

Supplementary information:
Model systems for dye-sensitized solar cells: cyanidin-silver
nanocluster hybrids at TiO₂ support

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Photovoltaic parameters

Incident photon to conversion efficiency (IPCE) of dye-sensitized solar cells (DSSC):

$$IPCE = LHE \cdot \Phi_{inj} \cdot \eta_c \quad (1)$$

can be evaluated as a product of light harvesting efficiency (LHE), electron injection efficiency (Φ_{inj}) and charge collection efficiency (η_c).

Light harvesting efficiency is approximated using the equation:

$$LHE = 1 - 10^{-f} \quad (2)$$

where f is the oscillator strength of maximum absorption. High performance of DSSCs is related to the large value of the LHE.

Electron injection efficiency (Φ_{inj}), is proportional to the driving force (ΔG^{inject}), which corresponds to the difference between oxidation potential energy of the excited state (E_{ox}^{dye*}), and reduction potential energy of conduction band (ECB).¹

$$E_{ox}^{dye*} = E_{ox}^{dye} - \lambda_{max} \quad (3)$$

where redox E_{ox}^{dye} is potential of the ground state, and λ_{max} is maximum absorption energy of the sensitizer. Finally, ΔG^{inject} can be written as:

$$\Delta G^{inject} = (E_{ox}^{dye} - \lambda_{max}) - ECB \quad (4)$$

$$E_{ox}^{dye} = -HOMO \quad (5)$$

Experimental ECB of TiO₂ has a value of -4 eV²⁻⁴ vs. vacuum.

To obtain effective photo-induced electron transfer from the dye sensitizer to the semiconductor, energy levels of the highest occupied (HOMO) and lowest unoccupied (LUMO) molecular orbitals of the dye are required to match the redox potential of the electrolyte, I⁻/I₃⁻ (-4.6 eV⁵ vs. vacuum) and the conduction band edge of the semiconductor.¹

Binding energy

The binding energy of the hybrid system E_b corresponds to the:

$$E_b = E_{hybrid} - (E_{cyanidin} + E_{NC}) \quad (6)$$

where E_{hybrid} , $E_{cyanidin}$ and E_{NC} are energies of optimized cyanidin-NC hybrid, cyanidin dye, and NC systems, respectively.

The binding (adsorption) energy E_{ads} of the cyanidin-Ag₉ at TiO₂ surface model is defined as:

$$E_{ads} = E_{complex} - (E_{TiO_2} + E_{cyanidin-Ag_9}) \quad (7)$$

$E_{complex}$ is the energy of adsorbed cyanidin-Ag₉, E_{TiO_2} is the energy of the optimized TiO₂ model and $E_{cyanidin-Ag_9}$ energy of the optimized cyanidin-Ag₉.

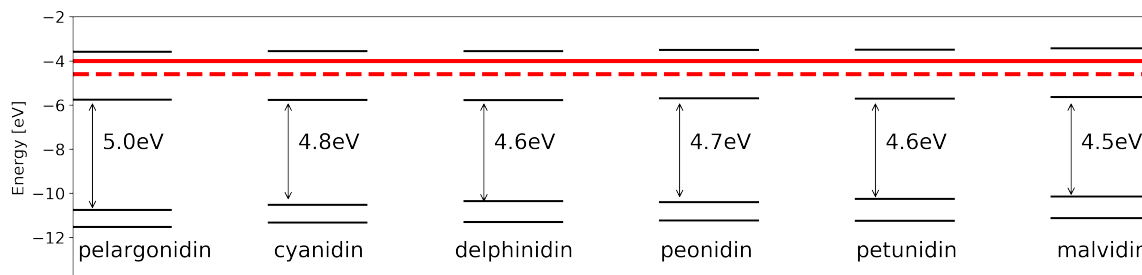


Figure S1: DFT calculated HOMO–LUMO (and HOMO-1, LUMO+1) energy gaps versus vacuum[eV] for the anthocyanidin dyes in comparison with the experimental TiO_2 conduction band edge (red line) and I^-/I_3^- redox level (dashed red line).

