$\beta$ (°)	$\gamma$ (°)
Experimental	13.661	14.720	17.277	86.23	68.98	80.01
DFT	13.782	14.629	17.554	85.72	69.37	79.52
DFT+vdW	13.316	14.531	17.015	86.26	68.42	79.88

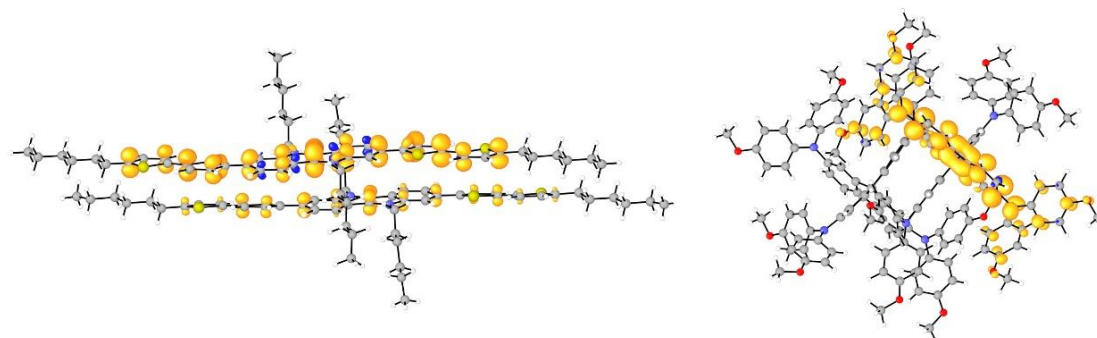
**Figure S1.** Band structure diagrams for spiro-OMeTAD (left) and IDIDF (right) calculated at the HSE06 level along the full k-path of the Brillouin zone  $X-\Gamma-Y|L-\Gamma-Z|N-\Gamma-M|R-\Gamma$  according to the centrosymmetric triclinic space group  $P\bar{1}$  of spiro-OMeTAD and IDIDF.



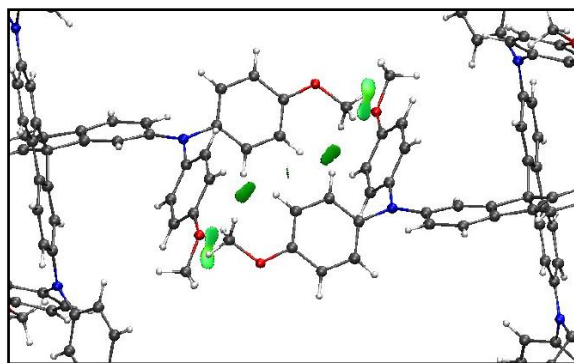
**Figure S2.** Highest-occupied molecular orbitals calculated at the XXX level for IDIDF (left) and spiro-OMeTAD (right).



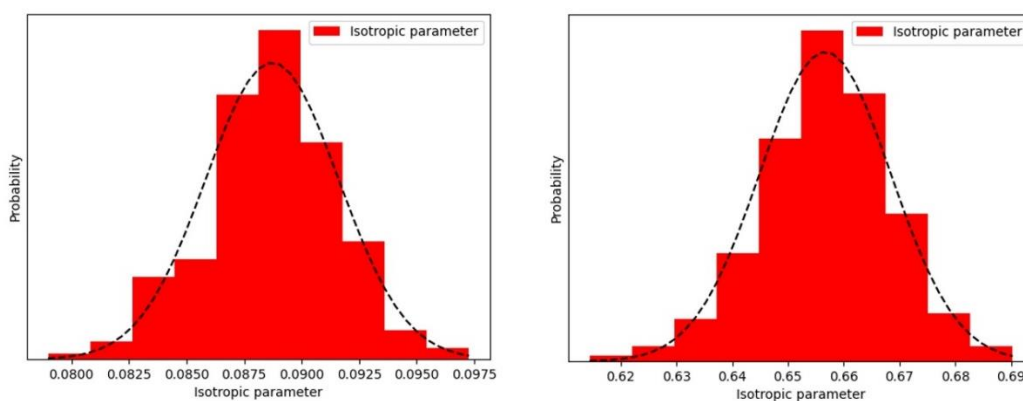
**Figure S3.** Crystal orbital representations of the valence band maximum (VBM) calculated at the HSE06 level of theory on the PBEsol-optimized geometries for IDIDF (left) and spiro-OMeTAD (right).



