

Supplementary Information to "C atoms versus Si atoms at the bridgehead positions of phenyl-decorated adamantane-type clusters: Influence on the nonlinear optical response"

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simone.sanna@theo.physik.uni-giessen.de

Atom coordinates

Second entry for each structure is from a HSE06 calculation.



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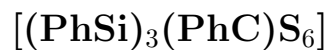
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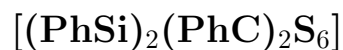
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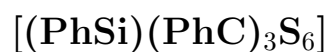
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Si H C S
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Direct

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0.5157083772565547	0.3606634645825662	0.5569077782692834
0.5189433093334989	0.3657745308344037	0.5937285407810473
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ATOMIC_POSITIONS angstrom

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H	19.7990548634	20.4555297983	15.2628496815
H	21.6511746275	21.2619630122	13.8256285403
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H	24.4693676923	19.0846347018	16.2537825668
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 C 20.8257269575 20.1590408457 15.4782714336
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 C 16.7194259418 22.0683173599 20.2807654031
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 C 17.7307272328 23.7357852884 21.7180079939
 C 18.9388008805 23.0345207420 21.6699076786
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 C 19.8010436038 16.6260160485 22.1742229299
 C 15.5456423931 17.2156229548 16.6893118595
 C 14.7105519916 16.2319455265 17.2478066179
 C 13.5791571291 15.7905994054 16.5597440979
 C 13.2686682094 16.3287606267 15.3079955862
 C 14.0899808093 17.3084482879 14.7432256085
 C 15.2223804710 17.7505136132 15.4292625324
 S 19.4213486844 20.4679401556 18.2335497326

S 20.4761795809 17.6364234463 18.6463875494
S 18.6358137449 18.1847716684 16.2520990865
S 18.5669494429 18.8442311753 20.7059425904
S 16.4864624443 19.6495977776 18.5864474650
S 17.6412934904 16.3970568899 19.0683381006

Si H C S
1 20 27 6

Direct

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0.6088909470025584 0.4932271742171996 0.4770954517809838
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0.5672042334313617 0.3574326855353835 0.5593322418466102
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0.6027035113416511 0.5101766770126395 0.4542073560920220
0.4829544042954723 0.5748048686101996 0.5445594449282362

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0.5127306455949346	0.5881013432856865	0.5633763182849120
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0.5481287536229187	0.4229380128516924	0.6196624546338501
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ATOMIC_POSITIONS angstrom

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H	19.7990548634	20.4555297983	15.2628496815
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H	24.0002949810	20.5682010238	14.3060187753
H	24.4693676923	19.0846347018	16.2537825668
H	22.6340294509	18.3123151966	17.7049279558
H	15.8457949785	21.7010111251	19.7437666204
H	15.6762934075	23.7868083581	21.0517307049
H	17.6536115435	24.6576016929	22.2957205890
H	19.8110692071	23.4053693994	22.2090488099
H	19.9800410272	21.3103360399	20.8930295876
H	20.0707023318	14.9043682211	19.2505093163
H	21.0827524374	13.2371727233	20.7803785616

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H C S

20 28 6

Direct

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0.5975810587040813	0.4707638289645169	0.6259681171441885
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0.4354791527663160	0.5910561246460271	0.5149226095920394
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0.3901574825063558	0.5663402689161375	0.5689339059805694
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HOMO / LUMO Plots

