ELECTRONIC SUPPORTING INFORMATION

Controlled Radical Release with Iron Oxide Nanoparticles Grafted with Thermosensitive Alkoxyamines Triggered by External *Stimuli*

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Models & theory

1. Determination of the Specific Absorption Rate (SAR).

The SAR (in $W g^{-1}$) is defined as follows:

$$SAR_{Fe} = \frac{P}{m_{Fe}} \quad or \quad SAR_{NP} = \frac{P}{m_{NP}}$$
 (1)

With SAR_{Fe} and SAR_{NP} the SAR defined with the mass of Fe (m_{Fe}) or the mass of the magnetic nanoparticle (NP) (m_{NP}), respectively. *P* represents the power produced by the magnetic hyperthermia phenomenon. In this article, we will use m_{NP} as reference.

Usually, the SAR values are calculated in the adiabatic approximation (no energy exchange between the sample and the environment):

$$P = m_f c_f \cdot \frac{dT}{dt} \tag{2}$$

$$SAR_{NP} = \frac{m_f c_f}{m_{NP}} \cdot \frac{dT}{dt}$$
(3)

where m_f and c_f are the mass and the weighted heat capacity of the fluid, respectively.

If we want to take into account the possibility of energy exchange between the sample and the environment, we need to rewrite the equation (2). The Newton temperature equation law is a simple phenomenological model, which tends to reproduce the thermal evolution of such case. Equation (2) becomes:

$$m_f c_f \cdot \frac{dT}{dt} = P - h \cdot A \cdot (T - T_0) \tag{4}$$

with *h* is a phenomenological term for the energy transfer between the sample and the environment, *A* represents the area where the transfer occurs and T_0 is the temperature of the environment.

The solution of the differential equation (4) is:

$$T(t) = \frac{P}{hA} + T_0 - \frac{P}{hA} \cdot \exp\left(-\frac{hA}{m_f c_f}t\right)$$
(5)

However, in some cases this phenomenological law cannot be sufficient. More generally, the energy transfer between the sample and the environment can be represented by an unknown function, which depends on *T* and *T*₀. We will define this function as $f(T,T_0)$. Then, equation (2) becomes:

$$m_f c_f \cdot \frac{dT}{dt} = P - f(T, T_0) \tag{6}$$

If we consider that T_0 is fixed and the difference between T and T_0 is small, then $f(T,T_0)$ has only one variable f(T) and can be write using the Taylor series. For order 1 we have:

$$m_f c_f \cdot \frac{dT}{dt} = P - (T - T_0) \frac{\partial f(T_0)}{\partial T}$$
(7)

Equation (7) is in fact the Newton equation law, and admits an analytical solution. If equation (7) is not sufficient to reproduce the experimental behavior, we can use the order 2 of the Taylor series of f(T):

$$m_f c_f \cdot \frac{dT}{dt} = P - (T - T_0) \frac{\partial f(T_0)}{\partial T} - \frac{(T - T_0)^2 \partial^2 f(T_0)}{2 \partial T^2}$$
(8)

This differential equation is more complex. We have chosen to solve this equation numerically, using the finite difference method:

$$T(t + \Delta t) = T(t) + \frac{\Delta t}{m_f c_f} \left(P - \left(T(t) - T_0 \right) \cdot A - \left(T(t) - T_0 \right)^2 \cdot B \right)$$
(9)
With, $A = \frac{\partial f(T_0)}{\partial T}$ and $B = \frac{1 \partial^2 f(T_0)}{2 \partial T^2}$.

It is interesting to note that:

$$m_f c_f = \frac{m_{NP} \cdot c_{V,f}}{C_{NP}} \tag{10}$$

with m_{NP} , $c_{V,f}$ and C_{NP} are the mass of the nanoparticles, the volume heat capacity of the fluid and the concentration of nanoparticles, respectively.

Equations (7) and (8) were used to fit the experimental data shown on Figures 3, S11, S12. The extracted SAR_{NP} from equation (7) and (8) are equal to 259 ± 1 W/g and 420 ± 1 W/g, respectively. The value of 420 ± 1 W/g will be held since the equation (8) provides the best fit. The parameters used for these fits are C_{NP} = 4 mg mL⁻¹ and $c_{V,f}$ = 1.538 J.mL⁻¹.K⁻¹.