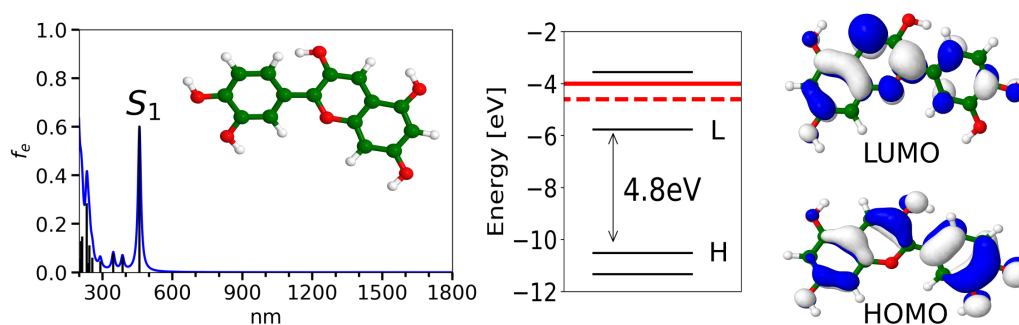


Figure S2: TDDFT calculated absorption spectrum for cyanidin dye employing CAM-B3LYP/def2-SVP method. The structure has been optimized at PBE/def2-SVP level of theory. DFT HOMO, LUMO orbitals, and HOMO, LUMO, HOMO-1 and LUMO+1 energy gaps versus vacuum. Experimental TiO_2 conduction band edge (full red line) and I^-/I_3^- redox level (dashed red line).

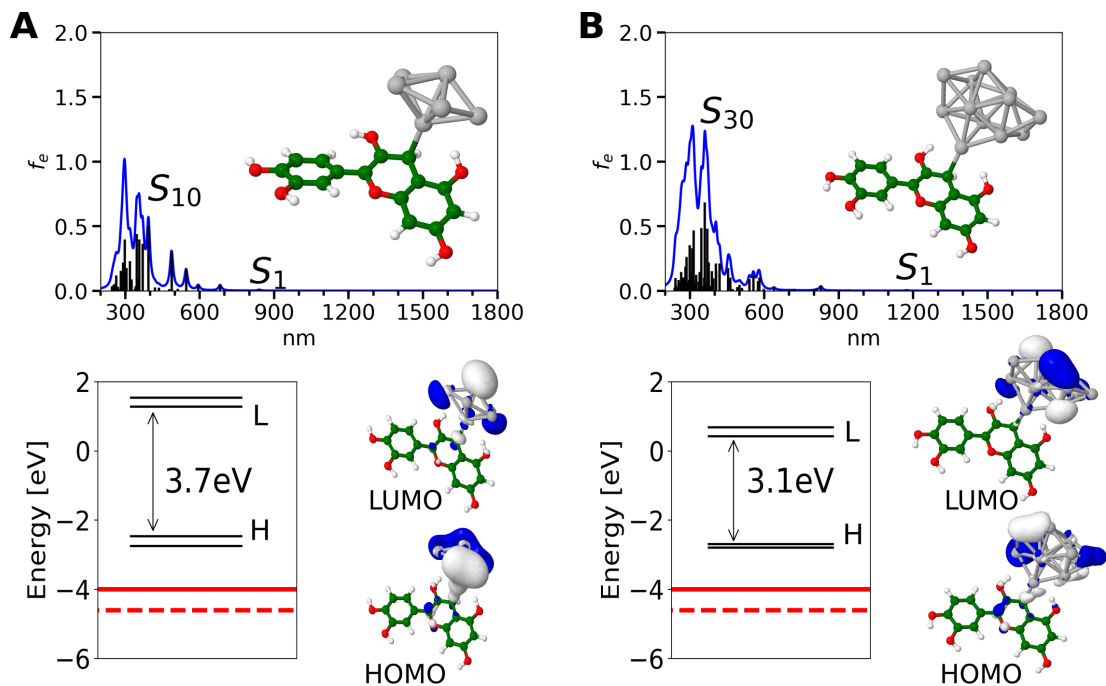


Figure S3: TDDFT calculated absorption spectra for the: A) (cyanidin-Ag₆)⁻, B) (cyanidin-Ag₁₂)⁻ employing CAM-B3LYP/def2-SVP method. Structures have been optimized at PBE/def2-SVP level of theory.