The isosurface for the HOMO / LUMO plots is $0.001 \text{ eV}/\text{\AA}^3$

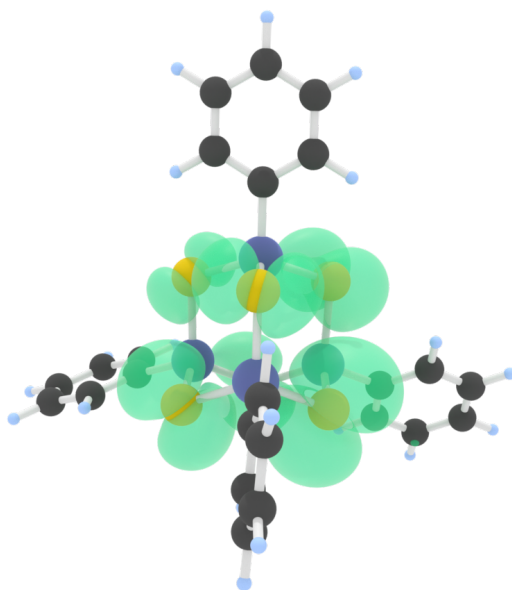


Figure 1: HOMO of $[(\text{PhSi})_4\text{S}_6]$

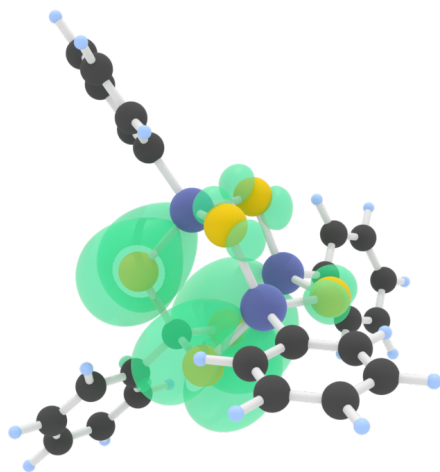


Figure 2: HOMO of $[(\text{PhSi})_3(\text{PhC})\text{S}_6]$

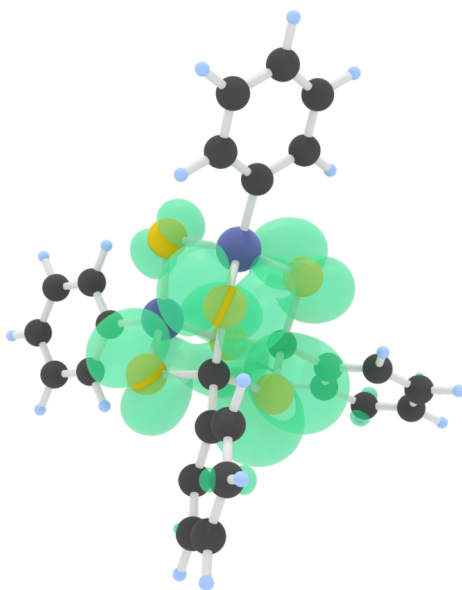


Figure 3: HOMO of $[(\text{PhSi})_2(\text{PhC})_2\text{S}_6]$

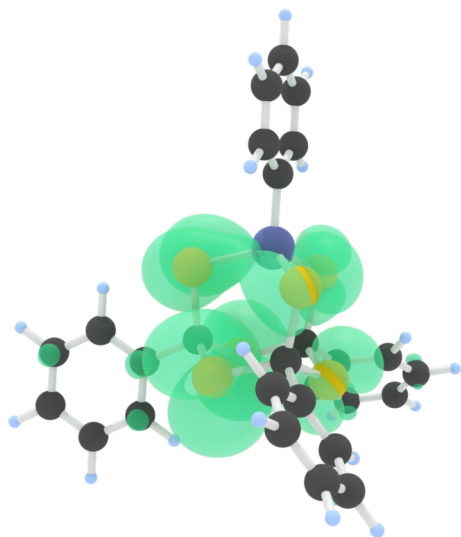


Figure 4: HOMO of [(PhSi)(PhC)₃S₆]

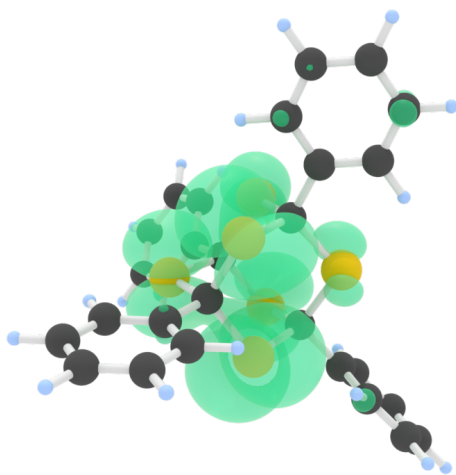


Figure 5: HOMO of [(PhC)₄S₆]

