

## Process-Dependent Hypersonic Phonon Dispersion of Brush Particle Metamaterials

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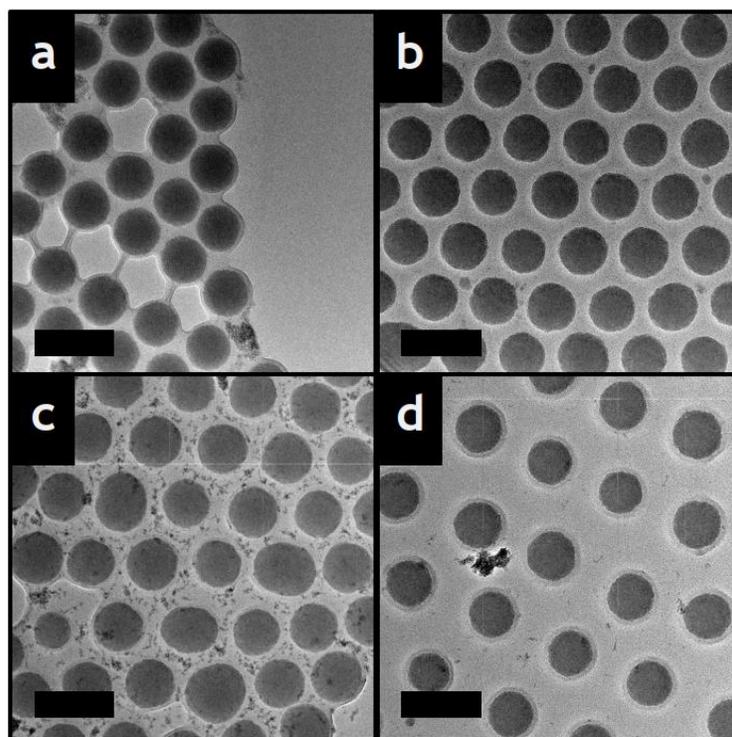
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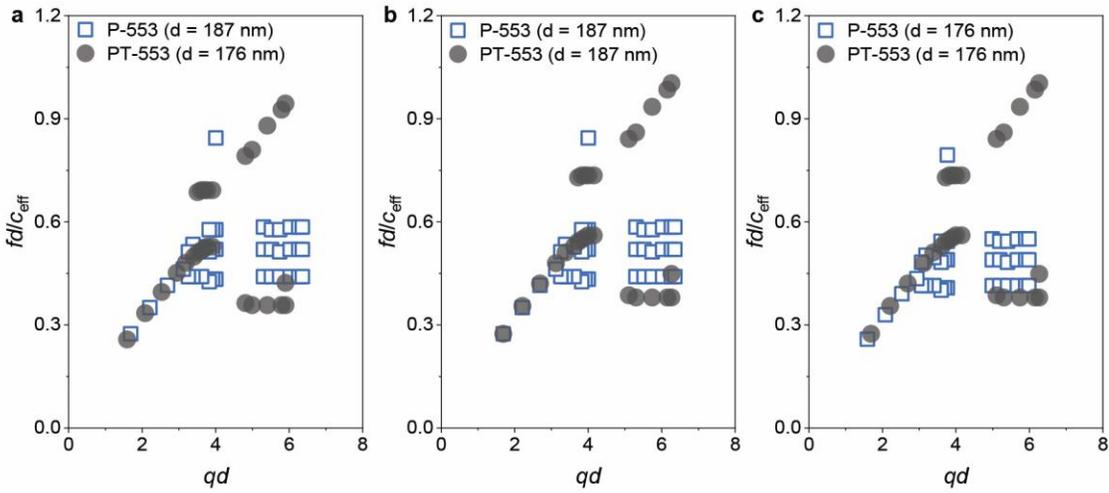
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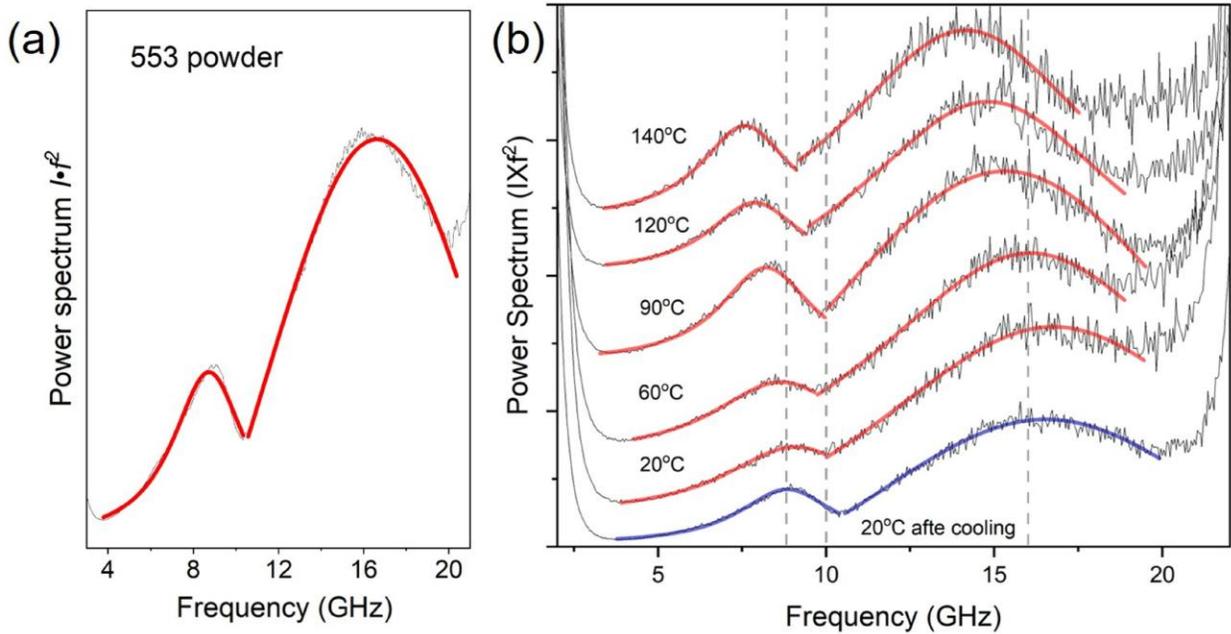
### Supporting Information