**Figure S4.** Spin density (isovalue = 0.01) calculated for the most interacting dimers of IDIDF (left) and spiro-OMeTAD (right) where one neutral molecule is replaced by its minimum-energy geometry in the cation state.



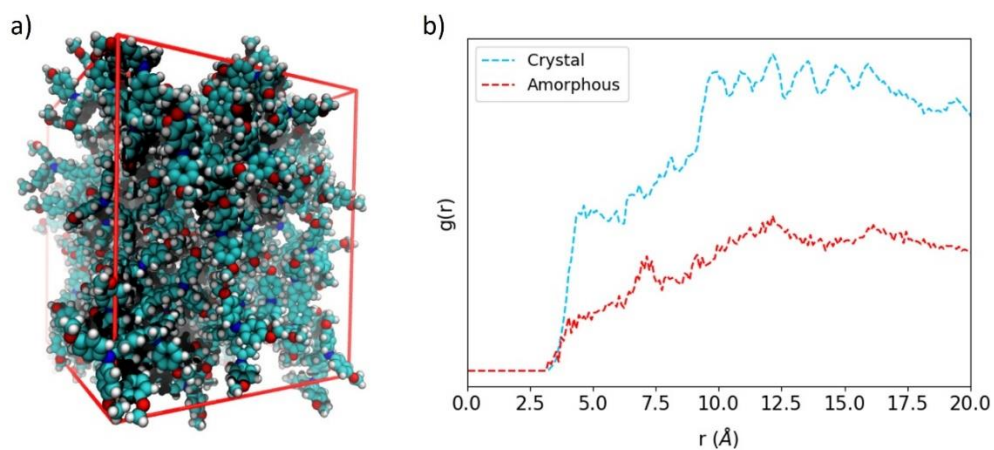
**Figure S5.** Non-covalent surfaces calculated for the  $\bar{1}\bar{1}\bar{1}$  dimer of spiro-OMeTAD, which possesses an electronic coupling of 2 meV.