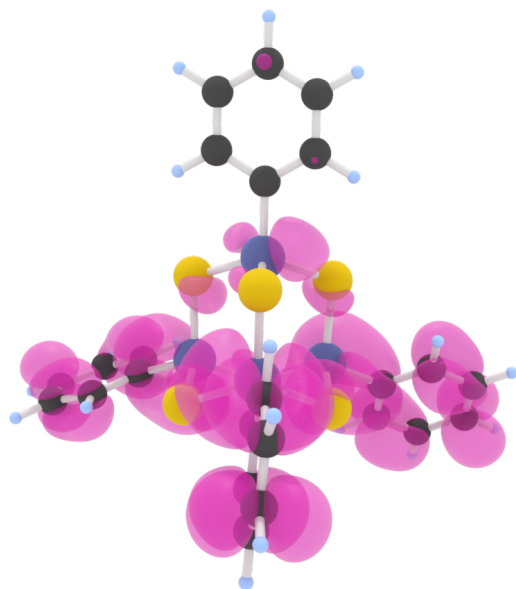


Figure 6: LUMO of $[(\text{PhSi})_4\text{S}_6]$

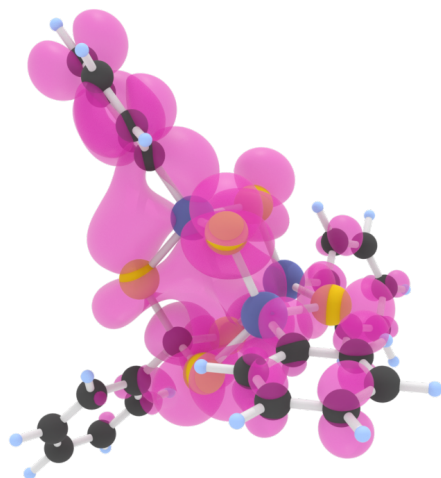


Figure 7: LUMO of $[(\text{PhSi})_3(\text{PhC})\text{S}_6]$

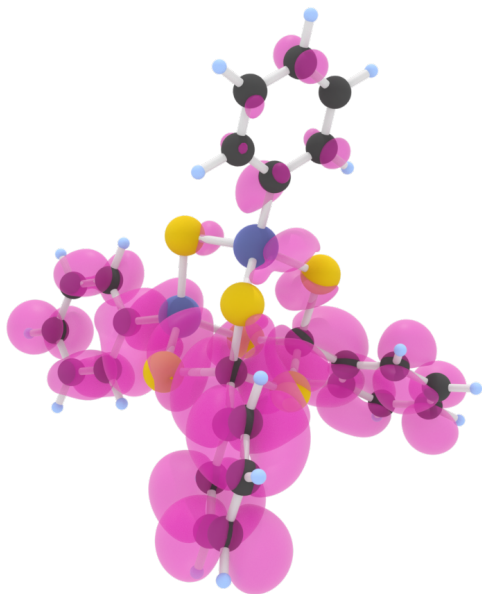


Figure 8: LUMO of [(PhSi)₂(PhC)₂S₆]

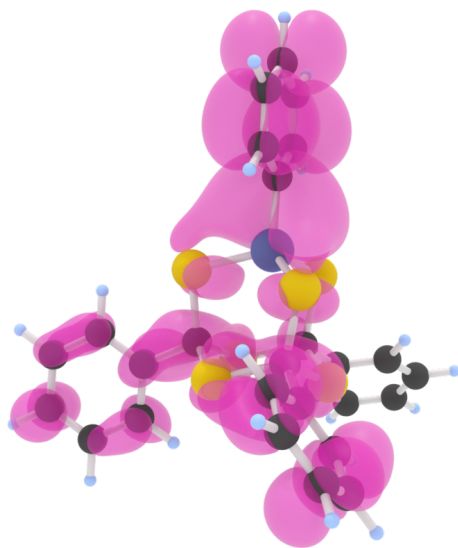


Figure 9: LUMO of [(PhSi)(PhC)₃S₆]