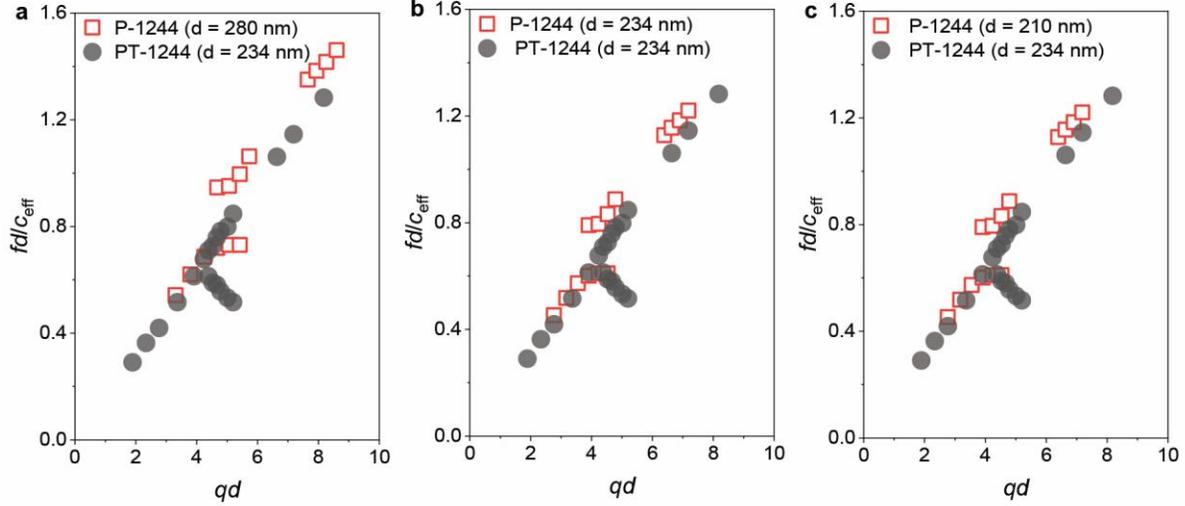
**Fig. S1.** Transmission electron microscope (TEM) images of two of the SiO<sub>2</sub>-PMMA samples (N=553 and N=1244) in (c) and (d); the other two are for (a) N=254 and (b) N=337. Scale bars are 200 nm. TEM was utilized to estimate the spacing  $d$  and polymer thickness  $h$  in 4<sup>th</sup> and 6<sup>th</sup> column of Table 1, respectively.