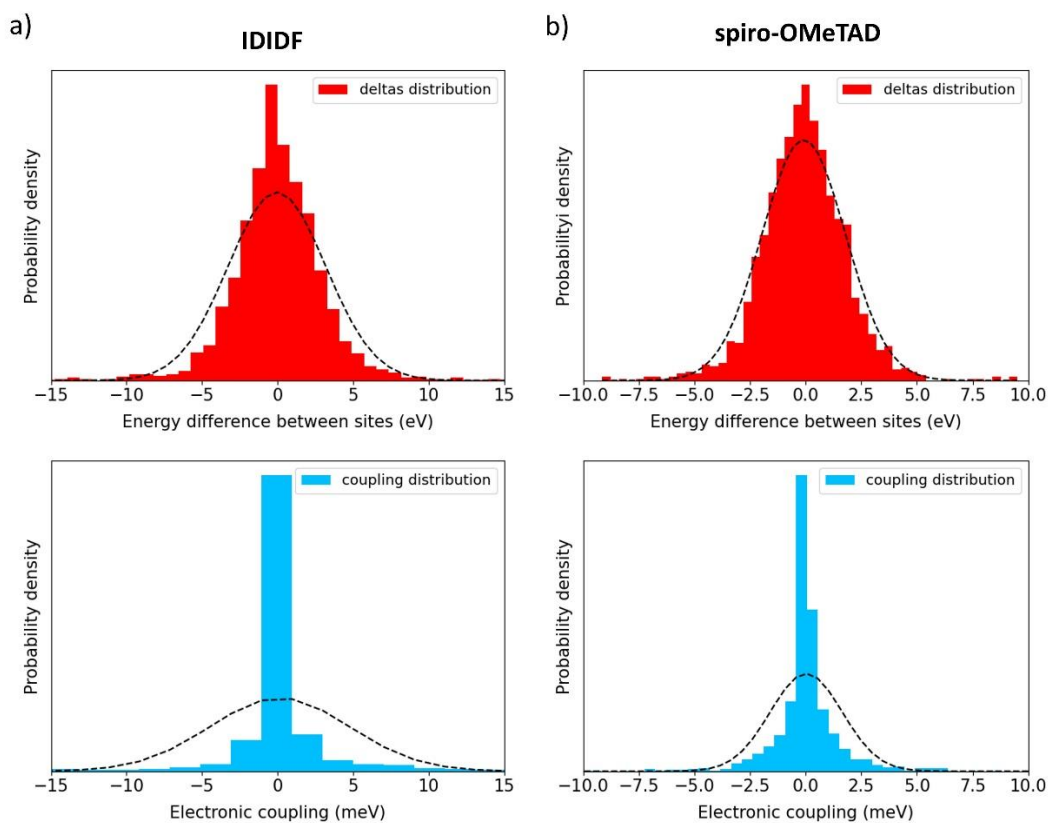
**Figure S6.** Normal distribution of the isotropic ISO parameter calculated for one molecule of IDIDF (left) and spiro-OMeTAD (right) HTMs along the molecular dynamics of the crystalline phase. The distribution parameters are:  $\bar{X} = 0.0887$ ,  $\sigma = 0.003$ ; and  $\bar{X} = 0.6567$ ,  $\sigma = 0.012$  for IDIDF and spiro-OMeTAD, respectively.

**Table S3.** Maximum values for the hole mobility calculated in each crystallographic axis and the mean value obtained considering all the crystallographic planes explored (*ab*, *bc*, and *ac*) in the static crystals (0 K) and considering dynamic disorder (298 K) for IDIDF and spiro-OMeTAD.

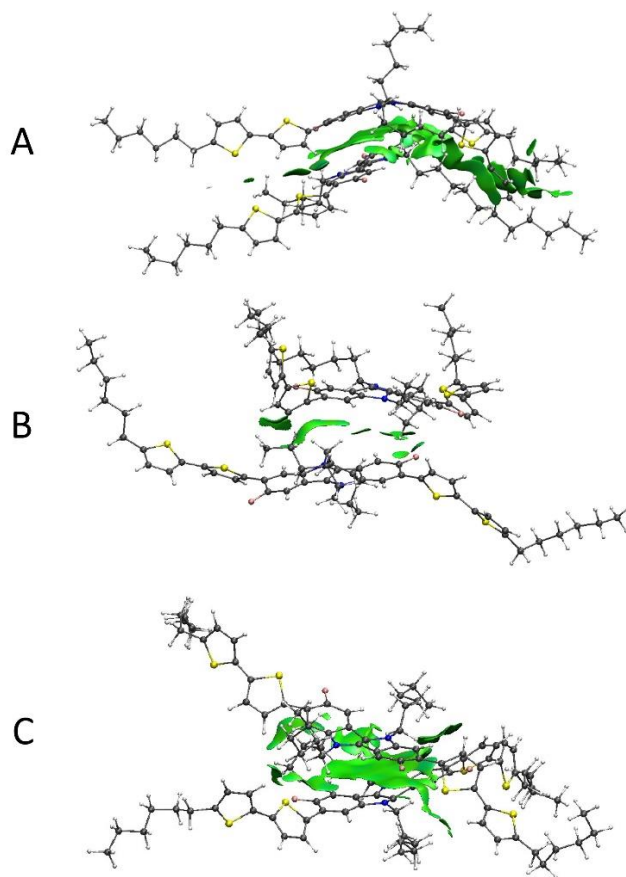
	$\mu_{a\text{-axis}}$	$\mu_{b\text{-axis}}$	$\mu_{c\text{-axis}}$	$\mu_{\text{mean}}$
IDIDF				
Crystal (0 K)	6.806	0.509	2.938	3.361
Crystal (298 K)	2.285	0.161	0.901	1.078
spiro-OMeTAD				
Crystal (0 K)	0.301	0.115	0.093	0.150
Crystal (298 K)	0.041	0.002	0.004	0.014