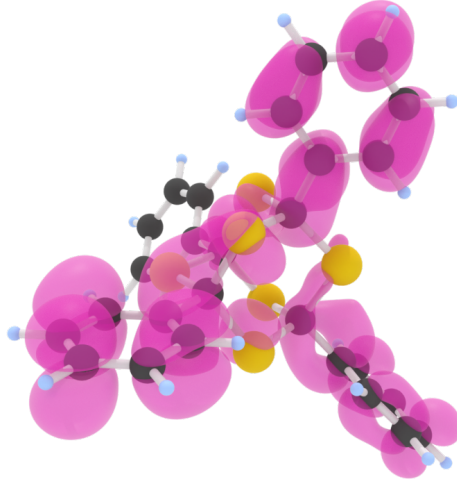


Figure 10: LUMO of $[(\text{PhC})_4\text{S}_6]$

Computational approach and main computational parameters for non-linear optical calculations.

Within the present approach, the nonlinear optical susceptibility $\chi^{(n)}$ is obtained in the time domain from the dynamical polarization $\vec{P}$ of the system, using the expansion of $\vec{P}$ in power of the applied field $\vec{E}$. The Cartesian component α reads

$$P_\alpha = \sum_{\beta} \chi_{\alpha\beta}^{(1)} E_\beta + \sum_{\beta,\gamma} \chi_{\alpha\beta\gamma}^{(2)} E_\beta E_\gamma + \sum_{\beta,\gamma,\delta} \chi_{\alpha\beta\gamma\delta}^{(3)} E_\beta E_\gamma E_\delta + \dots \quad (1)$$

The starting point of the calculation is the the zero-field Kohn-Sham Hamiltonian

$$\hat{H}^0 = -\frac{\hbar^2}{2m} \sum_i \nabla_i^2 + \hat{V}_{\text{el}} + \hat{V}_{\text{H}}[\rho_0] + \hat{V}_{\text{xc}}[\rho_0], \quad (2)$$

where $\hat{V}_{\text{el}}$ is the electron-ion interaction, $\hat{V}_{\text{H}}$, and $\hat{V}_{\text{xc}}$ are the Hartree and the exchange-correlation potentials, respectively.

The (external, monochromatic) laser excitation

$$\vec{E}(t) = \vec{E}_0 \sin(\omega_L t) \quad (3)$$

with the frequency ω_L is used for the definition of the non-hermitian field coupling operator

$$\hat{w}_{\vec{k}} = \frac{ief}{4\pi} \sum_m \sum_{\alpha=1}^3 (\vec{a}_\alpha \cdot \vec{E}) N_{\vec{k}_\alpha} \sum_{\sigma=\pm} \sigma |\tilde{v}_{\vec{k}_\alpha, m}^\sigma\rangle \langle v_{\vec{k}, m}|. \quad (4)$$

In this equation

$$|\tilde{v}_{\vec{k}_\alpha, n}^\pm\rangle = \sum_m (S(\vec{k}, \vec{k}_\alpha^\pm)^{-1})_{mn} |v_{\vec{k}_\alpha, m}^\pm\rangle. \quad (5)$$

The matrix S is built up by the overlaps of the periodic part of the occupied Bloch functions $v_{\vec{k}, m}$:

$$S_{mn}(\vec{k}, \vec{k} + \vec{q}_\alpha) = \langle v_{\vec{k}, m} | v_{\vec{k} + \vec{q}_\alpha, n} \rangle. \quad (6)$$