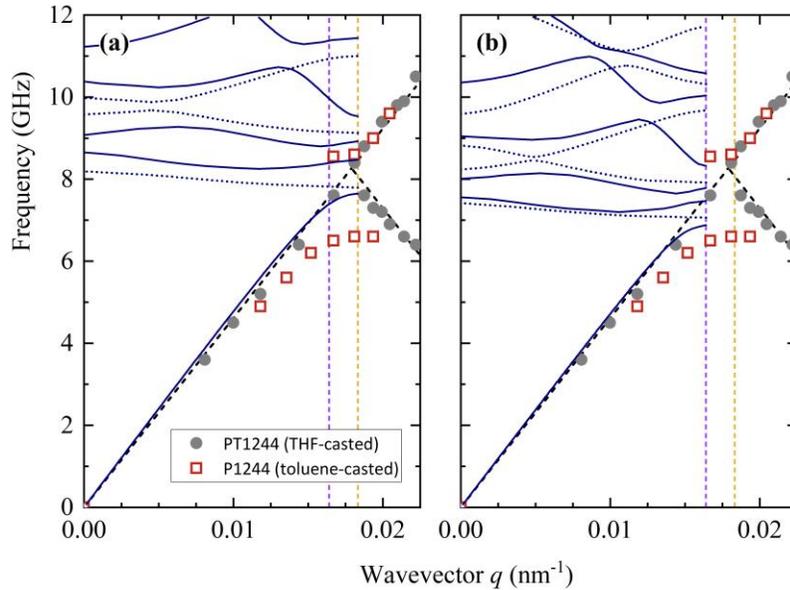
**Fig. S2:** Reduced band diagram  $fd/c_{\text{eff}}$  vs  $qd$  for P-553 (open squares) and PT-553 (solid symbols) using  $c_{\text{eff}} = 2590$  m/s and  $3000$  m/s, respectively. **(a)**  $d = 187$  nm for the P-553 and  $d = 176$  nm for PT-553. **(b)**  $d = 187$  nm and **(c)**  $d = 176$  nm for both systems. A good superposition of the acoustic branch is obtained in all cases, indicating that the long-wavelength regime is insensitive to the choice of diameter and mainly determined by the experimental  $c_{\text{eff}}$ . For PT-553, a band gap appears at  $fd/c_{\text{eff}} \approx 0.65$  and  $qd \approx 3.8$ . In contrast, superposition of the flat modes is achieved only when the TLM diameter  $d = 176$  nm is assumed for both films, highlighting the sensitivity of the short-wavelength regime to particle size and local structural differences related to metastability. Due to the later, it is conceivable that the sound velocity in the cluster of PGN's is lower than  $c_{\text{eff}}$  rendering the superposition of the band diagram incomplete.



**Fig. S3.** Power vibration spectrum for P533 powder at **(a)**  $20$  °C and **(b)** different temperatures below and above  $T_g$  ( $\sim 110$  °C) of PMMA. The low-frequency modes in Figure 2b yielding to a broad peak at  $\sim 8.7$  GHz in the power spectrum in **(a)** are attributed to inter-PGN interactions. The spectra of solvent-cast films in **(b)** remain essentially unchanged upon annealing up to  $140$  °C (above  $T_g$ ) and after cooling, indicating robust and reversible PGN interactions associated with trapped, metastable configurations.



**Fig. S4.** Reduced band diagram  $fd/c_{\text{eff}}$  vs  $qd$  for  $\text{SiO}_2$ -PMMA-1244 drop casted from THF (P-1244) and toluene (PT-1244) solutions. (a) For PT-1244,  $d=234$  nm and for P-1244,  $d=280$  nm (Table 1). (b) For both PT-1244 and P-1244,  $d=234$  nm. (c) For PT-1244,  $d=234$  nm and for P-1244,  $d=210$  nm. Using  $d=280$  or  $234$  nm for P-1244 fails to reproduce the narrow Bragg gap of PT-1244 at  $fd/c_{\text{eff}} \approx 0.65$ ,  $qd \approx 4$  and the gap is shifted to higher reduced frequency values. A satisfactory superposition is obtained only when a smaller effective local diameter ( $d_{\text{loc}} \approx 210$  nm) is assumed for P-1244, consistent with the TLM prediction.



**Fig. S5.** Calculated band structure assuming core-to-core (first-neighbour) distance (a)  $d=210$  nm (TLM value) and (b)  $d=234$  nm (TEM value), along  $[111]$  direction of  $fcc$  lattice, assuming PBC for silica particles embedded in PMMA host, compared to the experimental data for both PT-1244 and P-1244 films. Coloured (orange and violet) vertical dashed lines indicate the corresponding BZ edges. Blue solid/dotted lines: longitudinal/deaf (inactive) modes.