# Supplemental Information: A multiscale approach to uncover the self-assembly of ligand-covered palladium nanocubes 

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We chose the Paul force field over commonly used OPLS force field because OPLS tends to model the alkane ligand as "brush-like" rods instead of interdigitating woven[1, 2]. We tested both Paul and OPLS force field in vacuum and visualized their difference in ligand modeling in Figure 1. It is clear that OPLS force field modeled dodecanethiol as "stiff" molecules standing up straight and Paul force field model them as interlacing molecules. We ran the comparison at $298,325,373 \mathrm{~K}$ and the difference persisted.


Figure S1: Visualization of dodecanethiol covered palladium nanoparticle using OPLS force field (a) and Paul force field (b) at room temperature.

Table S1: Harmonic bond parameter for palladium - sulfur[1]

|  | $k\left(\mathrm{kcal} / \mathrm{mol} / \AA^{2}\right)$ | $\mathrm{r}_{0}(\AA)$ |
| :---: | :---: | :---: |
| Pd-S | 265 | 2.4 |

Table S2: 12-6 Lennard-Jones parameters for Palladium atoms in a face-centered particle[3]

|  | $\varepsilon(\mathrm{kcal} / \mathrm{mol})$ | $\sigma(\AA)$ |
| :---: | :---: | :---: |
| Pd | 6.15 | 2.512 |

Table S3: 12-6 Lennard-Jones parameters of ligand molecule in opti-mized unitedatom model for simulations of polymethylene melts

|  | $\varepsilon(\mathrm{kcal} / \mathrm{mol})$ | $\sigma(\AA)$ |
| :---: | :---: | :---: |
| $\mathrm{CH}_{3}$ | 0.22644 | 4.009 |
| $\mathrm{CH}_{2}$ | 0.09344 | 4.009 |
| SH | 0.3 | 4.25 |

Table S4: Bond parameters of ligand molecule in optimized united-atom model for simulations of polymethylene melts

| Harmonic bond interaction | $k(\mathrm{kcal} / \mathrm{mol})$ | $\mathrm{r}_{0}(\AA)$ |
| :---: | :---: | :---: |
| C-C bond | 444 | 1.81 |
| S-C bond | 643 | 1.53 |

Table S5: Angle and dihedral parameters of ligand molecule in optimized unitedatom model for simulations of polymethylene melts ( $k_{i}$ is in $\mathrm{kcal} / \mathrm{mol}$ )[2]

| Angle and dihedral interaction | $\theta_{0}$ (degree) | $k_{1}$ | $k_{2}$ | $k_{3}$ |
| :---: | :---: | :---: | :---: | :---: |
| C-C-C angle | 110.01 | 120 | N/A | N/A |
| S-C-C angle | 113.4 | 125 | N/A | N/A |
| C-C-C-C dihedral | N/A | 1.6 | -0.867 | 3.24 |
| S-C-C-C dihedral | N/A | 1.6 | -0.867 | 3.24 |

Table S 6: 12-6 Lennard-Jones parameters of toluene seven point united-atom model of TraPPE-UA force field[4]

|  | $\varepsilon(\mathrm{kcal} / \mathrm{mol})$ | $\sigma(\AA)$ |
| :---: | :---: | :---: |
| CH (aromatic carbon) | 0.1003 | 3.695 |
| $\mathrm{CH}-[\mathrm{C}]-\mathrm{CH}_{3}$ | 0.04173 | 3.88 |
| $\mathrm{CH}_{3}$ | 0.1947 | 3.75 |

Table S7: Bond and angle parameters of toluene in united-atom model of TraPPEUA force field[4]

| Angle and dihedral interaction | $\mathrm{r}_{0}(\AA)$ | $\theta_{0}($ degree $)$ | $k(\mathrm{kcal} / \mathrm{mol})$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{CH}-\mathrm{CH}$ bond | 1.54 | N/A | rigid |
| $\mathrm{CH}-\mathrm{CH}_{3}$ bond | 1.40 | N/A | rigid |
| $\mathrm{CH}-\mathrm{CH}-\mathrm{CH}$ angle | N/A | 120 | rigid |
| $\mathrm{CH}-\mathrm{CH}-\mathrm{CH}_{3}$ angle | N/A | 120 | rigid |