DFT HOMO, LUMO, HOMO-1, and LUMO+1 energy gaps versus vacuum[eV] and HOMO-LUMO orbitals. Experimental TiO₂ conduction band edge (full red line) and I⁻/I₃⁻ redox level (dashed red line)

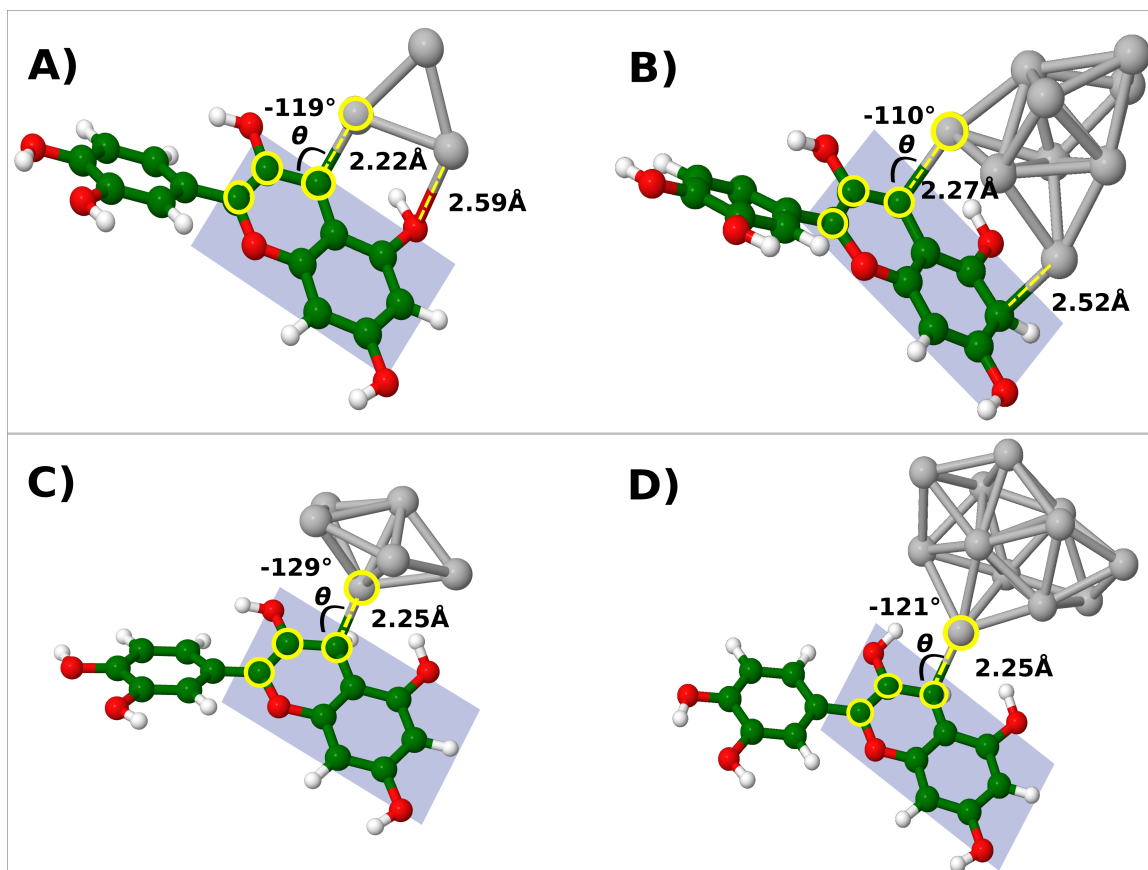


Figure S4: The bond lengths and dihedral angles between cyanidin and silver NC of hybrid systems are shown for A) cyanidin- Ag_3 , B) cyanidin- Ag_9 with an even number of electrons, and C) (cyanidin- Ag_6)⁻, D) (cyanidin- Ag_{12})⁻ with an odd number of electrons. The blue rectangle represents the plane of the indoline group of the cyanidin dye, θ is the dihedral angle between the plane and silver NC.