The latter are obtaining by solving the equation of motions

$$i\hbar \frac{d}{dt} |v_{\vec{k}, m}\rangle = (\hat{H}_{\vec{k}}^0 + \hat{w}_{\vec{k}}^\dagger(\vec{E})) |v_{\vec{k}, m}\rangle \quad (7)$$

in the time domain. Finally, the polarization

$$\vec{P}_\alpha = -\frac{ef}{2\pi v} \frac{\vec{a}_\alpha}{N_{\vec{k}_\alpha}^\perp} \sum_{\vec{k}_\alpha^\perp} \text{Im} \left[\sum_{i=1}^{N_{\vec{k}_\alpha^\perp}^\perp - 1} \text{Tr} \ln S(\vec{k}_i, \vec{k}_i + \vec{q}_\alpha) \right] \quad (8)$$

is calculated, from which the non-linear susceptibility can be obtained. The procedure is repeated for all required excitation frequencies.

In this computational procedure, general parameters were converged for the SHG response. Effective parameters chosen were identified through multiple convergence tests

weighting the numerical cost to accuracy of the results. Orbitals are set by an interval including a factor for the lowest and highest orbital. The interval of (0.2, 0.8) orbitals when the interval of (0, 1) orbitals range from 1 to $3 \times$ number of occupied orbitals for frequencies with energies ≥ 5 eV, linear interpolated between $2 \times$ number of occupied orbitals, and $3 \times$ number of occupied orbitals for energies within (2 eV, 3 eV), and $2 \times$ number of occupied orbitals for energies ≥ 0 eV. G-Vectors were reduced down to a minimum of ≈ 0.014 times, and in their maximum 0.5 times an energy cutoff of 80 Ry. Occasionally these parameters were/could be tweaked as they tend to overestimate the necessary accuracy needed for intended evaluation. The volumes within the molecules are simulated are 51179,6 Å³ dependent on molecule size. Field intensity was chosen to be 0.100000E+8 kW/cm².

Further numerical parameters

The spectral broadening within YAMBO is set to 0.2 eV.

Further computational parameters vary for each structure, due to the stepwise process.

Therefore, we give a range (min,max) for each structure.

[(PhSi)₄S₆]

Simulation times: 43.2999990474 fs , 43.2999990474 fs

Step time: 0.00999999978 fs , 0.00999999978 fs

Kinetic-energy cutoff (QE): 1088.4554498392001 eV

[(PhSi)₃(PhC)S₆]

Simulation times: 43.2999990474 fs , 43.2999990474 fs

Step time: 0.00999999978 fs , 0.00999999978 fs

Kinetic-energy cutoff (QE): 1088.4554498392001 eV

[(PhSi)₂(PhC)₂S₆]

Simulation times: 43.2999990474 fs , 43.2999990474 fs

Step time: 0.00999999978 fs , 0.00999999978 fs

Kinetic-energy cutoff (QE): 1088.4554498392001 eV

$[(\text{PhSi})_1(\text{PhC})_3\text{S}_6]$

Simulation times: 43.2999990474 fs , 43.2999990474 fs

Step time: 0.00999999978 fs , 0.00999999978 fs

Kinetic-energy cutoff (QE): 1088.4554498392001 eV

$[(\text{PhC})_4\text{S}_6]$

Simulation times: 43.2999990474 fs , 43.2999990474 fs

Step time: 0.00999999978 fs , 0.00999999978 fs

Kinetic-energy cutoff (QE): 1088.4554498392001 eV

k-Point Grid

The Gamma point sampling was employed for the structural optimization of the isolated clusters. A regular 2x2x2 k-point mesh is furthermore needed for the converged calculation of the dynamical polarization and was thus employed for the calculation of the optical response of all investigated structures.

Coloring of the orbital (overlaps)

The overlap $(\rho_1 \cdot \rho_2)$ in the range from the numerical values $0.0 \left(\frac{\text{eV}}{\text{\AA}^3}\right)^2$ to $0.05 \left(\frac{\text{eV}}{\text{\AA}^3}\right)^2$ is colored from "low" to "high" by the RGB values from (1.0, 1.0, 1.0) to (0.0, 0.0, 0.0). This color code is The presented isosurfaces is the $\frac{1}{32}$ of the maximum value of the distribution.