Table S8: Free energy barrier of transition from face-to-face to brick-wall configuration in vacuum and toluene. All units are in $\mathrm{kcal} / \mathrm{mol}$.

|  | C-6 | C-9 | C-12 | C-15 |
| :---: | :---: | :---: | :---: | :---: |
| Vacuum | 1356 | 441 | 329 | 204 |
| Toluene | 1641 | 478 | 251 | N/A |



Figure S2: Comparison of free energy landscapes provided by TI in vacuum (blue) and in a coarse-grained toluene solvent (orange). (a)-(c) show results for C-6, C-9 and C-12 ligands, respectively.

(a) [Left] Steered Molecular Dynamics (SMD) values for the free energy landscape as a function of time. [Right] Corresponding Thermodynamic Integration results for the same process as a function of distance, which are linearly related. The SMD-derived peak at about 20 ns occurs when the system adopts a brickwall configuration. Note that it is typical for the energy values to differ significantly between SMD and TI, with the latter typically assumed to be more accurate.

(b) [Left] Degree of rotation of the top particle relative to the bottom layer of the Pd particles during SMD simulations for the results shown in Figure 3(a). The particle rotates by $\sim 11^{\circ}$ at about 10 ns and maintains that tilt throughout the rest of the transition. [Right] Lack of correlation between the degree of rotation of the top particle (shown on the left) and the free energy of the system.

Figure S3: Thermodynamic profiles of the Steered Molecular Dynamic simulation of the configurational change in a 3-nanoparticle system.


Figure S4: Visualization of the steered Molecular Dyanmics simulation trajectory of 3-nanoparticle system undergoes face-to-face to brick-wall configuration change. Palladium atoms are displayed as dots and ligand molecules are displayed as lines.


Figure S5: Snapshot of an starting (a), intermediate (b) and the final (c) frame of the simulation of described in Figure 2 in the method section. In (a), the pulled nanoparticle is at the initial face-to-face configuration. In (b), it has arrived at the brick-wall configuration in its local lattice site, exhibiting a $15-18^{\circ}$ rotation with respect to the bottom layer of Pd nanoparticles. Note, the movement of the nanoparticle nearest to the moving particle towards the moving nanoparticle. In (c), the pulled nanoparticle has reached its destination but maintained $\sim 11^{\circ}$ of rotation in spite of the spring force continuing to exist. All solvent molecules are hidden for clarity.

(a) Free energy landscape of the SMD simulation il- (b) Rotation of nanopaticle B with respect to the botlustrated in Figure 5. The significant increase in free tom layer of the superlattice as a function of time energy reflects the difficulty of overcoming the inter- during the simulation, showing a maximum rotation molecular forces needed to continuously pull B to- of $15-18^{\circ}$ after $4-6 \mathrm{~ns}$, a time corresponding to the wards A. brick wall position. It then returns to a tilt of about $11^{\circ}$ as B moves into its final 22 F configuration adjacent to A .

(c) Relationship between the degree of rotation and (d) Relationship between the degree of rotation and the displacement of nanopaticle B. It shows that the free energy change during the pulling B to the nanopaticle reached a maximum rotation of 15 - healed position showing that the maximum rotation $18^{\circ}$ around 20 to $28 \AA$ corresponding to the brick- $\left(\sim 18^{\circ}\right)$ corresponds to the lowest free energy and that wall configuration. It then returns to a tilt of about the particle remains tilted by about $11^{\circ}$ as it returns $11^{\circ}$ as B moves into its final F2F configuration adja- to the F2F position cent to A.

Figure S6: Thermodynamic profiles of the Steered Molecular Dynamic simulation illustrated in Figure 5 Parts (a)-(d) described above.

## References

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