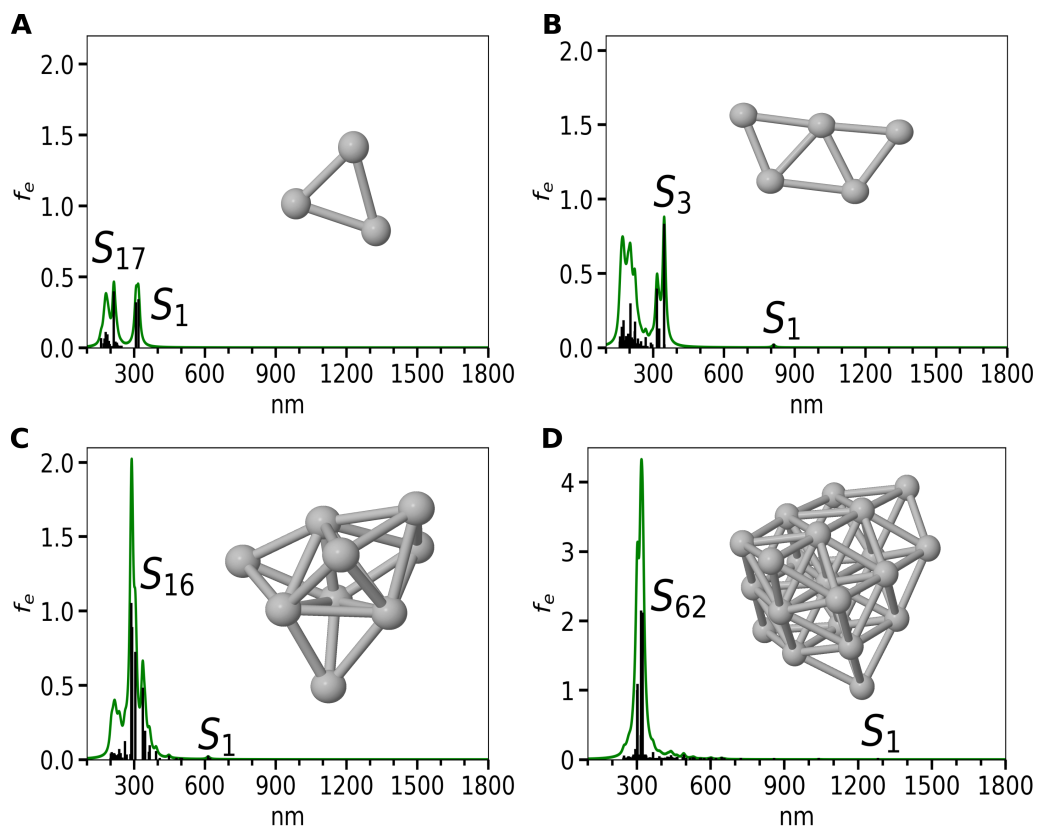


Figure S5: TDDFT calculated absorption spectra for the: A) Ag_3^+ , B) Ag_5^+ , C) Ag_9^+ and D) Ag_{21}^+ clusters with structures from cyanidin-NC hybrids at CAM-B3LYP/def2-SVP level of theory.