Visual representation of the core distortion

The bridgehead atoms $\text{Si}_{4-x}\text{C}_x$ $x = 0, \dots, 4$ (i.e., the core skeleton of the investigated compounds) define a more or less distorted tetrahedron (or rather a trigonal pyramid), as shown in the following picture on the example of the $[(\text{PhSi})_4\text{S}_6]$ molecular cluster.

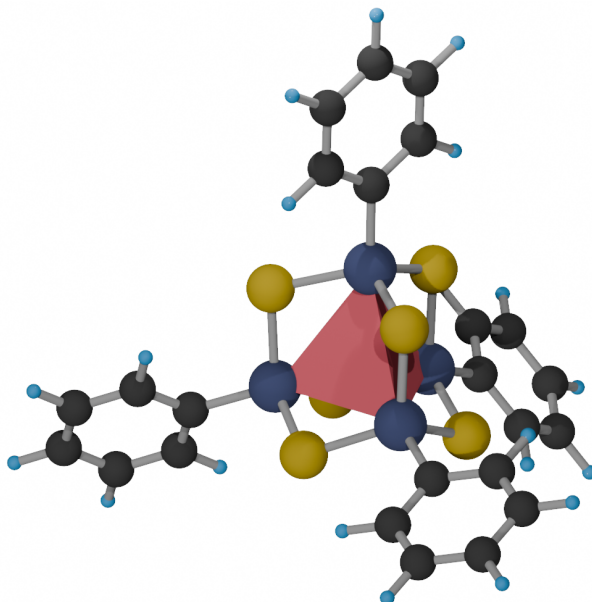


Figure 11: $[(\text{PhSi})_4\text{S}_6]$ molecular cluster, with the tetrahedron described by the bridgehead atoms highlighted in the cluster core.

Structural distortions can be visualized, e.g., by the deviation of the calculated bridgehead atom positions from an ideal tetrahedron. Keeping the color coding employed in Fig. 2 or 3 of the manuscript ($[(\text{PhC})_4\text{S}_6]$ in **violet**, $[(\text{PhSi})(\text{PhC})_3\text{S}_6]$ in **red**, $[(\text{PhSi})_2(\text{PhC})_2\text{S}_6]$ in **green**, $[(\text{PhSi})_3(\text{PhC})\text{S}_6]$ in **orange**, and $[(\text{PhSi})_4\text{S}_6]$ in **blue**) the visual comparison of the tetrahedron distortion is obtained:



Figure 12: Comparison of the tetrahedra described by the bridgehead atoms highlighted in the cluster core.

The shaded region represents an ideal, regular tetrahedron with the volume of the pyramid inscribed in $[(\text{PhC})_4\text{S}_6]$ (in violet), showing that nearly no distortion occurs for this compound. The same holds for $[(\text{PhSi})_4\text{S}_6]$ (in blue). All other structures show a more or less pronounced distortion, instead.

Hybrid-DFT

Further calculations using HSE06 hybrid functional were conducted to address the question whether a more precise description of the electronic structure would change the interpretations/conclusions made in the manuscript.

The effect of the hybrid calculations on the electronic structure is both to open the HOMO-LUMO gap (see table below) and to slightly broaden the energy window in which

orbital groups are distributed. This is illustrated on the example of the (occupied) energy levels we associate with the first SHG peak.

Table 1: HOMO-LUMO gaps as calculated within DFT-PBE and HSE06 for the investigated compounds. All values are given in eV.