The ILP may be calculated as:

$$ILP = \frac{SAR}{f \cdot H^2} \tag{10}$$

2. Fits for the ΔT as a function of the concentration curve for magneto- and photothermia experiment

The simplify heat equation for a magneto- or photothermal can be write as follows:

$$m_f c_f \cdot \frac{dT}{dt} = P - h \cdot (T - T_0) \tag{11}$$

where m_f and c_f are the mass and the weighted heat capacity of the fluid, respectively. The term *h* represents the energy exchange rate between the liquid and the environment. *P* represents the thermal power source, which is produced by the conversion of magnetic or light energy.

The term ΔT represents the difference between the initial temperature and the stabilized temperature when T reaches the plateau at t $\rightarrow \infty$. The plateau at t $\rightarrow \infty$ means that $dT(t \rightarrow \infty)/dt = 0$. Then if we call T_∞ the temperature reached on the plateau we have:

$$P - h \cdot (T_{\infty} - T_{0}) = 0$$
(12)

and so:

$$\Delta T = \frac{P}{h} \tag{13}$$

In the case of magnetothermia, P is directly proportional to the quantity of nanoparticles in the solution, due to the deep penetration of the magnetic field inside the solution. That means the variation of ΔT as a function of the nanoparticle concentration is linear.

In the context of photothermia, the absorption of light by the entire solution is observed to exhibit a nonlinear dependence on the concentration of nanoparticles. Considering a laser beam with a direction along the x-axis. The power intensity of the beam at the position x in the solution is defined by:

$$P_{beam}(x) = P_{laser} \cdot \exp\left(-a \cdot x\right) \tag{14}$$

with *a* the absorption coefficient, and P_{laser} the output power of the laser. The total absorbed power after the light has traversed a distance of *d* is given by the following equation:

$$P_{abs} = P_{laser} - P_{beam}(d) \tag{15}$$

If we develop, we have:

$$P_{abs} = P_{laser} \cdot (1 - \exp(-a \cdot d)) \tag{16}$$

The heat source power, as expressed in Equation (13), is proportional to the absolute pressure, P_{abs} , as demonstrated by the following relationship:

$$P = P_{abs} \cdot \eta \tag{17}$$

with η the light to heat conversion coefficient.

If we make the classical approximation that the coefficient absorption *a* is proportional to the nanoparticle concentration then we have:

$$a = \alpha \cdot c_{NP} \tag{18}$$

with α the concentration normalized coefficient absorption, and c_{NP} the concentration in mg/mL of nanoparticles in the solution.

Then, from the previous equations, we can deduce the variation of ΔT as a function of the nanoparticle concentration for a photothermal experiment.

$$\Delta T = \frac{\eta}{h} P_{laser} \cdot \left(1 - \exp\left(-\alpha \cdot d \cdot c_{NP} \right) \right)$$
(19)

3. Investigation of the homolysis kinetic of the radical release by EPR

A suspension of IONP@alkoxyamine nanoparticles **3** (10⁻³ M) in *tert*-butylbenzene was submitted to the action of an ac magnetic field (\approx 20 mT at a frequency of 350 kHz). At each timepoint, 100 µL was withdrawn and cooled in an ice bath to stop the reaction. After magnetic separation, 50 µL of the supernatant was analysed by EPR.

A correlation between intensities measured by EPR and concentrations of alkoxyamines was performed using a calibration at three concentrations (10^{-4} M, 10^{-5} M, and 10^{-6} M) of alkoxyamine in *tert*-butylbenzene, in order to obtain an intensity value (I_{ref}) for a reference concentration (C_{ref}). Those solutions had been incubated at 120 °C overnight and measured by RPE. From each supernatant analysed during the kinetic study, the intensity of the first peak is measured, and the conversion rate in radicals was estimated using this following equation:

$$\frac{C}{C_0} = \frac{I \times C_{ref}}{C_0 \times I_{ref}}$$

with $C_0 = 1$ mM is the initial concentration, and *I* is the intensity of the first peak measured by EPR at each timepoint. When plotted as a function of time, the kinetic constant k_d can be determined with the slope of the linear regression. The temperature at the surface of the nanoparticles can be estimated by using the Arrhenius law:

$$k_d = A \times exp(-E_a/RT)$$
 and
 $T = \frac{-E_a}{R \cdot Ln(\frac{k_d}{A})}$

where E_a is the activation energy, $A = 2.4 \cdot 10^{14} \text{ s}^{-1}$ is the pre-exponential factor in the case of the alkoxyamine homolysis, R is the perfect gas constant and k_d is a homolysis constant.