## **Supplementary Information**

Distinct ferrovalley characteristic for Janus RuClX (X = F, Br) monolayer

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Figures S1(a) and S1(b) show the spin charge density distribution of RuClF and RuClBr monolayers with FM and AFM states, respectively. It can be seen that the spin charge density of both spin-up (yellow in Figures S1) and spin-down (cyan in Figures S1) orientations is mainly concentrated on the Ru atoms, while the spin charge density of Cl, F or Br atoms are small enough to be neglected. This proves that the moment of RuClX monolayer is mainly contributed by the Ru atoms.



**Figure S1.** The spin charge density of the FM and AFM states for (a) RuClF and (b) RuClBr monolayers, respectively. The yellow and cyan represent spin up and spin down charge density, respectively. The isovalue is 0.04 e/Bohr<sup>3</sup>.

Figures S2(a) and (b) show the projected density of states of Ru-d, Cl-p and F/Br-p orbitals in RuClF and RuClBr monolayers. It is found that  $d_{xy}$  and  $d_{x^2-y^2}$ ,  $d_{xz}$  and  $d_{yz}$ ,  $p_x$  and  $p_y$  orbitals are degenerate. By comparing Figure S2(a) and Figure S2(b), it can be found that the strength of the unoccupied states (at the energy positive of 1 eV)  $d_{z^2}$  orbitals is smaller in Figure S2(b). Thus, the interaction between the occupied spin-up states of  $d_{yz}$  and the unoccupied spin-down states of  $d_{z^2}$  orbitals in RuClBr monolayer would generate a smaller out-of-plane magnetic anisotropy.





**Figure S2.** (a) The projected density of states (PDOS) of Ru, Cl and F atoms in RuClF monolayer. (b) The PDOS of Ru, Cl and Br atoms in RuClBr monolayer. Fermi energy level is set to zero.

Figures S3(a)-S3(e) show the d-orbitals projected band structures of Ru atoms in the RuClF monolayer. As shown in Figures S3(a) and S3(e), the valley states at the K and K' points are mainly contributed by the occupied  $d_{xy}$  and  $d_{x^2-y^2}$  orbitals of Ru atoms. The  $d_{z^2}$  orbitals of Ru atoms are mainly distributed near the K and K' points of the conduction band, as displayed in Figure S3(c).



Figure S3. (a)-(e) Projected band structures of Ru-d orbitals for the RuClF monolayer. Fermi energy level is set to zero.

Figure S4(a)-S4(e) show the d-orbitals projected band structures of Ru atom in the RuClBr monolayer. From Figures S4(a) and S4(e), it is found that the valley states at the K and K' points are mainly contributed by the unoccupied  $d_{xy}$  and  $d_{x^2-y^2}$  orbitals of Ru atoms. As shown in Figure S4(c), the  $d_{z^2}$  orbitals of Ru atoms are mainly distributed near the K and K' points of the valence band.



**Figure S4.** (a)-(e) Projected band structures of Ru-d orbitals for the RuClBr monolayer. Fermi energy level is set to zero.

The 4d orbitals of Ru atom, the 2p orbitals of F atom, the 3p orbitals of Cl atom, and 4p orbitals of Br atom are considered as the projected orbitals in calculating Wannier functions for the present RuClF and RuClBr monolayers, respectively. The corresponding parameters of the frozen and outer energy windows for Wannier fitting are list in Table S1.

	6	5
System	Frozen energy windows (eV)	Outer energy windows (eV)
RuClF	-10.50, -1.98	-12.0, 2.0
RuClBr	-8.90, -1.0	-9.30, 8.0

Table S1. The parameters of the frozen and outer energy windows for Wannier fitting.

Figure S5 presents the calculated data by the VASP method and the fitting results using

WANNIER functions for RuClF and RuClBr monolayers. It can be seen that the WANNIER fitted band structures are nearly identical with the VASP calculated results.



**Figure S5.** The calculated data by the VASP method and the fitting results using WANNIER functions for (a) RuClF and (b) RuClBr monolayers.

Figure S6 shows the final Wannier spread for each of wave function for RuClF monolayer. The Wannier spread for each of wave function is smaller than the lattice constant. The optimized lattice constant of RuClF monolayer is 3.34 Å.