Compound	PBE [eV]	HSE06 [eV]
$[(\text{SiPh})_4\text{S}_6]$	3.80	5.06
$[(\text{SiPh})_3(\text{CPh})\text{S}_6]$	3.28	4.56
$[(\text{SiPh})_2(\text{CPh})_2\text{S}_6]$	3.25	4.53
$[(\text{SiPh})(\text{CPh})_3\text{S}_6]$	3.20	4.46
$[(\text{CPh})_4\text{S}_6]$	3.15	4.44

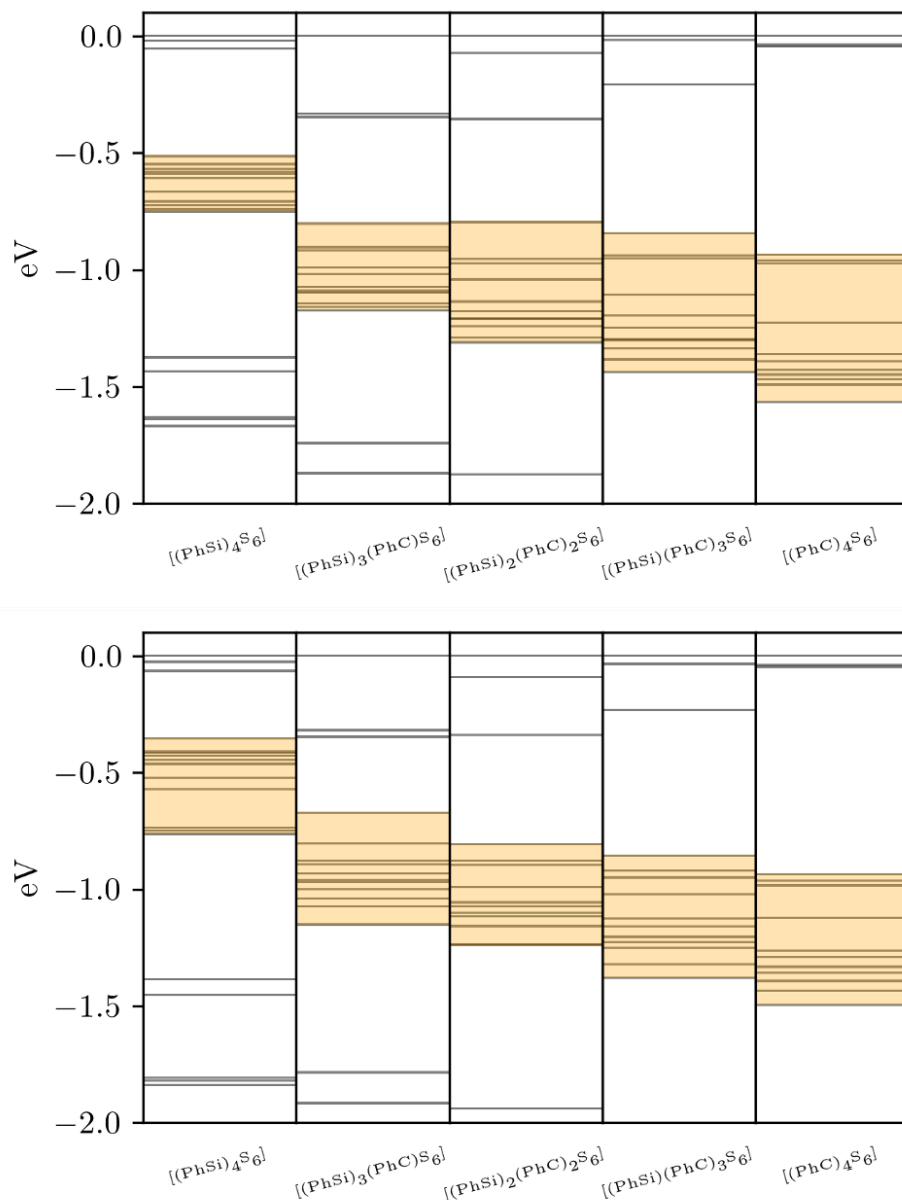


Figure 13: Occupied electronic levels of the investigated molecular clusters as calculated within DFT-PBE (upper panel) and hybrid-DFT in the HSE06 formulation. The orbital group associated to the SHG response in the region below 3 eV is highlighted. It contains 11 individual electronic levels and widens from left to right. The same 11 levels can be discriminated both in the PBE and HSE06 calculations. The origin of the energy scale is the HOMO level.