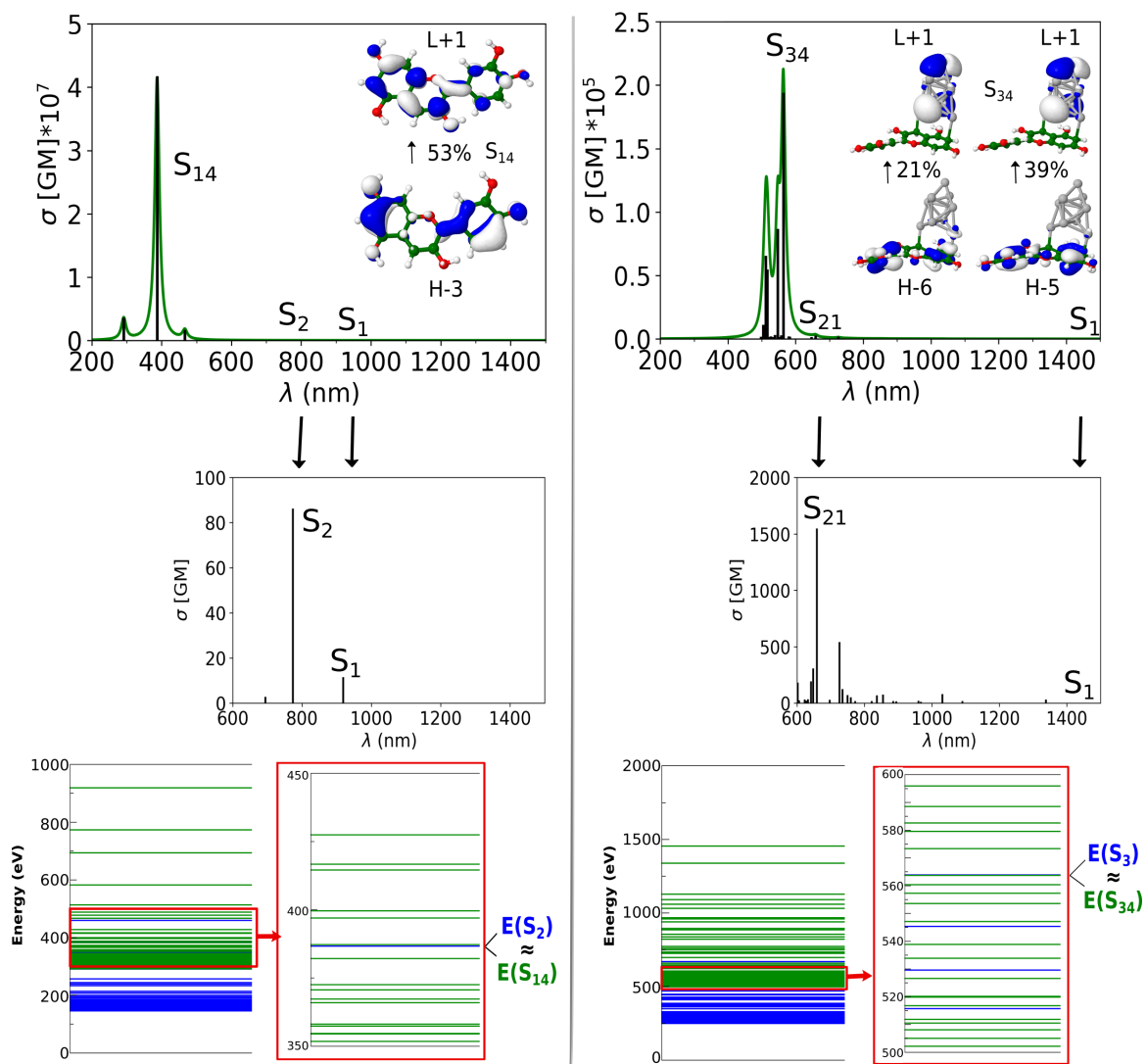


Figure S6: TPA spectra obtained by TDDFT for cyanidin (left) and cyanidin-Ag₉ (right). MO analysis is shown. TPA cross sections for states in NIR regime are also presented (middle). The resonance between OPA (blue) and TPA (green) states is included (bottom).