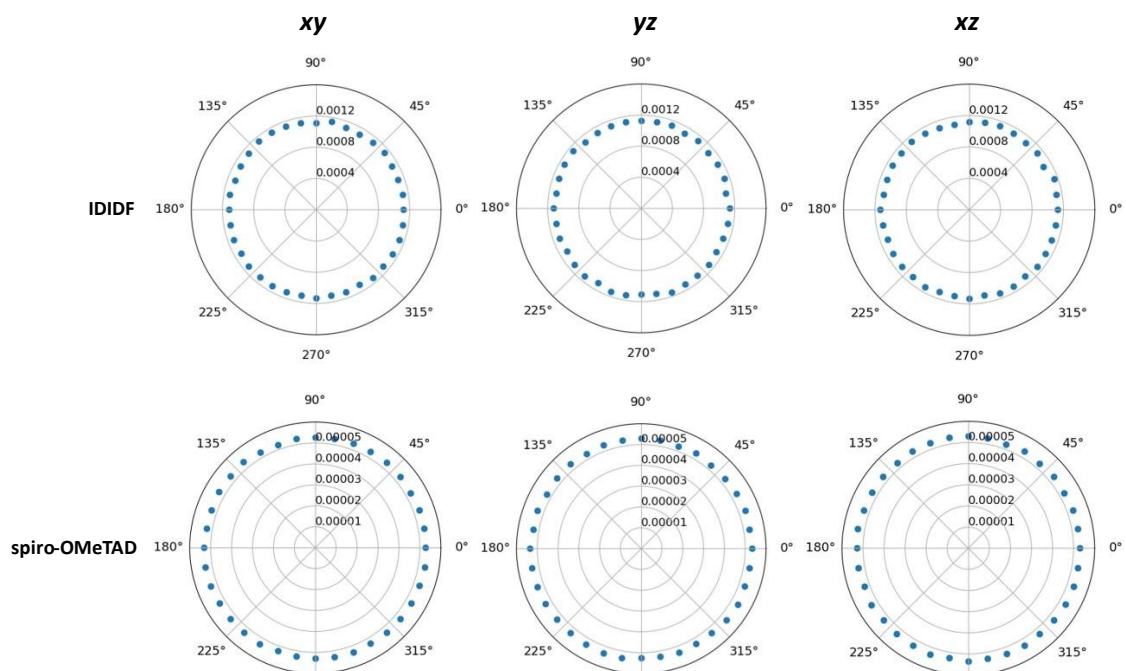
**Figure S7.** a) Representative snapshot of the amorphous spiro-OMeTAD phase. b) Radial distribution function of the centroid of a central molecule with respect to the other molecules in crystalline and amorphous spiro-OMeTAD.



**Figure S8.** Normal distributions for the site energy (top) and dimer electronic coupling (bottom) of IDIF (a) and spiro-OMeTAD (b) obtained along the molecular dynamics simulations of the amorphous materials.



**Figure S9.** Non-covalent NCI surface plots of the reduced density gradient calculated using promolecular densities for dimers A, B and C of amorphous IDIDF indicated in Figure 9 of the main text.



**Figure S10.** Hole mobilities calculated along the  $xy$ ,  $yz$  and  $xz$  planes for the amorphous materials of IDIDF and spiro-OMeTAD.

## References

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