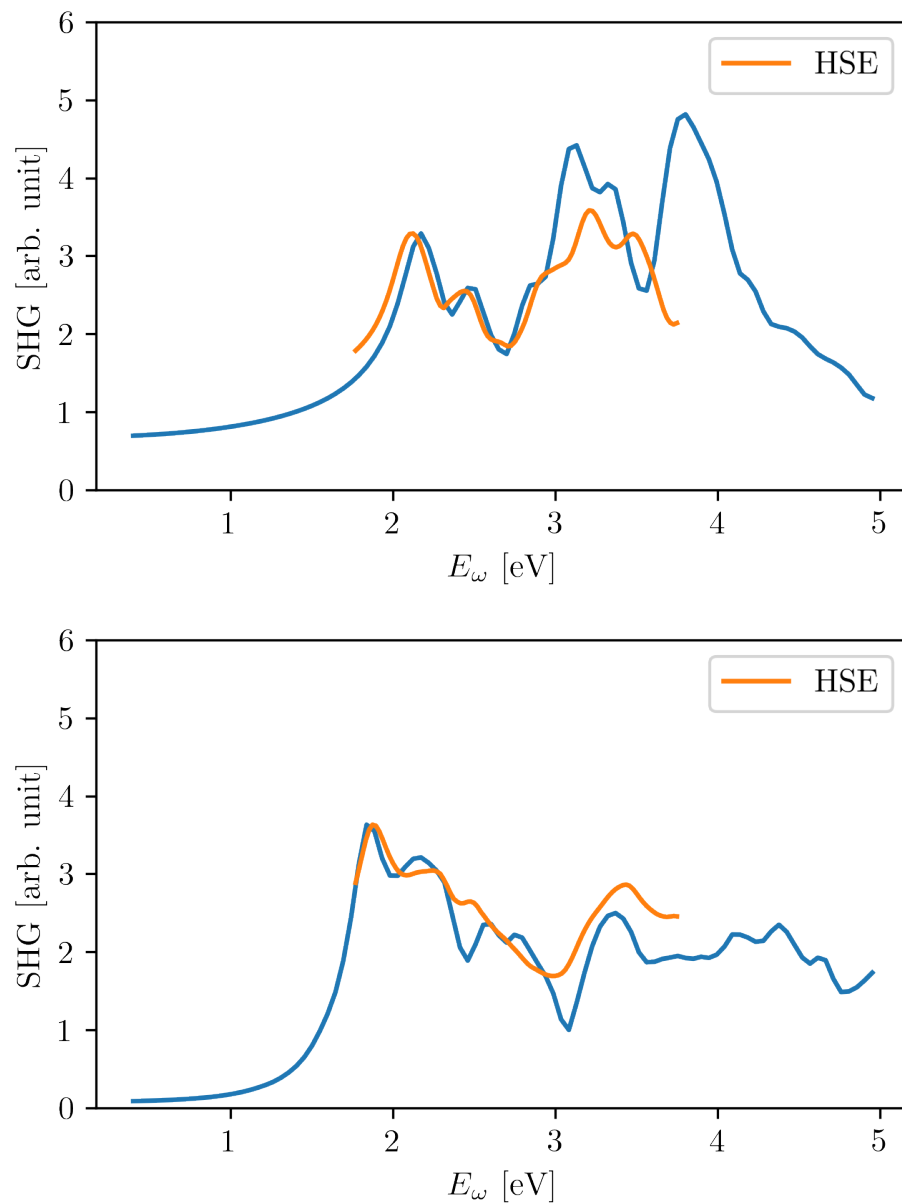


Figure 14: Comparison of the SHG spectra as calculated within DFT-PBE (blue line) and hybrid-DFT in the HSE06 formulation (orange line) for $[(\text{SiPh})_4\text{S}_6]$ (upper panel) and $[(\text{CPh})_4\text{S}_6]$ (lower panel). The HSE06 spectra are shifted by a value corresponding to the HOMO-LUMO gap difference between PBE and HSE06 for a better comparison.