L T 110	ar state	-							
WF	centre	and	spread	1	(	0.000001,	0.000180,-11.602537	)	0.89308211
WF	centre	and	spread	2	(	0.000001,	0.001903,-11.619107	)	1.08020221
WF	centre	and	spread	3	(	0.000001,	-0.001055,-11.611377	)	0.88166220
WF	centre	and	spread	4	(	0.000003,	-0.005065,-11.630551	)	1.16104057
WF	centre	and	spread	5	(	0.000000,	0.001058,-11.611375	)	0.88161185
WF	centre	and	spread	6	(	0.000001,	0.005087,-11.630554	)	1.16093300
WF	centre	and	spread	7	(	0.000000,	-0.046873,-11.607261	)	0.91706057
WF	centre	and	spread	8	(	0.000001,	-0.105075,-11.621500	)	1.20804216
WF	centre	and	spread	9	(	0.000000,	0.046605,-11.607258	)	0.91709940
WF	centre	and	spread	10	(	-0.000001,	0.102713,-11.621476	)	1.20879943
WF	centre	and	spread	11	(	-0.000003,	1.925647,-10.341487	)	0.59824342
WF	centre	and	spread	12	(	-0.000003,	1.925687,-10.338726	)	0.60468571
WF	centre	and	spread	13	(	0.000012,	1.937572,-10.367278	)	0.63620621
WF	centre	and	spread	14	(	0.000012,	1.936781,-10.362392	)	0.62871263
WF	centre	and	spread	15	(	-0.000002,	1.913677,-10.367278	)	0.63621039
WF	centre	and	spread	16	(	-0.000002,	1.914503,-10.362398	)	0.62875026
WF	centre	and	spread	17	(	-0.000004,	1.925707, 10.233828	)	1.41587931
WF	centre	and	spread	18	(	-0.000005,	1.925909, 10.217947	)	1.42087307
WF	centre	and	spread	19	(	0.000004,	1.939550, 10.249597	)	1.31243121
WF	centre	and	spread	20	(	0.000005,	1.942181, 10.236011	)	1.30918615
WF	centre	and	spread	21	(	-0.000003,	1.911679, 10.249595	)	1.31244779
WF	centre	and	spread	22	(	-0.000004,	1.909105, 10.236010	)	1.30920922
Sur	n of cer	ntres	s and spre	ads	(	0.000013,	23.107477,*********	)	22.12236888

Figure S6. The final Wannier spread for each of wave function for RuClF monolayer.

Figure S7 shows the final Wannier spread for each of wave function for RuClBr monolayer. The Wannier spread for each of wave function is smaller than the lattice constant. The optimized lattice constant of RuClBr monolayer is 3.66 Å.

E T 110	ar state								
WF	centre and	i spread	1	(	0.000000,	0.000057,	11.681388	)	0.96537273
WF	centre and	l spread	2	(	0.000001,	0.000203,	11.683048	)	1.09302935
WF	centre and	l spread	3	(	0.000001,	-0.007088,	11.688042	)	0.94655624
WF	centre and	i spread	4	(	0.000001,	-0.012083,	11.700482	)	1.23992672
WF	centre and	l spread	5	(	0.000000,	0.007088,	11.688039	)	0.94641136
WF	centre and	l spread	6	(	0.000000,	0.012080,	11.700483	)	1.23968387
WF	centre and	i spread	7	(	0.000000,	-0.040588,	11.682123	)	0.95732745
WF	centre and	l spread	8	(	0.000001,	-0.090871,	11.685562	)	1.32573392
WF	centre and	l spread	9	(	0.000000,	0.040463,	11.682122	)	0.95733716
WF	centre and	l spread	10	(	0.000000,	0.090251,	11.685556	)	1.32589401
WF	centre and	l spread	11	(	-0.000006,	2.114131,	10.124889	)	1.36969418
WF	centre and	l spread	12	(	-0.000004,	2.114212,	10.119799	)	1.38526594
WF	centre and	l spread	13	(	0.000005,	2.139114,	10.181995	)	1.39106285
WF	centre and	l spread	14	(	0.000005,	2.140306,	10.170200	)	1.37824615
WF	centre and	l spread	15	(	-0.000004,	2.088914,	10.181994	)	1.39107555
WF	centre and	l spread	16	(	-0.000004,	2.087798,	10.170198	)	1.37827961
WF	centre and	l spread	17	(	-0.000004,	2.114181,	-10.087780	)	1.73742665
WF	centre and	l spread	18	(	-0.000004,	2.114263,	-10.078330	)	1.75165003
WF	centre and	l spread	19	(	0.000004,	2.138787,	-10.127120	)	1.72192726
WF	centre and	l spread	20	(	0.000004,	2.141363,	-10.111972	)	1.71407209
WF	centre and	l spread	21	(	-0.000003,	2.089195,	-10.127117	)	1.72196884
WF	centre and	l spread	22	(	-0.000003,	2.086722,	-10.111967	)	1.71414369
Su	n of centre	es and sp	reads	(	-0.000010,	25.368499,	117.181634	)	29.65208566

Figure S7. The final Wannier spread for each of wave function for RuClBr monolayer.

As observed in Figure S8(a), with 0.05 holes per f.u. are doped, the Fermi energy level passes

through the K' valley in RuClF monolayer. Therefore, the carriers are dominated by the K' valley. The Fermi energy level passes through the K valley in RuClBr monolayer with doping 0.05 electrons per f.u., as shown in Figure S8(b). At this time, the carriers are excited by the K valley. In the first-principles calculations, carrier doping is simulated by removing or adding electrons from the system and using a homogeneous background charge to maintain charge neutrality<sup>[1,2]</sup>.



**Figure S8.** (a) The band structures with SOC of RuClF monolayer with doping 0.05 holes per f.u.. (b) The band structures with SOC of RuClBr monolayer with doping 0.05 electrons per f.u.. The red arrow represents spin up, and the blue arrow represents spin down. Fermi energy level is set to zero.

## References

- X. Li, X. Wu, J. Yang, Half-Metallicity in MnPSe<sub>3</sub> Exfoliated Nanosheet with Carrier Doping, Journal of the American Chemical Society, 136 11065 (2014).
- [2] Y. Chen, Q. Fan, Y. Liu, G. Yao, Electrically tunable magnetism and unique intralayer charge transfer in Janus monolayer MnSSe for spintronics applications, *Physical Review B*, **105** 195410 (2022).