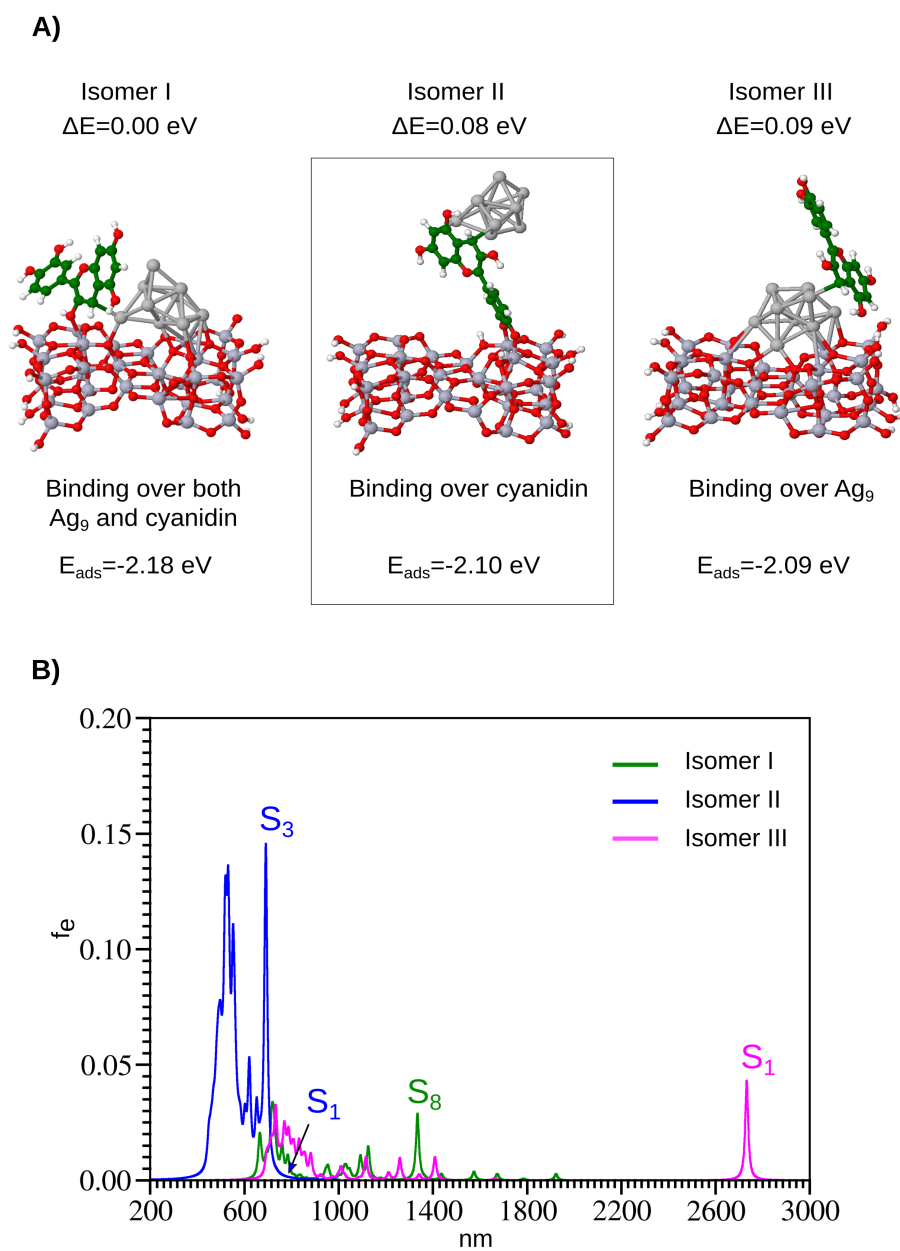


Figure S7: A) Isomers of {cyanidin-Ag₉}-TiO₂ with different adsorption modes. B) TDDFT calculated absorption spectra for isomers I, II and III of {cyanidin-Ag₉}-TiO₂ at CAM-B3LYP/def2-SVP level of theory. S₁ state of isomer I is located at 6090 nm ($f_e=0.027$).

Table S1: Excited states corresponding to the largest calculated absorptions for {cyanidin-Ag₉}-TiO₂ at CAM-B3LYP/def2-SVP level of theory with calculated wavelengths λ , oscillator strengths f_e and predominant transitions

Excited state	$\lambda[nm]$	f_e	Transitions
S ₃	691	0.1425	H → L+25 (39%) H → L+26 (26%) H → L+29 (11%)
S ₇	620	0.0440	H → L+2 (38%) H → L+14 (13%) H → L+12 (12%)
S ₁₄	551	0.0742	H → L+49 (20%) H → L+51 (16%) H → L+52 (10%)
S ₁₈	531	0.0868	H → L+19 (19%) H-1 → L+25 (15%)

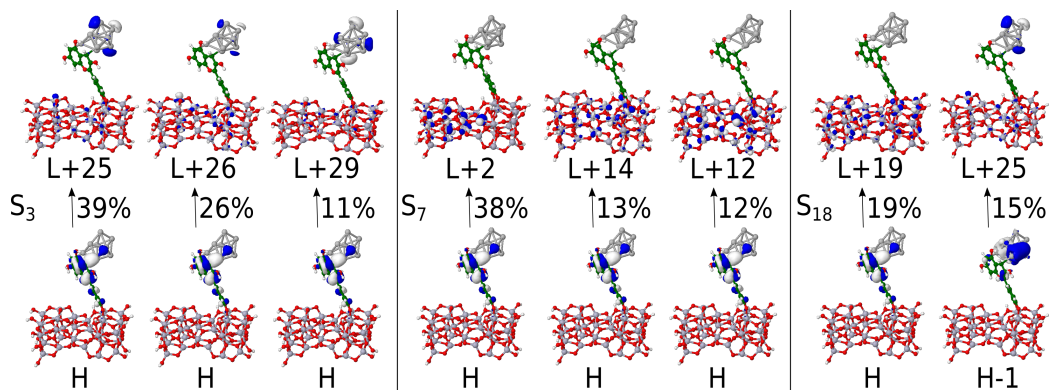


Figure S8: Transitions representing S₃, S₇ and S₁₈ states are shown. S₃ state corresponds to transitions from HOMO orbital delocalized on the hybrid to higher LUMO orbitals delocalized partly at support and partly at Ag₉. The transitions associated with the excited states higher in energy (S₇, S₁₈) demonstrate clear charge transfer from HOMO (HOMO-1) to higher LUMO orbitals delocalized entirely at TiO₂ model.

References

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