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## Supporting Information for

Nanohybrids with switchable multicolor emission for anticounterfeiting

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Figure S1. (a-g) UCL spectra of different samples. Inset of (b-k): Digital photographs of different samples under 980 nm laser diode.



Figure S2. (a-h)  $(Ahv)^2$ -hv plot of different CsPbX<sub>3</sub> QDs glasses: the oblique line is the fitted curve to estimate the  $E_g$ .



Figure S3. (a, b) XRD patterns of different CsPbX<sub>3</sub> QDs glasses and nanohybrids, respectively.



Figure S4. The partial magnification spectroscopy of the FTIR peaks shift.

	362 nm		450 nm		478 nm		524 nm	
UCP/QDs	Α	<b>τ(μs)</b>	Α	τ(μs)	А	τ(μs)	А	τ(μs)
UCP	1.07	430.28	1.07	426.55	1.17	983.20		
8:1	1.07	427.35	1.07	424.40	1.16	920.21	1.05	490.19
2:1	1.07	424.19	1.06	420.69	1.15	895.57	1.05	504.63
1:2	1.07	420.74	1.07	419.84	1.16	883.90	1.05	533.77
1:8	1.07	417.01	1.05	414.34	1.14	851.28	1.05	551.04

**Table S1** The fitting parameters of UCL decays lifetime for Tm/Yb: NaYF<sub>4</sub> UCP+CsPbBr<sub>3</sub> QDs glasses nanohybrids at different emission bands excited by 980 nm.

Both  $Tm^{3+}$  ions and  $CsPbBr_3$  QDs UCL lifetimes in **nanohybrids** V are obtained by singe exponential decay <sup>1</sup>:

$$I = A \exp(-t / \tau) \qquad (S1)$$

Where I is the luminescence intensity of  $Tm^{3+}$  ions or CsPbB<sub>3</sub> QDs; t is the time; A is fitting constants;  $\tau$  is fluorescence lifetimes for the exponential components.

The DCL lifetime of CsPbBr<sub>3</sub> QDs glasses are obtained by double exponential decay <sup>1</sup>:

$$I = A_1 \exp(-t / \tau_1) + A_2 \exp(-t / \tau_2)$$
 (S2)

![](_page_5_Figure_0.jpeg)

Figure S5. DCL decay curve of CsPbBr<sub>3</sub> QDs glasses.

## Note 1: Förster radius (R<sub>0</sub>)

Förster radius  $R_0$  is the distance between donor and acceptor when the energy transfer efficiency is 50%, the equation is defined as <sup>2</sup>:

$$R_0^{6} = \left[\frac{9(\ln 10)\kappa^2 \eta_D}{128\pi^5 N_A n^4} J\right]$$
 (83)

Known from the **Equation S3** that Förster radius  $R_0$  is proportional to the orientation factor of the interacting dipoles ( $\kappa^2$ ), the donor's UCL fluorescence quantum yield in the absence of acceptor ( $\eta_D$ ), spectral overlap integral ( $M^{-1}$  cm<sup>-1</sup> nm<sup>4</sup>) (*J*), and is inversely proportional to the Avogadro constant ( $N_A$ ), and the average refractive index of the medium (n). The spectral overlap integral *J* is calculated as <sup>2</sup>:

$$J(\lambda) = \int F_{\rm D}(\lambda) \varepsilon_{\rm A} \lambda^4 d\lambda \qquad (S4)$$

Where  $F_D(\lambda)$  is the UCL spectrum of the donor normalized to unit area, which is 5.4\*10<sup>16</sup>,  $\varepsilon_A$  is the molar extinction coefficient (M<sup>-1</sup> cm<sup>-1</sup>) of quantum dots, which is 5\*10<sup>6</sup> M<sup>-1</sup>cm<sup>-1</sup>,  $\lambda$  is the wavelength range of spectral overlap of donner and acceptor (300~600 nm). Herein, assuming that interacting dipoles  $\kappa^2$  is a random orientation of donor and acceptor, thus the value of  $\kappa^2$  is 2/3 (the most used value in the literature), the value of  $\eta_D$  is 0.01%, the average refractive index of the medium (glass and UCP) is assumed as 1.5.

## Note 2 Determination of the slope of the double logarithmic plots of the UCL intensity vs. excitation power density.

The relationship between UCL intensity and the laser power densities of 980 nm is defined as <sup>3</sup>,

 $I_{UCL} \propto (I_{NIR})^n$  (S5)

Where n refers to the number of photons.

## References

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